

Topological geons with self-gravitating phantom scalar field

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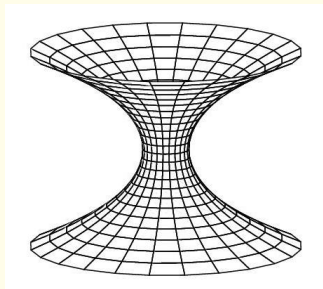
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Topological geons

A topological geon is the quotient manifold $G = M/\mathbb{Z}_2$ where M is a static spherically symmetric wormhole having the reflectional symmetry with respect to its throat.

The \mathbb{Z}_2 action on M is defined in the spherical coordinates by $(t, r, \theta, \varphi) \rightarrow (t, -r, \pi - \theta, \varphi + \pi)$.

In what follows, we assume that the topological geon's spacetime is asymptotically flat.



References



Sorkin R.D. Introduction to topological geons. *In: Proceedings of the Topological Properties and Global Structure of Space-Time*, ed. by P.G. Bergmann and V. de Sabbata. Erice, Italy, May 12-22, 1985. Pp. 249-270.



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Sakellariadou M. Production of Topological Defects at the End of Inflation. *Lecture Notes in Physics*, 2008, vol. 738, pp. 359-392. Available at: [arXiv:hep-th/0702003](https://arxiv.org/abs/hep-th/0702003)

Action

$$\Sigma = \frac{1}{8\pi} \int \left(-\frac{1}{2}S + \varepsilon \langle d\phi, d\phi \rangle - 2V(\phi) \right) \sqrt{|g|} d^4x,$$

where S — scalar curvature, $\varepsilon = \pm 1$, $V(\phi)$ is a self-interaction potential, and the angle brackets denote the scalar product with respect to the metric.

For topological geons $\varepsilon = -1$.

Energy condition $T_{ij}U^iU^j \geq 0$ does not hold for wormholes (see [Morris M.S., Torn K.S. *Wormholes in space-time and their use for interstellar travels* // Am. J. Phys., 1988, V.56, p. 395-412](#))

Metric and the quadratures

$$ds^2 = A^2 dt^2 - \frac{dr^2}{A^2} - C^2 (d\theta^2 + \sin^2 \theta d\varphi^2) \quad (1)$$

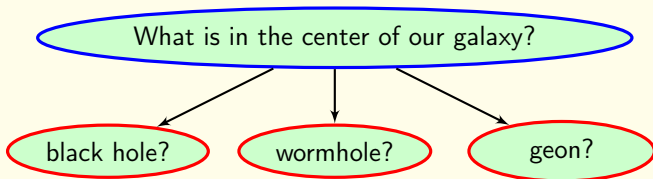
$$\phi' = \sqrt{-\varepsilon C''/C}, \quad A^2 = 2C^2 \int_r^\infty \frac{r-3m}{C^4} dr \quad (2)$$

$$\tilde{V}(r) = \frac{1}{2C^2} \left(1 - 3C'^2 A^2 - CC'' A^2 + 2C' \frac{r-3m}{C} \right) \quad (3)$$

For topological geons we assume $\varepsilon = -1$, $m = 0$, C — even function, $C'' \geq 0$, $C = r + o(1)$, $r \rightarrow \infty$.

Physical interpretation

Topicality of current research is explained by one of the tasks of the project RadioAstron.



How can we distinguish a black hole from a symmetric wormhole or a geon?

	black hole	geon
Distance orbits	Yes	No ($A^2 \sim 1 - \frac{a}{r^4}, r \rightarrow \infty$)
Angular velocity in ISCO ¹	$\omega \neq 0$	$\omega = 0$

¹Innermost Stable Circular Orbit

Classification of topological geons

The motion of test particles are defined by the equation

$$\left(\frac{dr}{ds}\right)^2 = E^2 - V_{\text{eff}}, \quad (4)$$

where $V_{\text{eff}} = A^2 \left(k + \frac{J^2}{C^2}\right)$, $k = 0, 1, -1$, E and J are constants (energy and angular momentum).

Classification

1. A^2 has a maximum at $r = 0$ and has no minima;
2. A^2 has a unique minimum at $r = 0$;
3. A^2 has a minimum at $r \neq 0$.

Instances

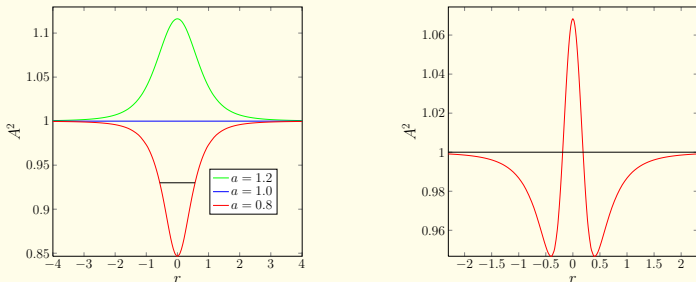


Figure: *Left:* Graphics of A^2 for $C = (r^4 + 2r^2 + a^4)^{1/4}$. Geon with $a = 1.2$ corresponds to the case 1 and has no any stable orbit. Geon with $a = 0.8$ corresponds to the case 2 and the motion of test particles is possible in the some neighborhood of topological threshold. $a = 1$ define Ellis solution $C = \sqrt{r^2 + a^2}$, $A^2 = 1$, $\phi = \arctan \frac{r}{a}$ with no one stable orbit. *Right:* Graphic of A^2 for $C = (r^6 + c^2r^4 + b^4r^2 + a^6)^{1/4}$ with $a = b = 0.5, c = 1$. This geon corresponds to the case 3.

Finite motion for the case 3 (I)

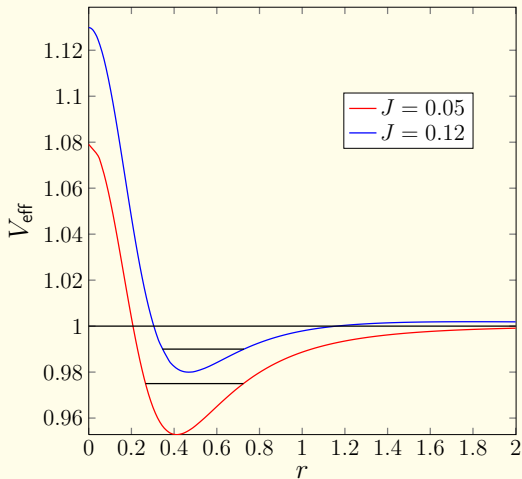


Figure: Graphics of V_{eff} for $C = (r^6 + c^2 r^4 + b^4 r^2 + a^6)^{1/4}$ with $a = b = 0.5$, $c = 1$.

Finite motion for the case 3 (II)

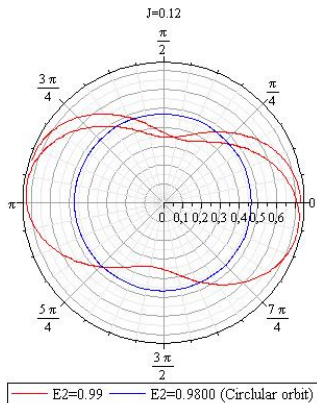
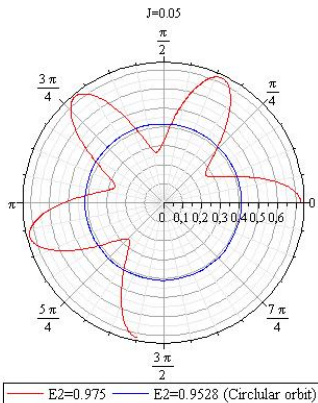


Figure: Orbits for geon with V_{eff} for $C = (r^6 + c^2 r^4 + b^4 r^2 + a^6)^{1/4}$ with $a = b = 0.5$, $c = 1$.