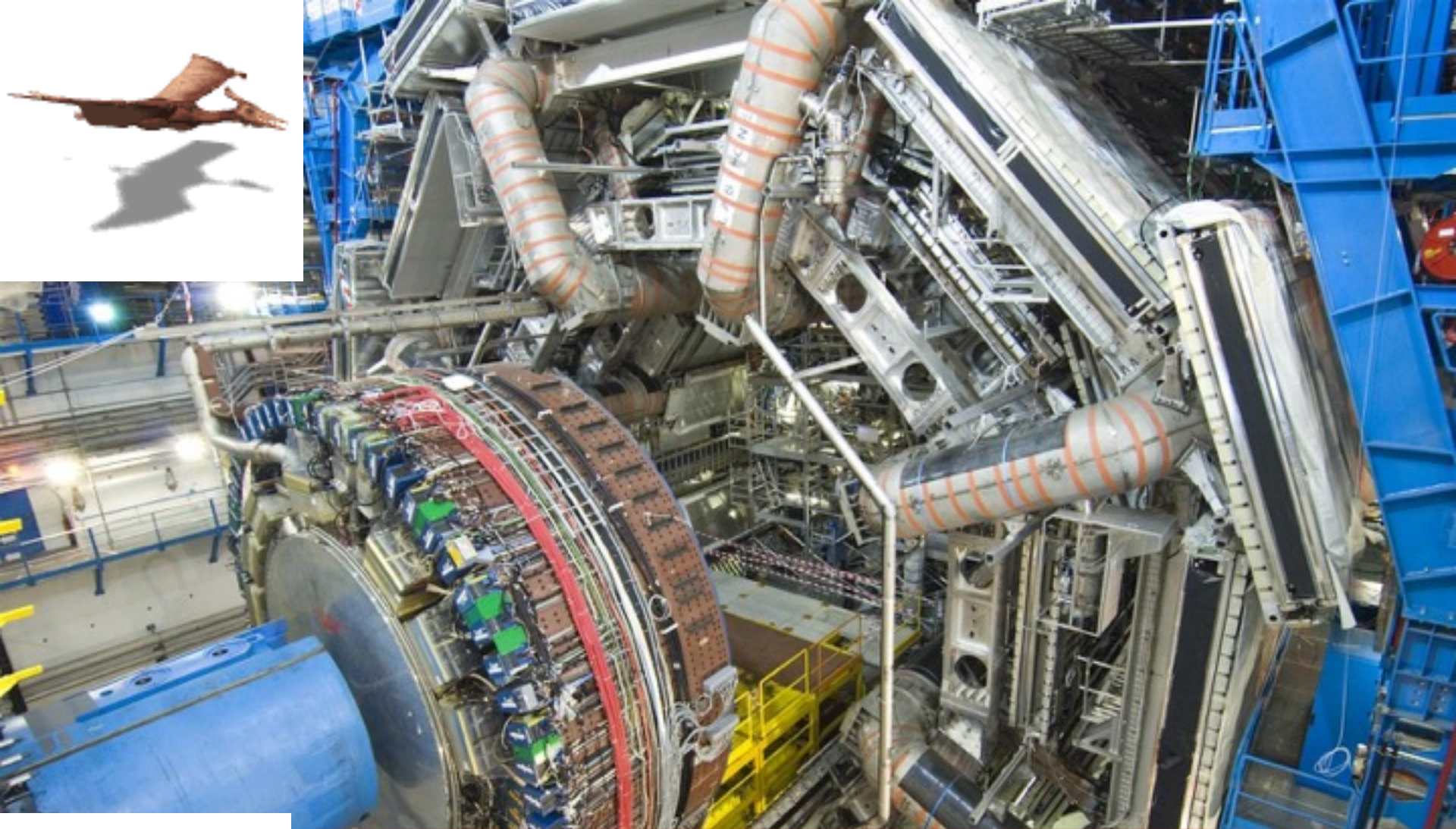


Physics at the LHC: a historical perspective and a pedagogical view on how this is done. The role of theory and experiment

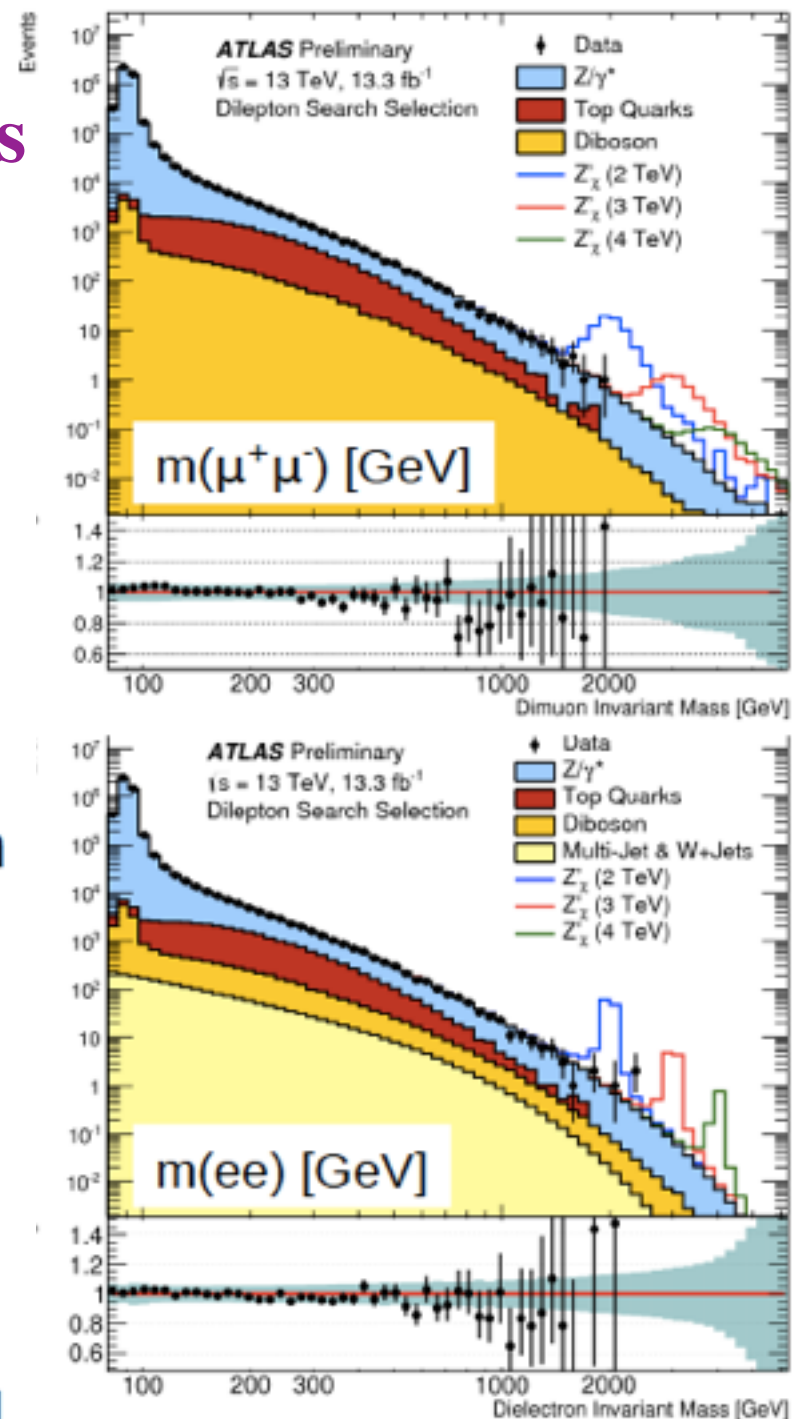


D. Froidevaux (CERN)

Physics at the LHC: a historical perspective and a pedagogical view on how this is done. The role of theory and experiment

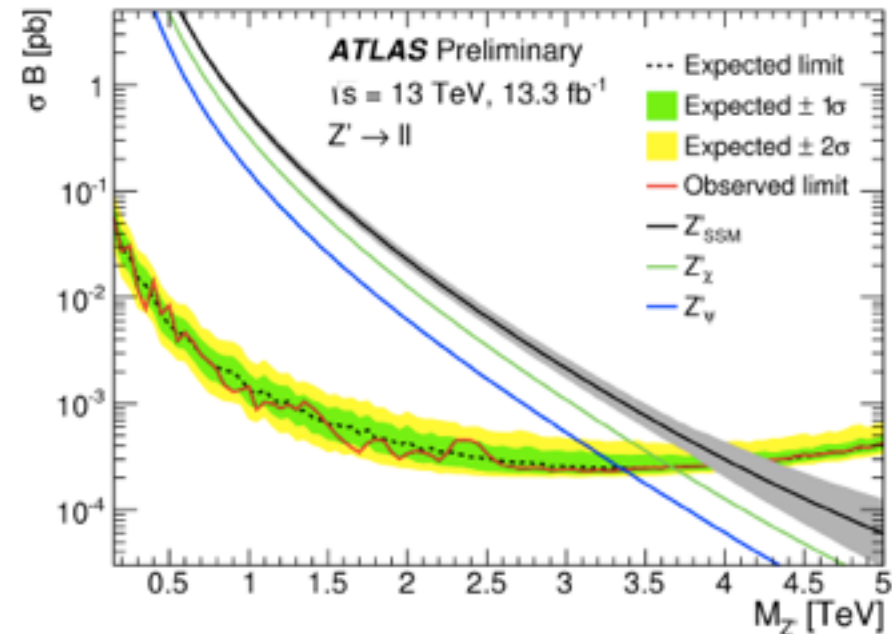
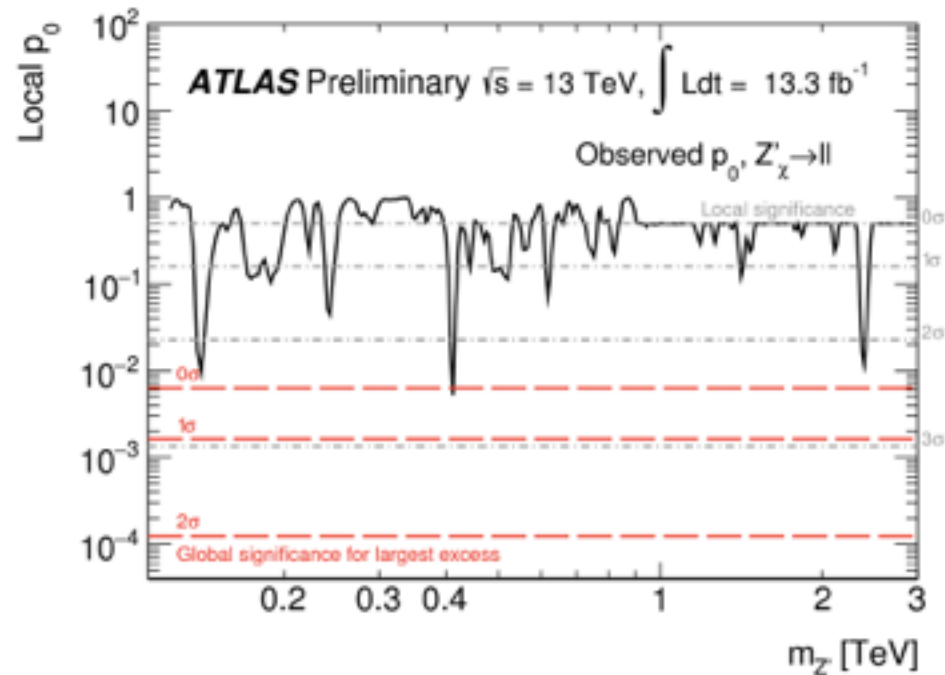
Search for high-mass resonances decaying to leptons

- Dimuon channel:
 - 30 μm muon spectrometer alignment critical (ATLAS)
 - Resolution 10-15% at $p_T = 1 \text{ TeV}$
- Dielectron channel:
 - Excellent resolution: $< 2\%$ at high momentum
 - Poor charge measurement \rightarrow no charge requirement
- Fit of the entire dilepton spectrum, incl. Z peak.



Search for Heavy Resonance: dilepton channel

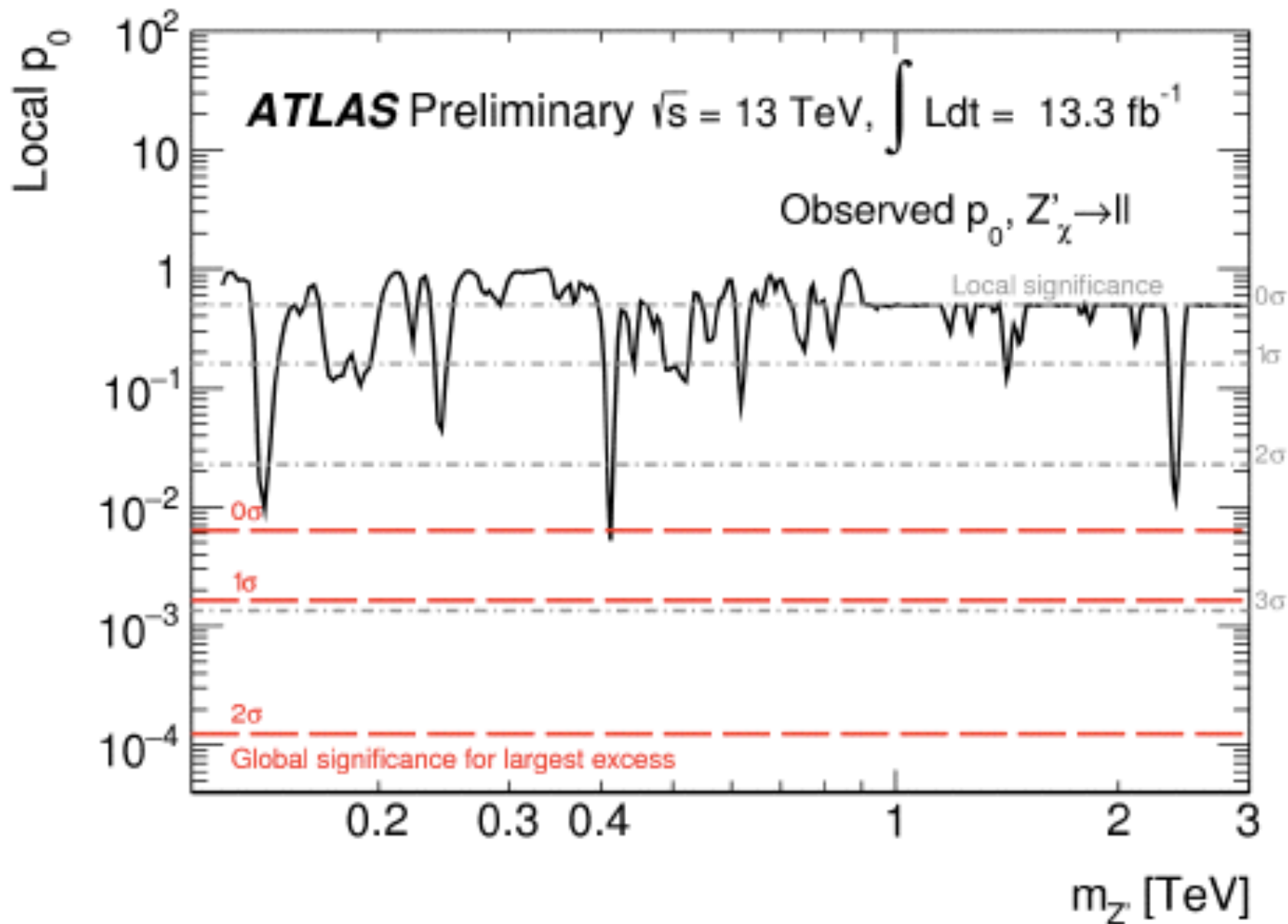
- Search for narrow resonance with binning optimized wrt detector resolution from 120 GeV to 5 TeV



Observed lower limit (TeV) at 95% CL:
 $m(\text{SSM } Z') > 4.05$ TeV

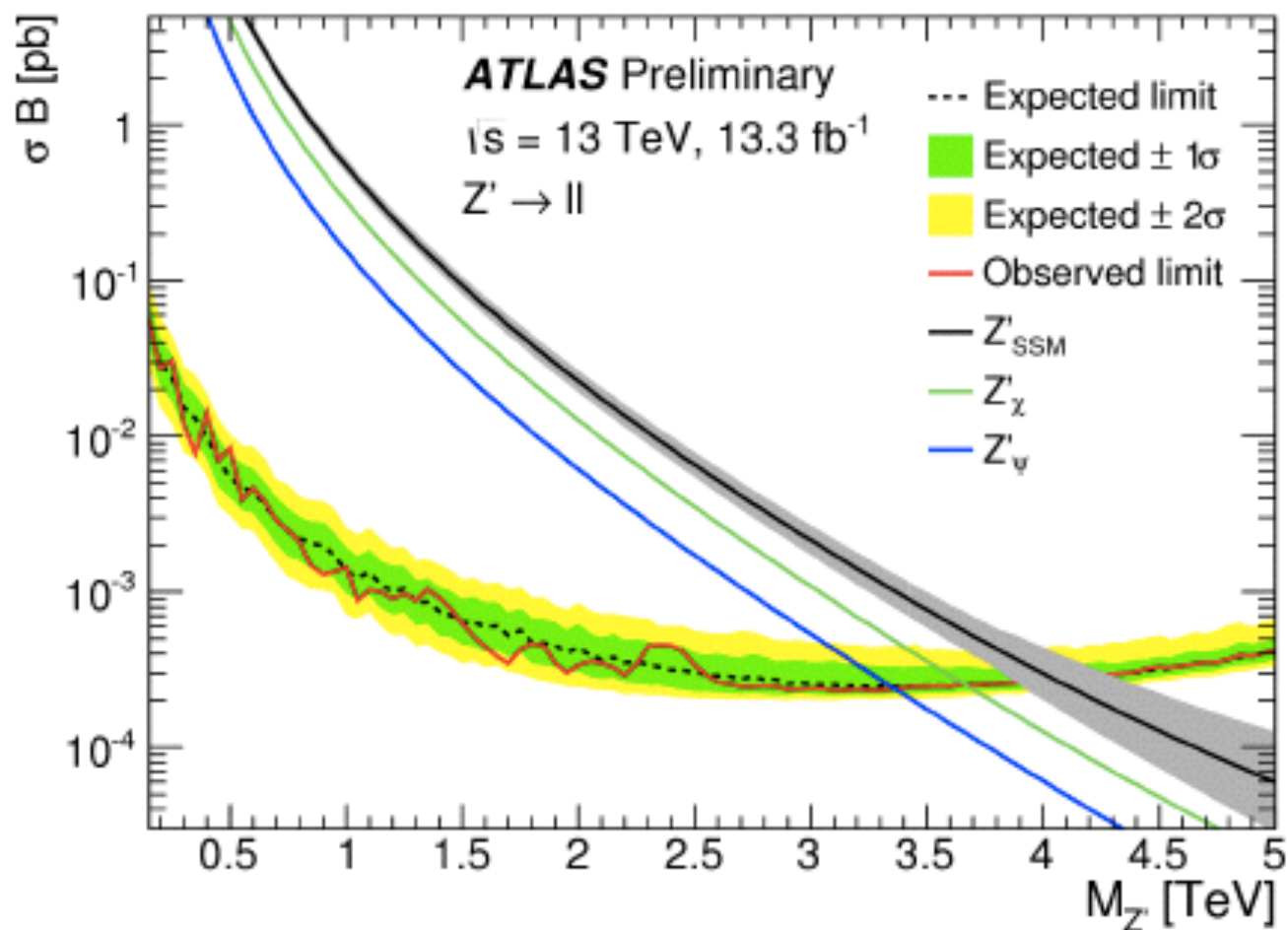
Search for Heavy Resonance: dilepton channel

Search for Heavy Resonance: dilepton channel



Search for Heavy Resonance: dilepton channel

Search for Heavy Resonance: dilepton channel

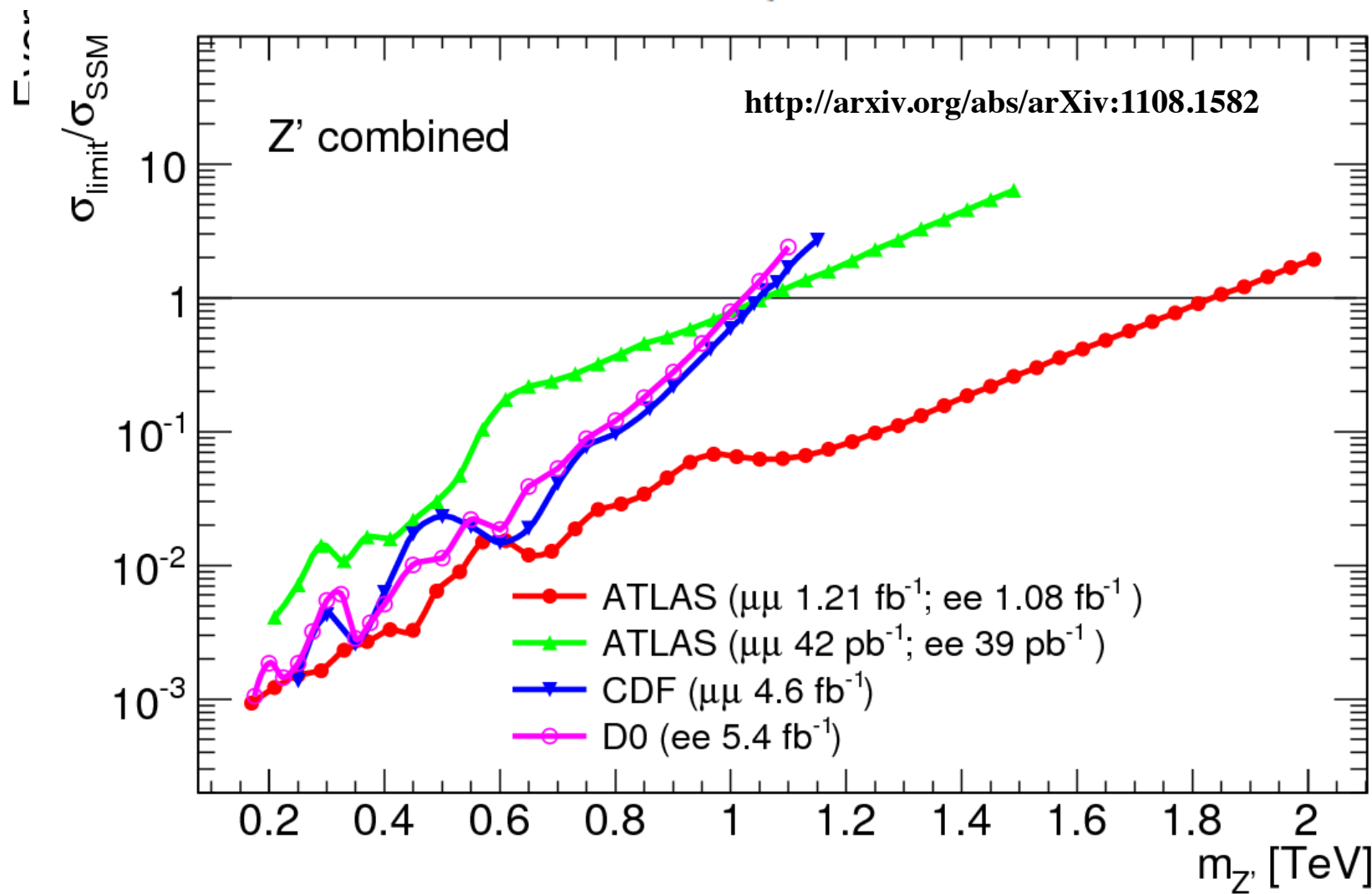


Observed lower limit (TeV) at 95% CL:
 $m(\text{SSM } Z') > 4.05 \text{ TeV}$

ATLAS status report

Now have covered a lot of phase space for many signatures

Search for dilepton resonances in pp collisions at $\sqrt{s} = 7$ TeV with the ATLAS detector

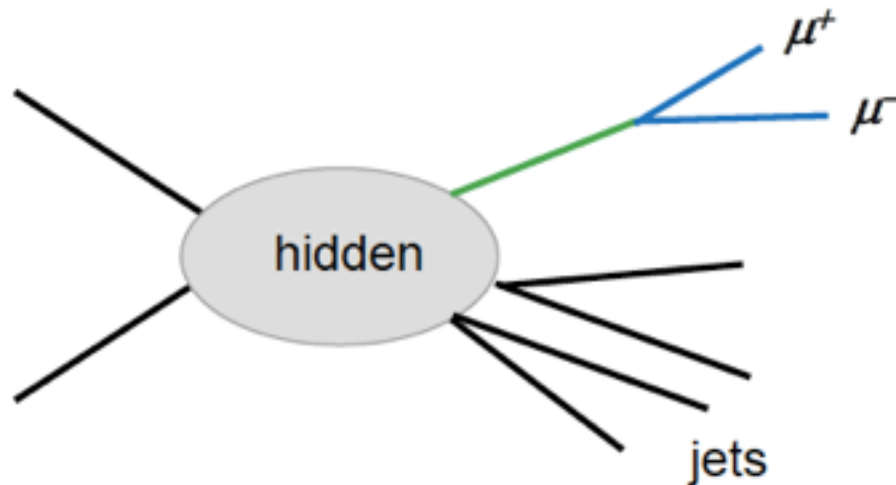


Physics Hidden in Inclusive Distributions

In various models (*common in hidden valleys/dark sectors*)

- Rare or unusual production of new neutral particle
- Bump, edge, endpoint, dip, wiggle is present
 - but swamped in inclusive background

MJS & Zurek '06
Han, Si, Zurek & MJS '07



Rare, prompt, light dilepton resonance along with hard jets

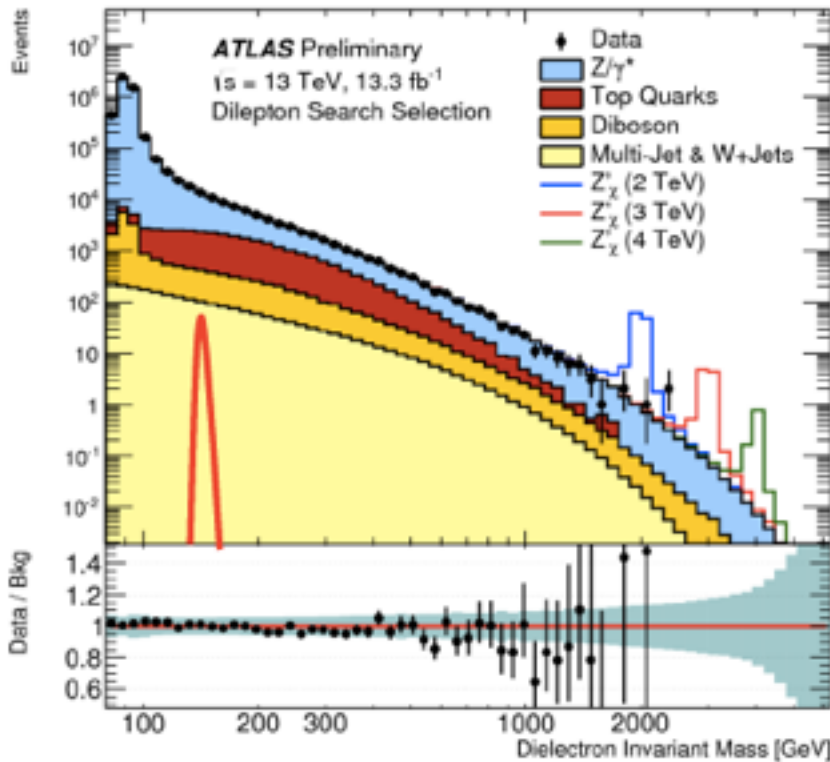
Physics Hidden in Inclusive Distributions

In various models (*common in hidden valleys/dark sectors*)

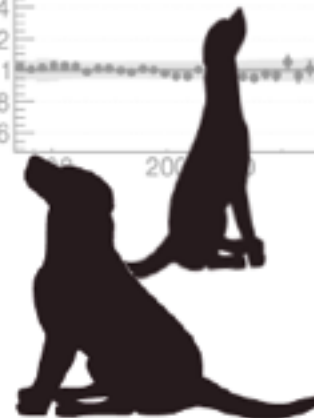
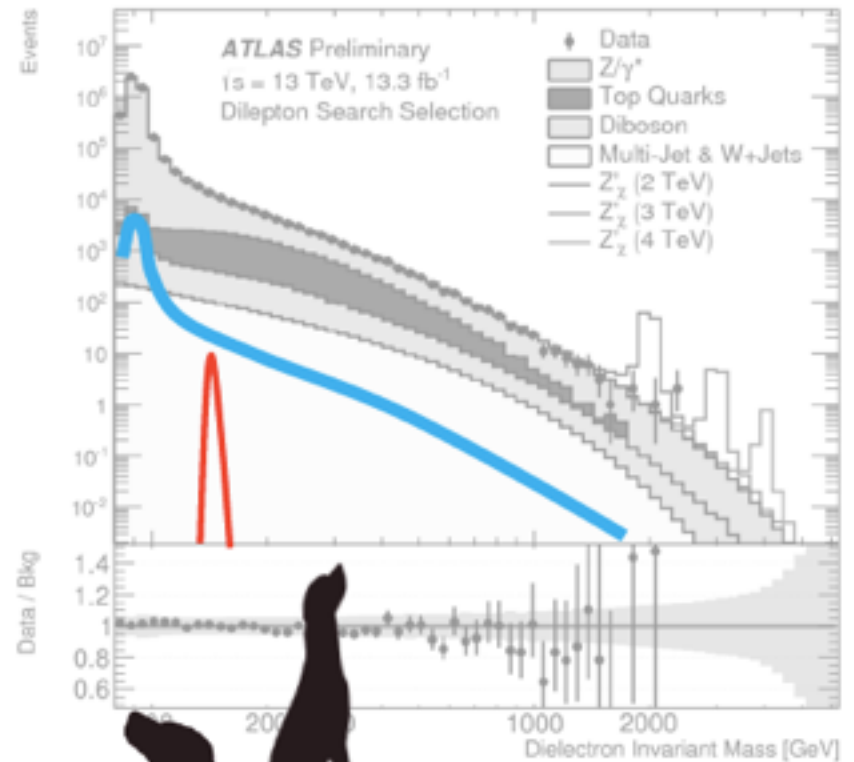
- Rare or unusual production of new neutral particle
- Bump, edge, endpoint, dip, wiggle is present
 - but swamped in inclusive background
- Point:
 - Higgs \rightarrow bb resonance is invisible in inclusive production
 - But a semi-exclusive search can reveal it
 - The selection criterion reduces background, keeps signal
- We should do this for other resonance searches as a matter of course!

Inclusive is not Conclusive

Inclusive



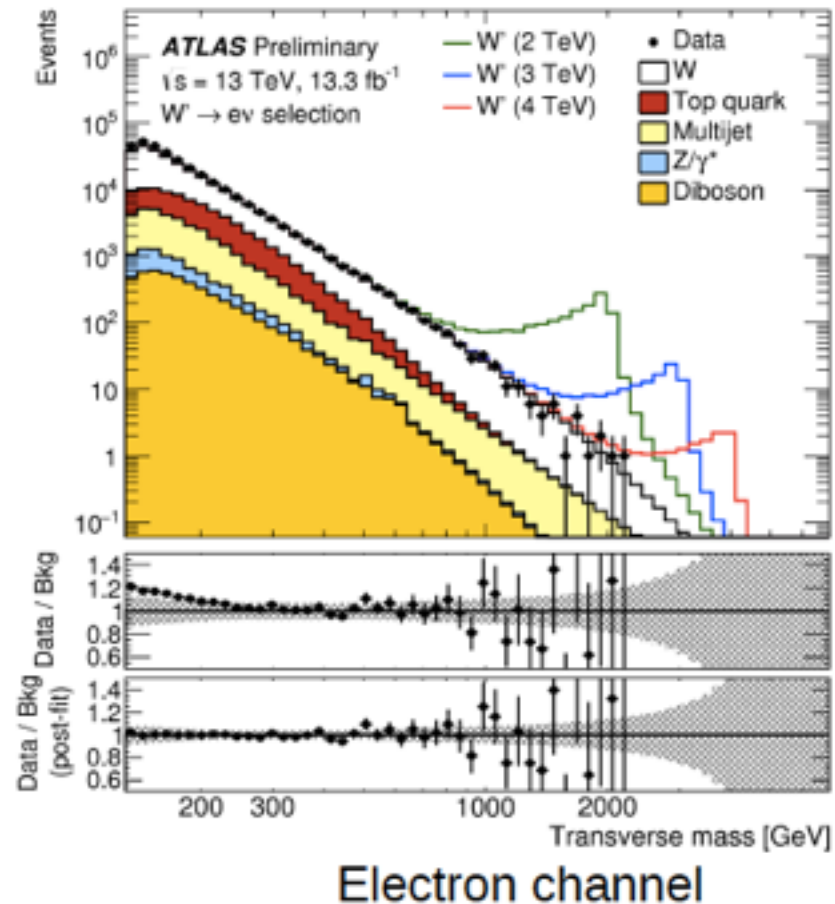
Require 4 jets and high HT and... presto!

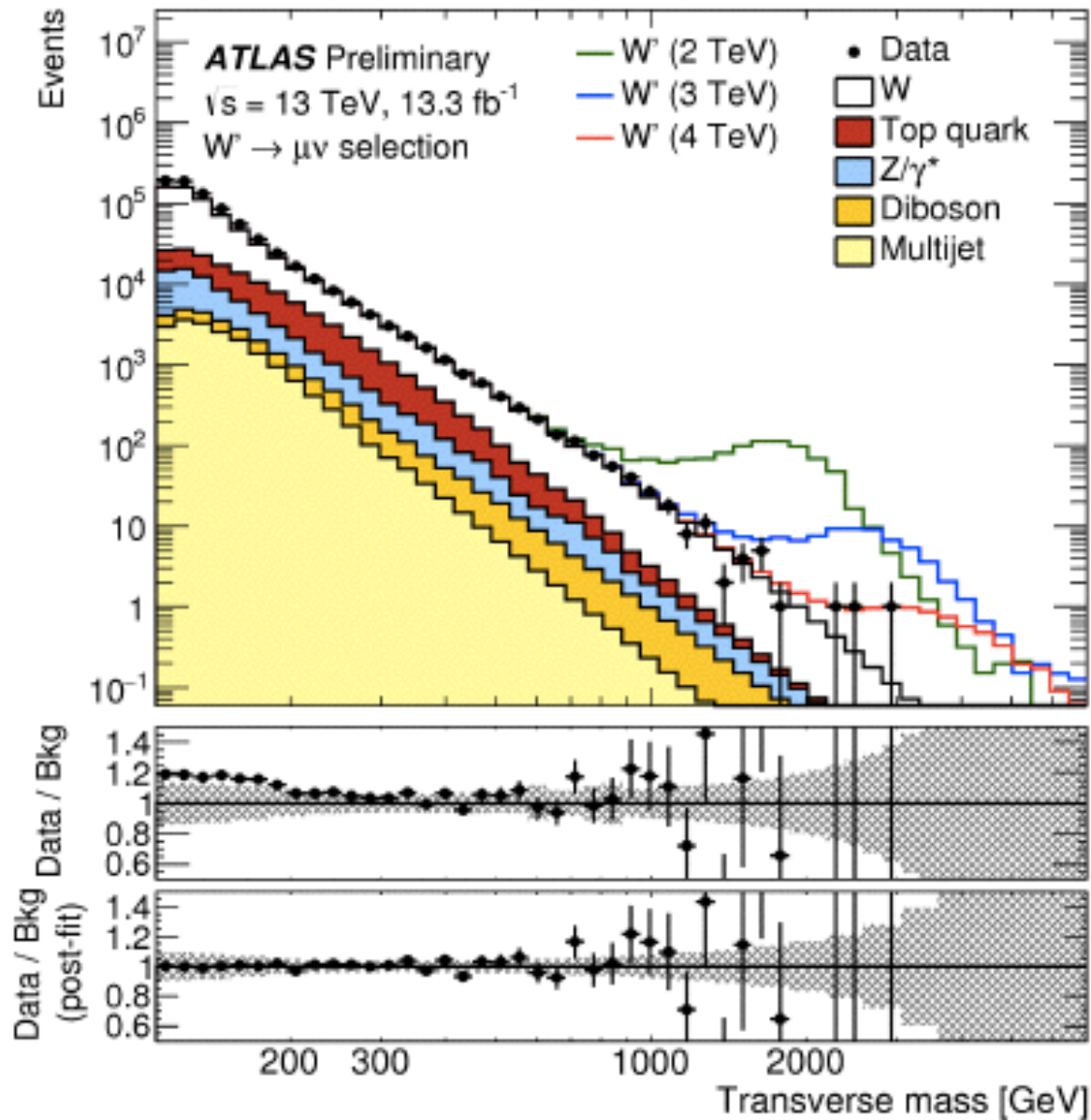


Search for Heavy Resonance: $W' \rightarrow l\nu$

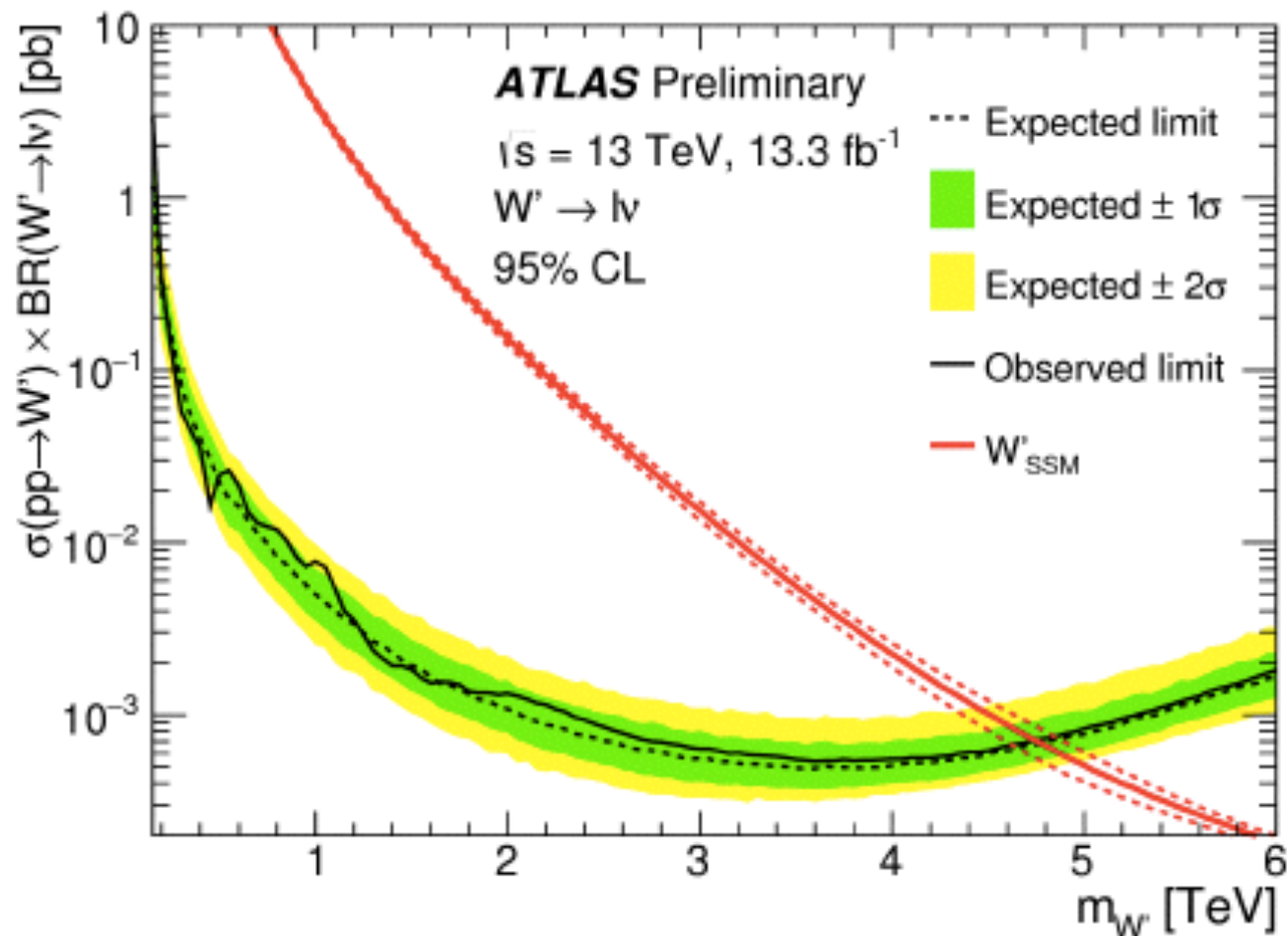
- W' : the charged equivalent of the Z'
- Bulk-RS: excited KK W
- Final state: 1 lepton + Missing E_T
- Look for Jacobian peak in transverse mass:

$$m_T = \sqrt{2p_T \cancel{E}_T (1 - \cos\Delta\phi_{l, \cancel{E}_T})}$$





Muon channel

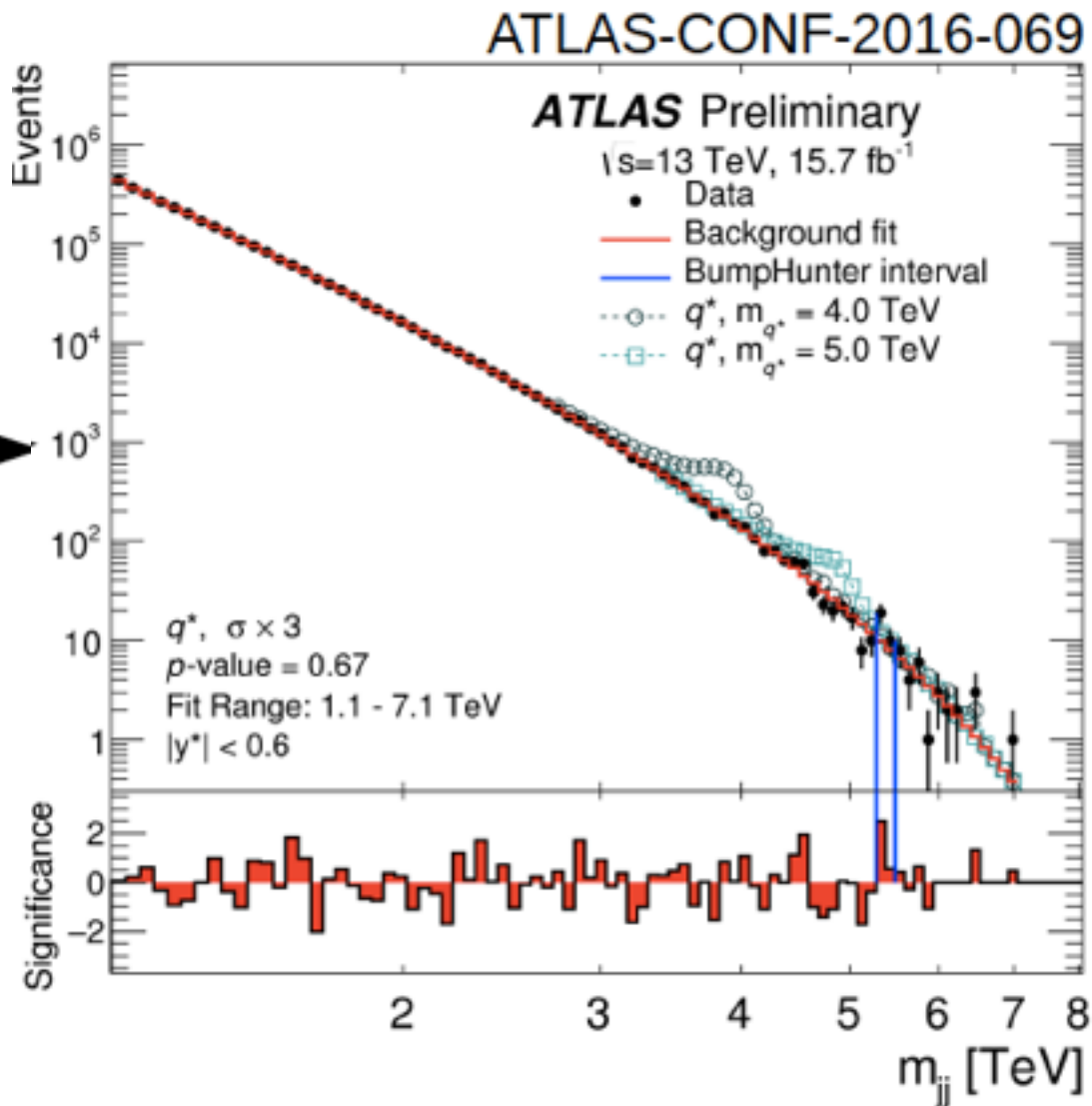


Sequential SM:

$m(W') > 4.74 \text{ TeV}$ at 95% C.L

Search for high-mass resonances decaying to jets

- W'/Z', excited quarks, strong gravity, DM-mediator
- Look for resonance above phenomenological fit of the data
- High-mass analysis [ATLAS-CONF-2016-069] limited by trigger
- To reach sensitivity at low-mass:
 - ISR photon or jet [ATLAS-CONF-2016-070]
 - Trigger-Level Analysis [ATLAS-CONF-2016-030]
- To enhance heavy-flavor:
 - B-tagging [ATLAS-CONF-2016-060]



Search for Heavy Resonance: Dijet High-mass

Trigger:

- 1-jet trigger $E_T \sim 380$ GeV (100% efficient)

Event selection:

- anti- k_T $R=0.4$ jets
- Leading jet $p_T > 440$ GeV, sub-leading $p_T > 60$ GeV
- $|y^*| < 0.6$ (or 1.2 depending on model)

Selection implies:

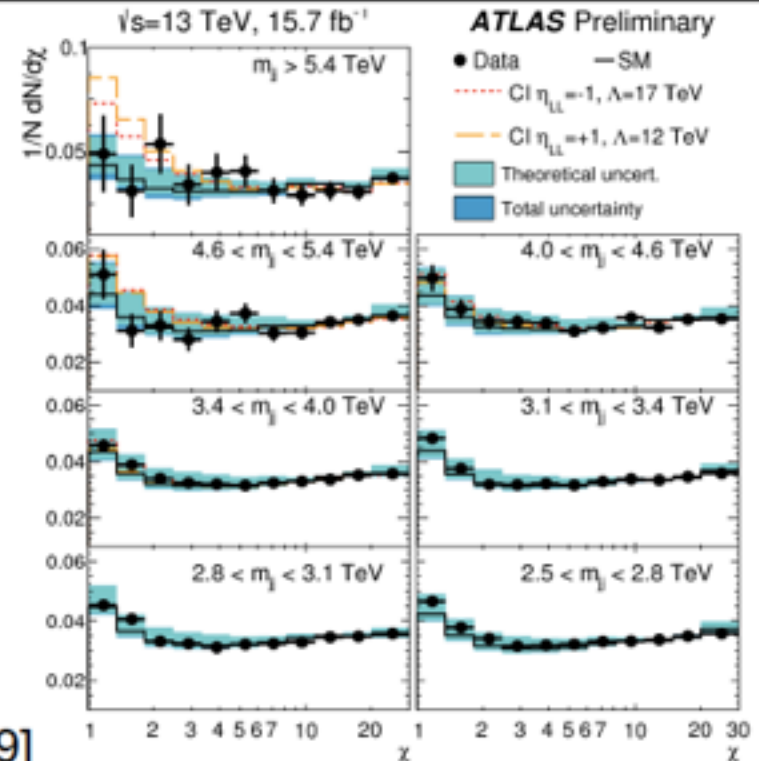
- $m(jj) \gtrsim 1.1$ TeV

Angular analysis:

- $|y^*| < 1.7$ and $|y_B| < 1.1$
- $m(jj) > 2.5$ TeV

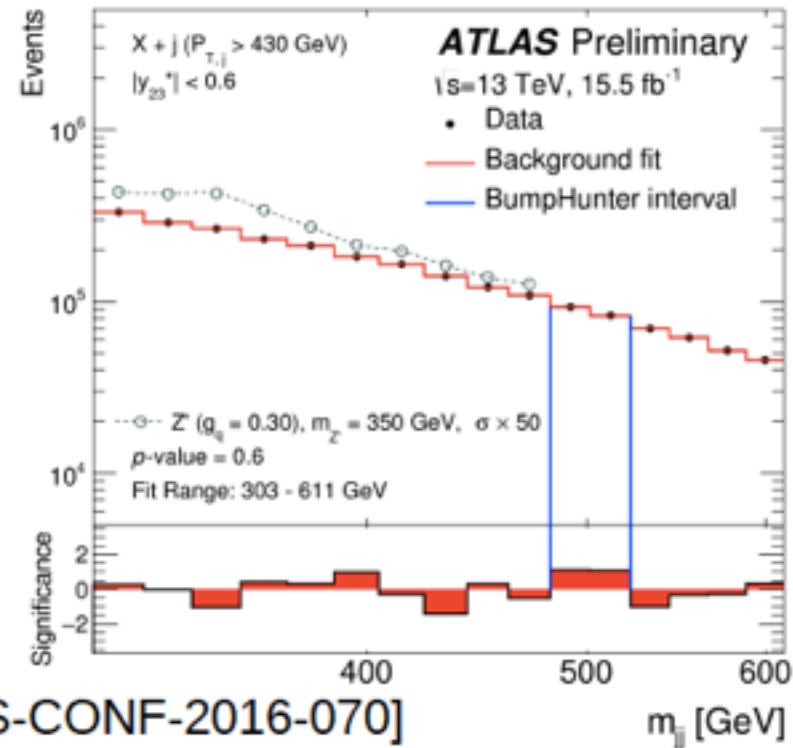
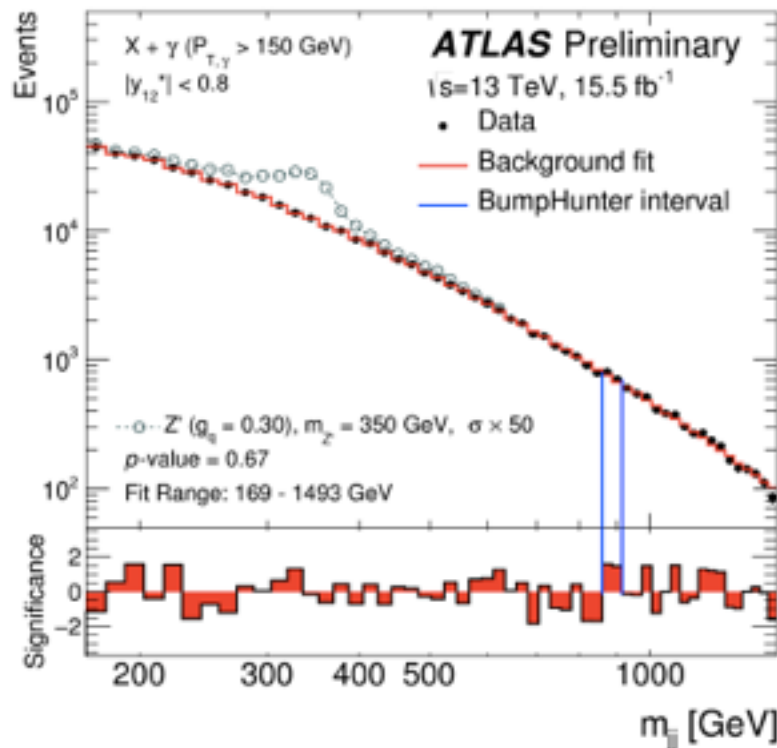
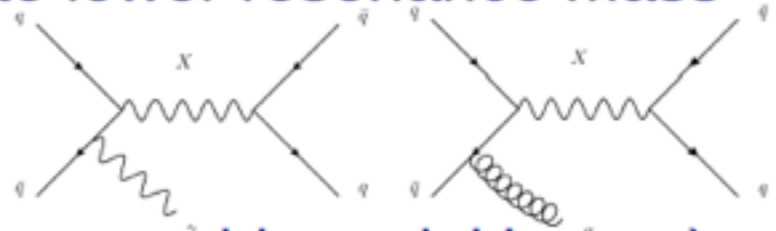
[ATLAS-CONF-2016-069]

Model	95% CL exclusion limit	
	Observed	Expected
Quantum black holes, ADD (BLACKMAX generator)	8.7 TeV	8.7 TeV
Excited quark	5.6 TeV	5.5 TeV
W'	2.9 TeV	3.3 TeV
W^*	3.3 TeV	3.3 TeV
Contact interactions ($\eta_{LL} = +1$)	12.6 TeV	13.7 TeV
Contact interactions ($\eta_{LL} = -1$)	19.9 TeV	23.7 TeV



Search for Low-mass Resonance: Dijet + ISR

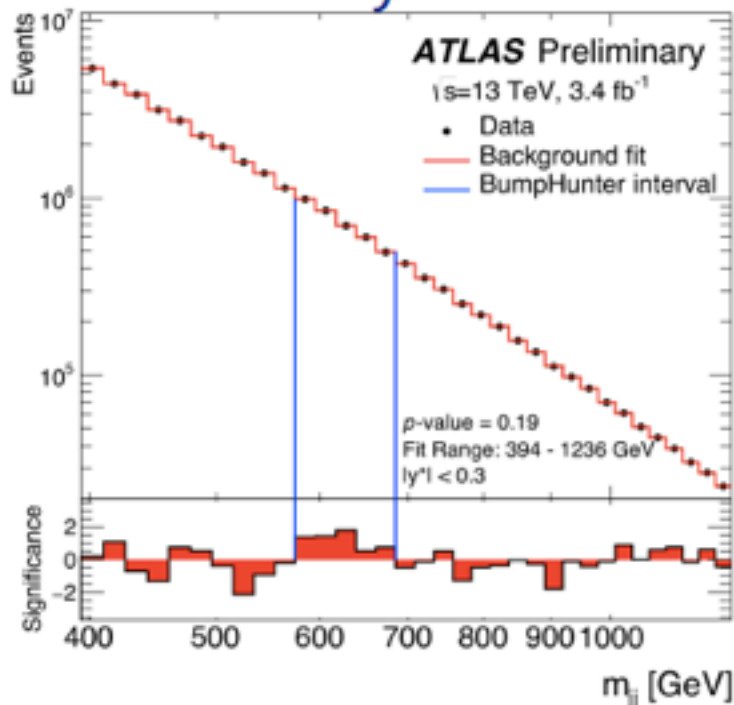
- Same trigger and offline thresholds (380 → 440 for jets, 140 → 150 for photon), but reach to lower resonance mass by requiring ISR photon or jet
- Photon: 200-1500 GeV
- Jet: 300-600 GeV (fails above due to combinatorial issues)



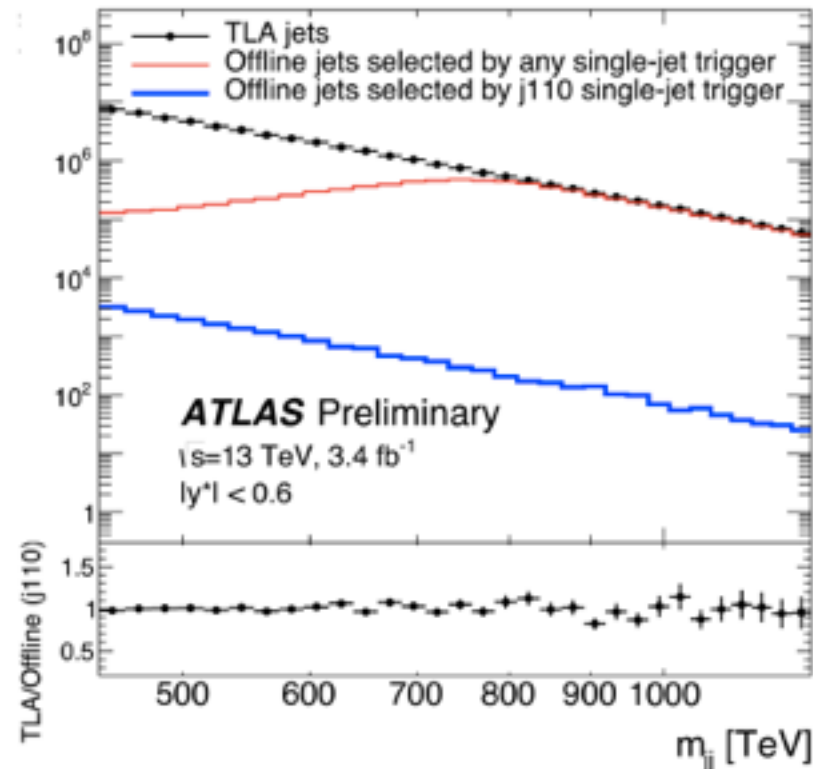
[ATLAS-CONF-2016-070]

Search for Low-mass Resonance: Dijet Trigger-Level Analysis

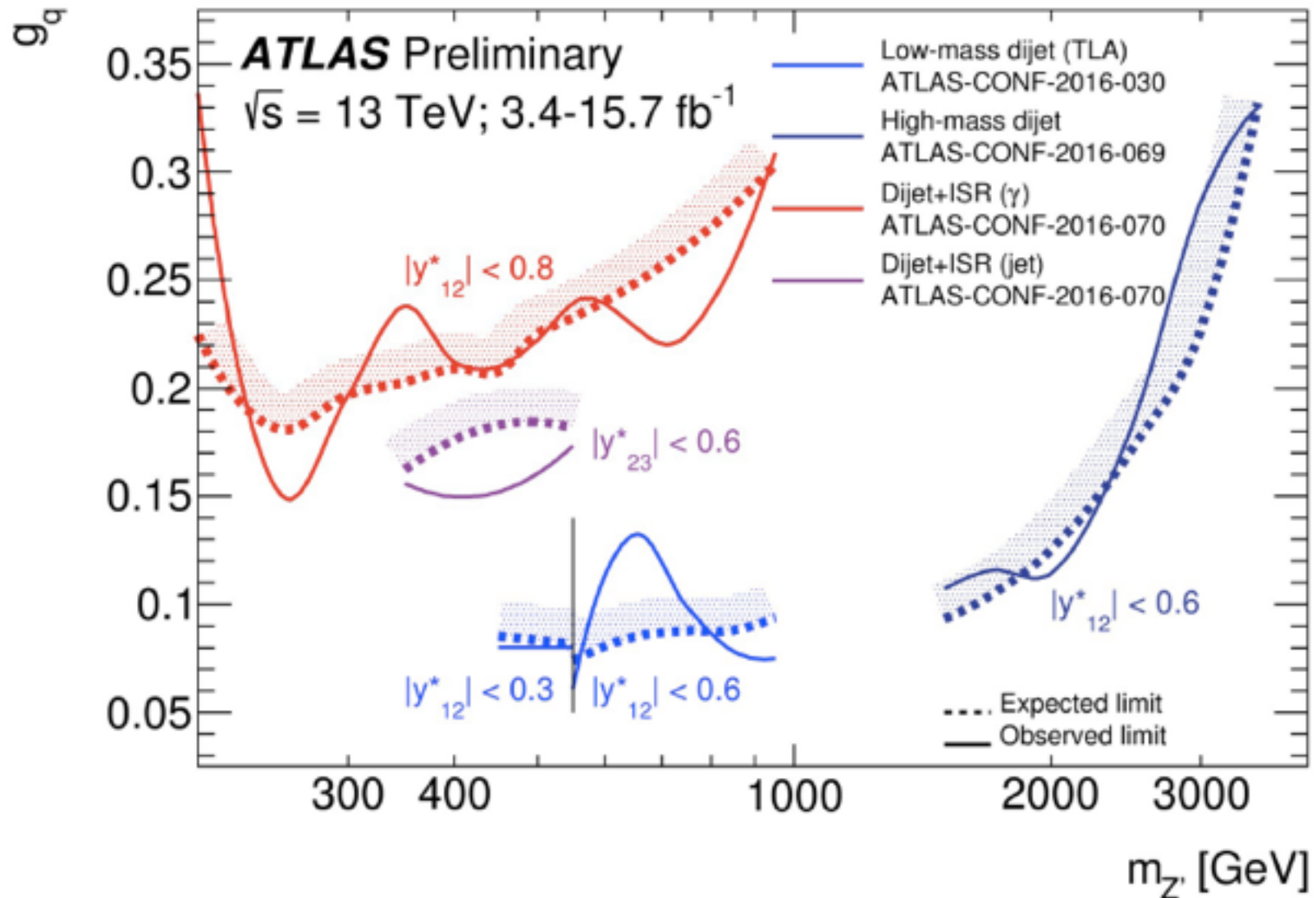
- Lower trigger-threshold by keeping only partial information of the event
- Addition jet calibration and cleaning applied online
- 2015 data only



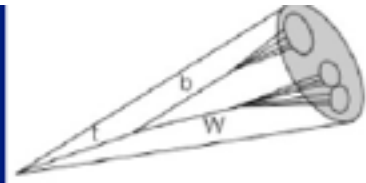
[ATLAS-CONF-2016-030]



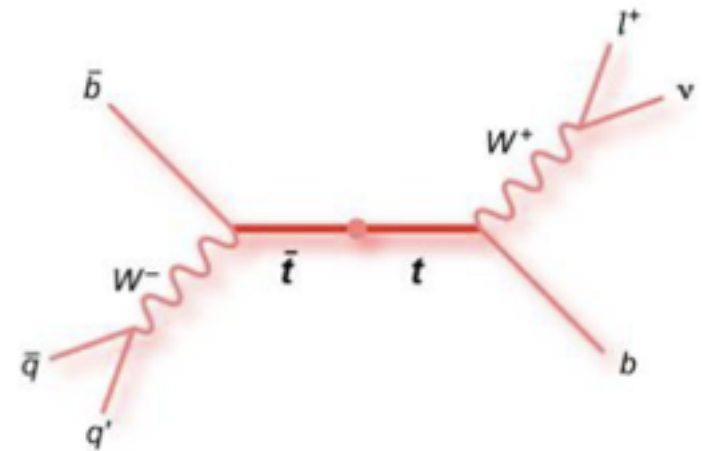
Dijet searches: a summary



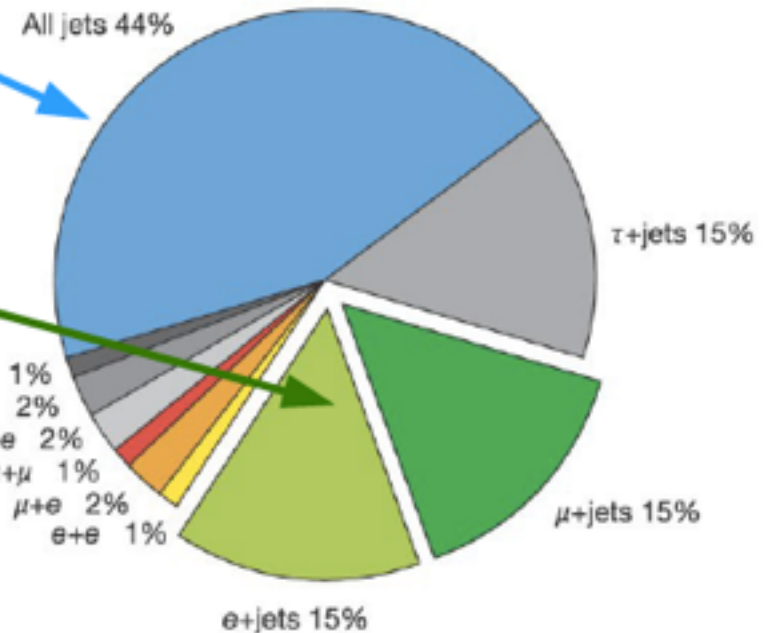
Top-antitop Resonance



- Event Topology:
 $t\text{-}\bar{t} \rightarrow Wb Wb$
- Final state depends on W decays:

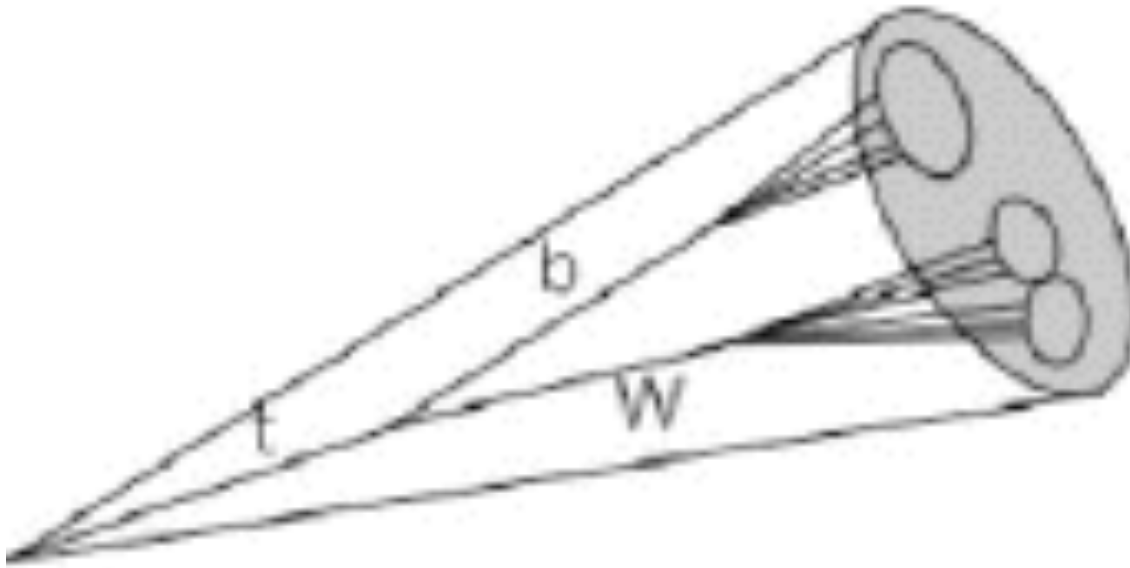


- All-hadronic final state:
 $2 W \rightarrow jj$
 $2 b\text{-jets} + 4 \text{ light jets (+0}\nu)$
- Lepton+Jets final state:
 $1 W \rightarrow l\nu, 1 W \rightarrow jj$
 $1l + 2 b\text{-jets} + 2 \text{ light jets (+1}\nu)$
- Dilepton final state:
 $\text{Both } W \rightarrow l\nu (l = e \text{ or } \mu)$
 $2l + 2 b\text{-jets (+2 neutrinos)}$

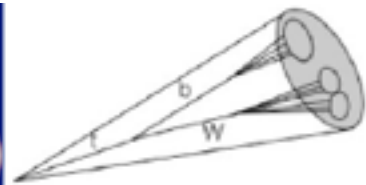


Larger Branching Ratio
but more background

Top-antitop Resonance



Top-antitop Resonance Lepton+Jets Channel (ATLAS)

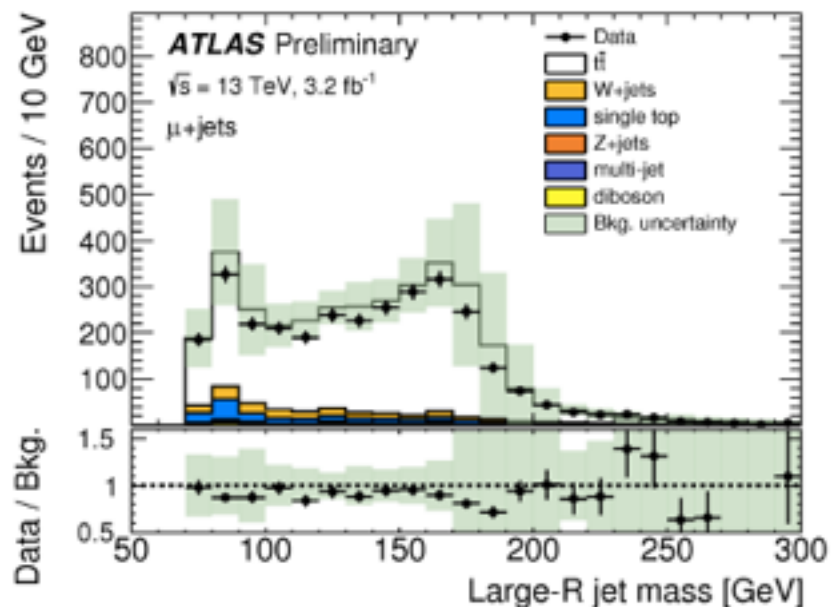
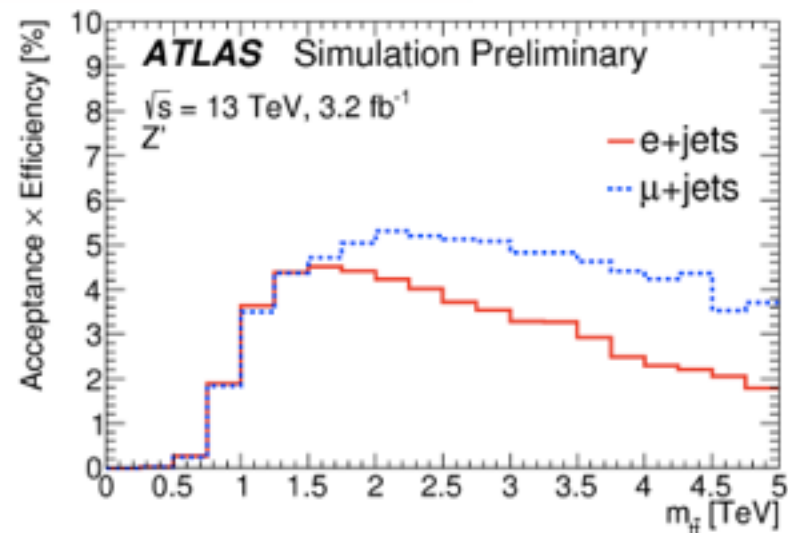


- First 13 TeV result, boosted channel only:
 - anti-kT R=1.0, $p_T > 300$ GeV, $|\eta| < 2$
 - Top-tagging (80% efficiency working point) using jet mass and n-subjettiness ratio τ_{32}

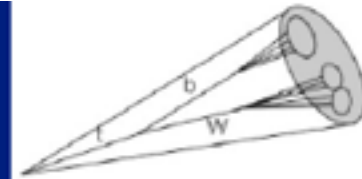
- Improve efficiency at high t-tbar mass with:

- Lepton “mini-isolation”:
 - smaller isolation cone at high momentum
 - tolerate muons inside jets at high pT
- R=0.2 track-jet b-tagging (at least one b-tag)

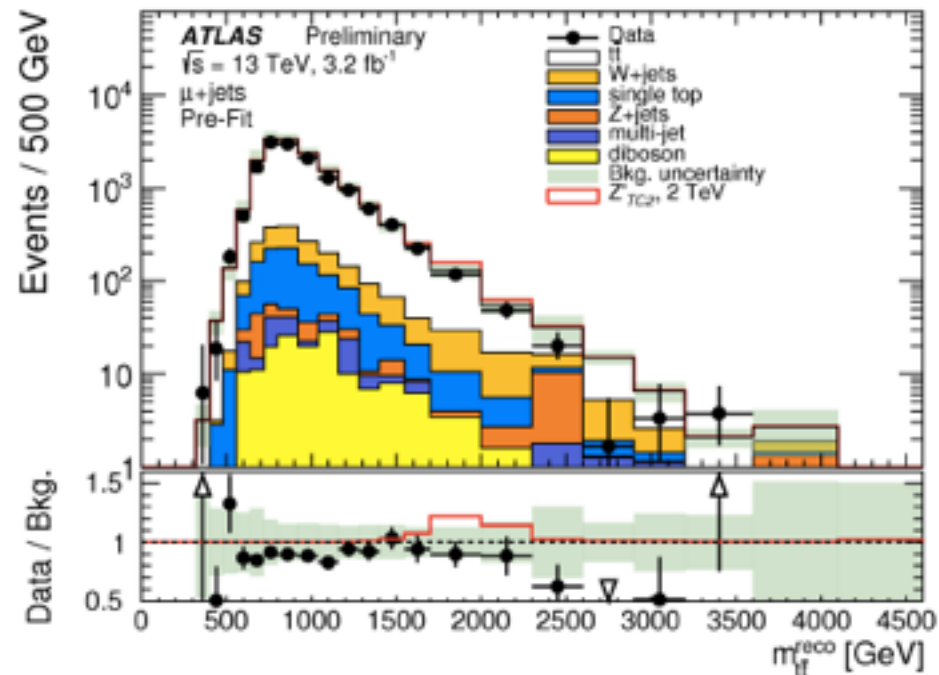
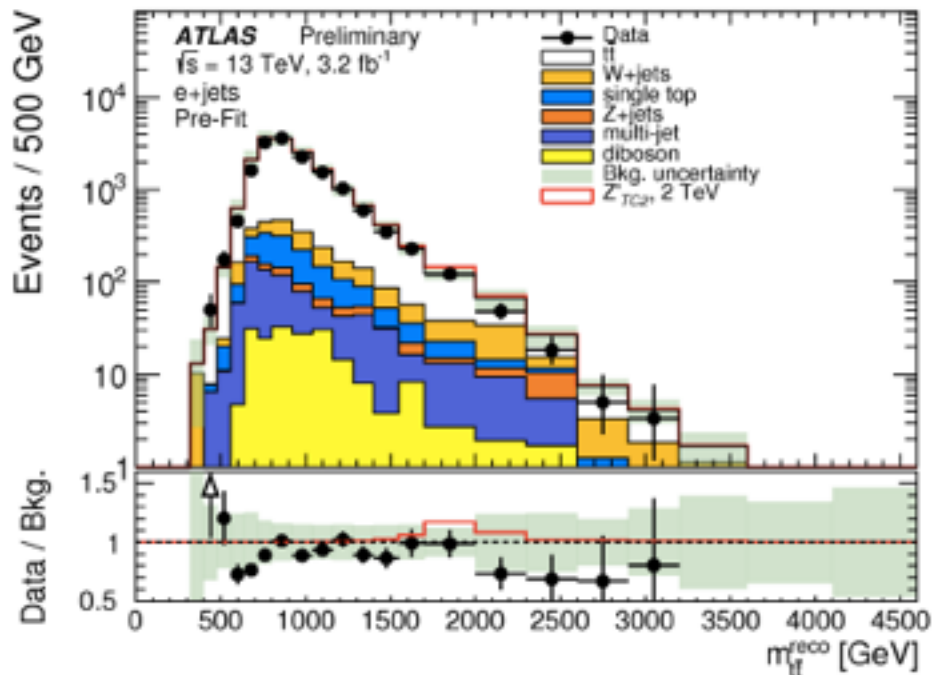
[ATLAS-CONF-2016-014]



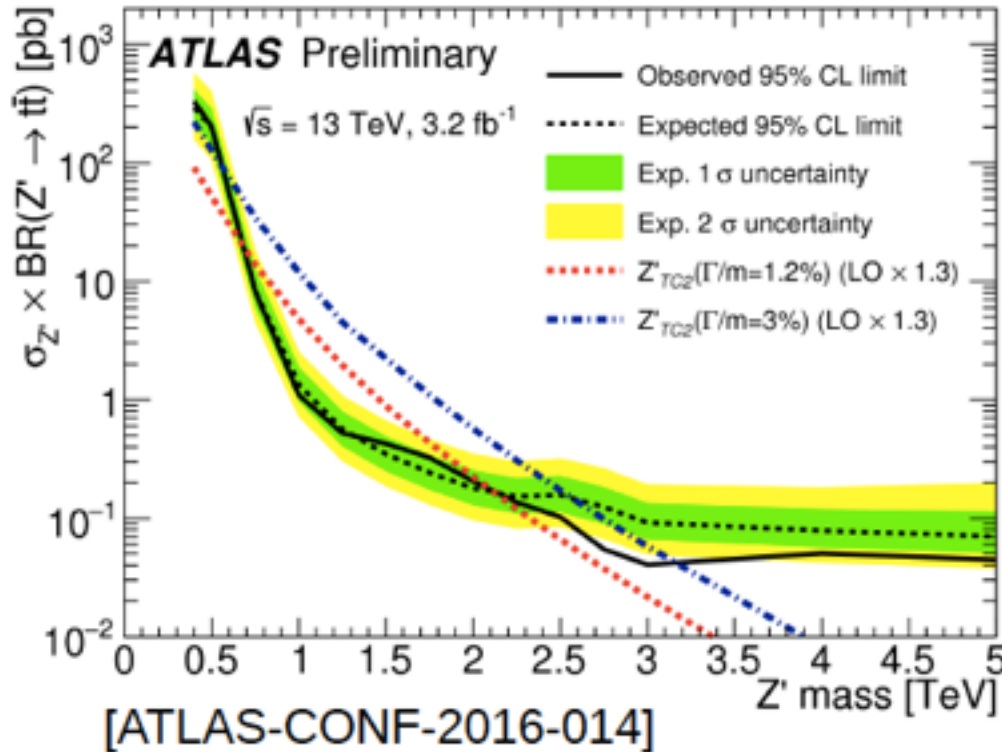
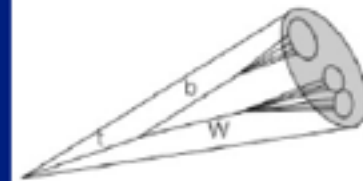
Top-antitop Resonance L+Jets Channel



- Top-antitop mass spectrum at 13 TeV:



Top-antitop Resonance L+Jets Channel



- Limit with 3 /fb at 13 TeV:

$m(Z' \text{ 1.2\% width}) > 2.0 \text{ TeV}$

- Already surpassing recent ATLAS Run 1 paper: <http://arxiv.org/abs/1505.07018>

$m(Z' \text{ 1.2\% width}) > 1.8 \text{ TeV}$

Pre-fit impact on μ :

□ $\theta_0 = +\Delta\theta$ □ $\theta_0 = -\Delta\theta$

Post-fit impact on μ :

■ $\theta_0 = +\Delta\theta$ ■ $\theta_0 = -\Delta\theta$

● Nuis. Param. Pull

Large-R jet cross calibration

Large-R jet run 1 extrapolation

Luminosity

b-tagging light EV 0 high pt

Multi-jet (e)

tt prod. cross section

b-tagging c EV 0 high pt

b-tagging light EV 0 low pt

b-tagging light EV 0 med. pt

Muon trigger eff. stat. unc.

b-tagging c EV 0 low pt

b-tagging eff. EV 0

Muon identification eff. syst. unc.

b-tagging c EV 0 med. pt

PDF EV 5

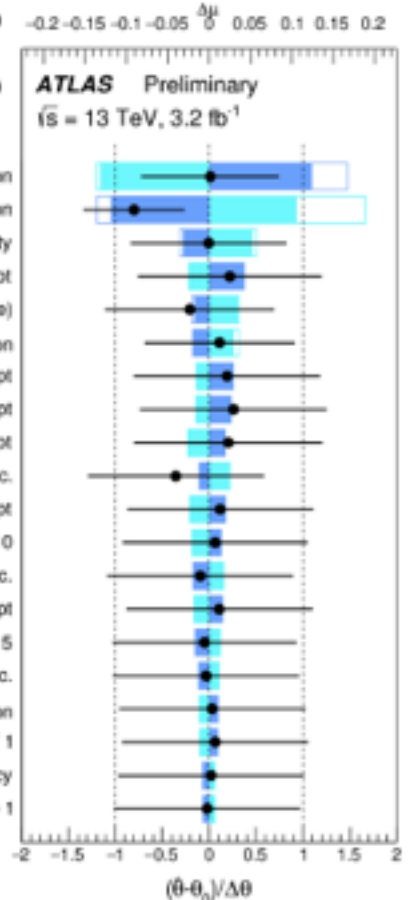
Muon isolation eff. syst. unc.

tt electroweak correction

b-tagging c EV 1

Electron isolation efficiency

Small-R Jet Energy Scale 1



Heavy Resonances: a summary

- Heavy resonance searches benefit the most (and the fastest) from increase in center-of-mass energy
- Corollary: future improvements will take a lot more time!

95% CL limit (TeV)	CDF	Run 1 '12	Moriond '15	ICHEP '16	300 fb ⁻¹ 14 TeV	3000 fb ⁻¹ 14 TeV
Z' → ll	1.1	2.9	3.4	4.1	6.5	7.8
q* → dijet	0.9	4.1	5.2	5.6	7.4	8
Z' → tt	0.9	1.8		2.0	3.3	5.5

ATLAS upgrade: ATL-PHYS-PUB-2013-003, ATL-PHYS-PUB-2015-004

CDF:

<http://arxiv.org/abs/1101.4578> (4.6 fb⁻¹)

<http://journals.aps.org/prd/abstract/10.1103/PhysRevD.83.031102> (5.3 fb⁻¹)

<http://prd.aps.org/abstract/PRD/v79/i11/e112002> (1.1 fb⁻¹)

<http://arxiv.org/abs/1211.5363> (9.5 fb⁻¹)

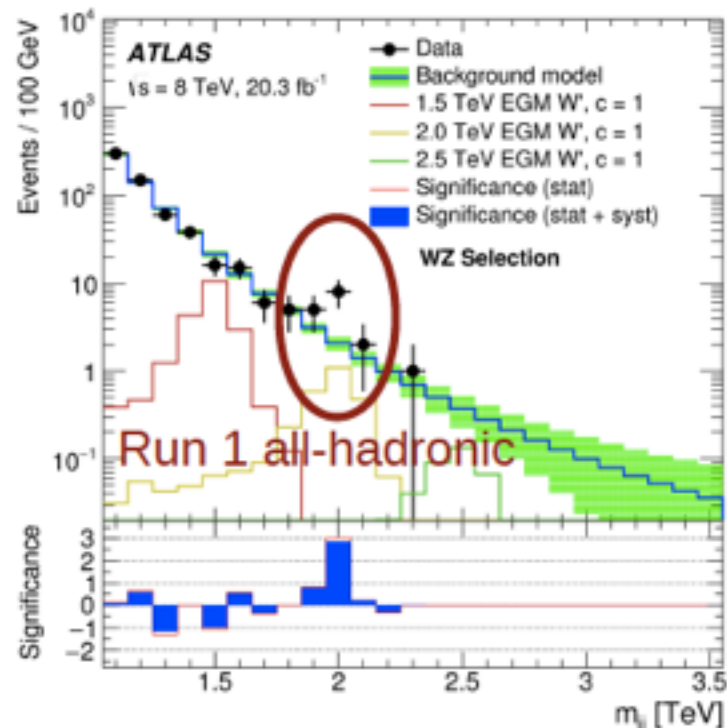
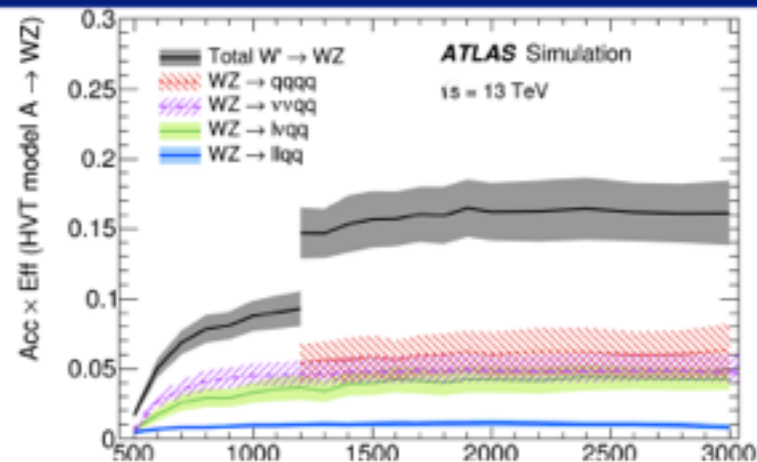
Search for WW/WZ/ZZ resonances

- Many channels:

- $WW \rightarrow l\nu l\nu$ or $l\nu qq$ or $(qq)(qq)$
- $WZ \rightarrow l\nu ll$ or $l\nu qq$ or $llqq$ or $(qq)(qq)$
- $ZZ \rightarrow (ll)(ll)$ or $(ll)(\nu\nu)$ or $llqq$ or $(qq)(qq)$
- (qq) can be either resolved into two jets, or merged into one fat jet

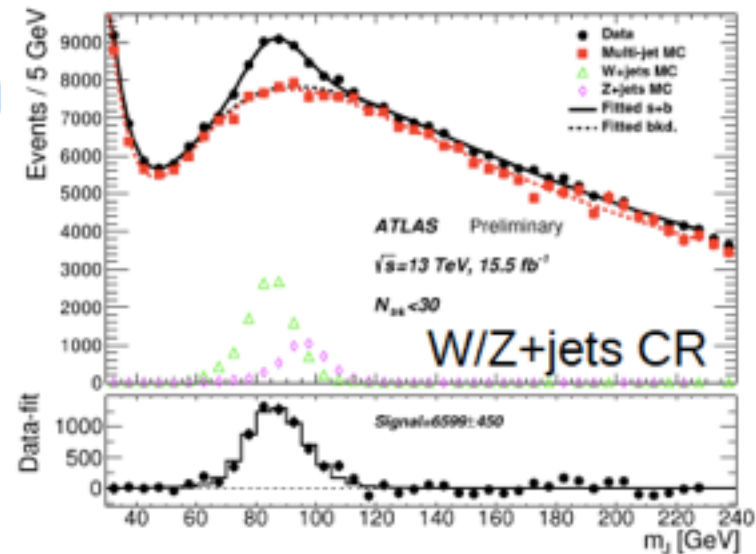
- Interpretation in terms of HVT W'/Z' , bulk RS graviton, and 2HDM

- Got some attention between Run 1 and Run 2 due to the all-hadronic ATLAS result (although not seen in other channels)

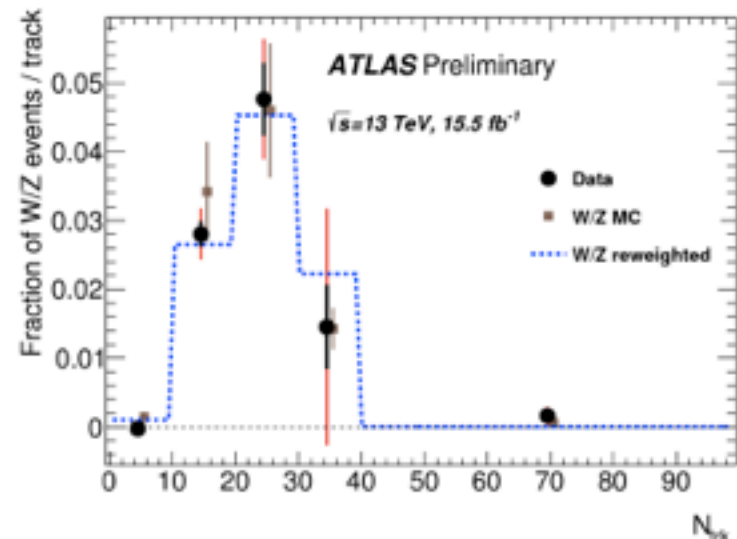


Diboson all-hadronic channel

- Look for pair of vector bosons, each decaying to pair of quarks:
 - WW or WZ or $ZZ \rightarrow (qq)(qq)$
- High-momentum → boosted topology:
 - anti- k_T $R=1.0$, $|\eta| < 2$
 - $p_T > 200$ GeV, $m > 30$ GeV after trimming
 - Boson-tagging:
 - mass window of 15 GeV around W or Z mass
 - p_T -dependent cut on $D2(\beta=1)$
 - $N(\text{tracks}) < 30$



[ATLAS-CONF-2016-055]



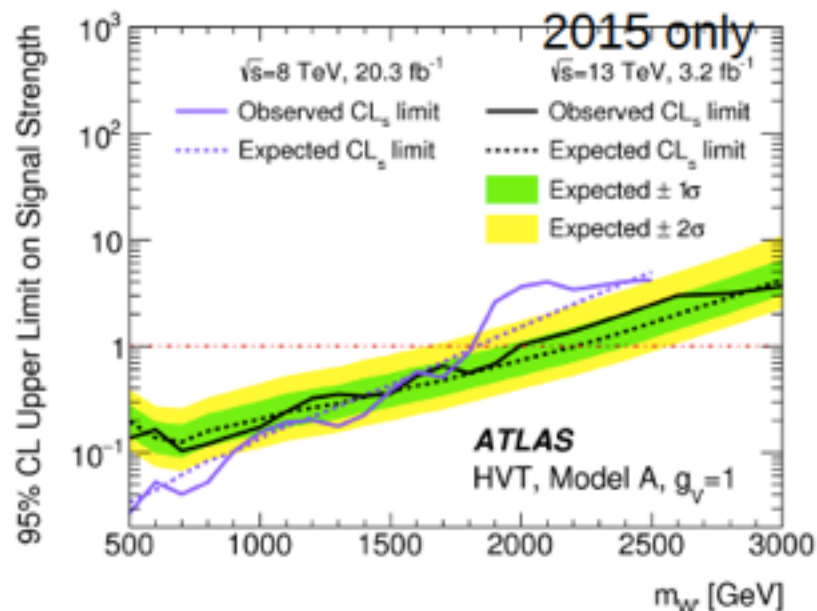
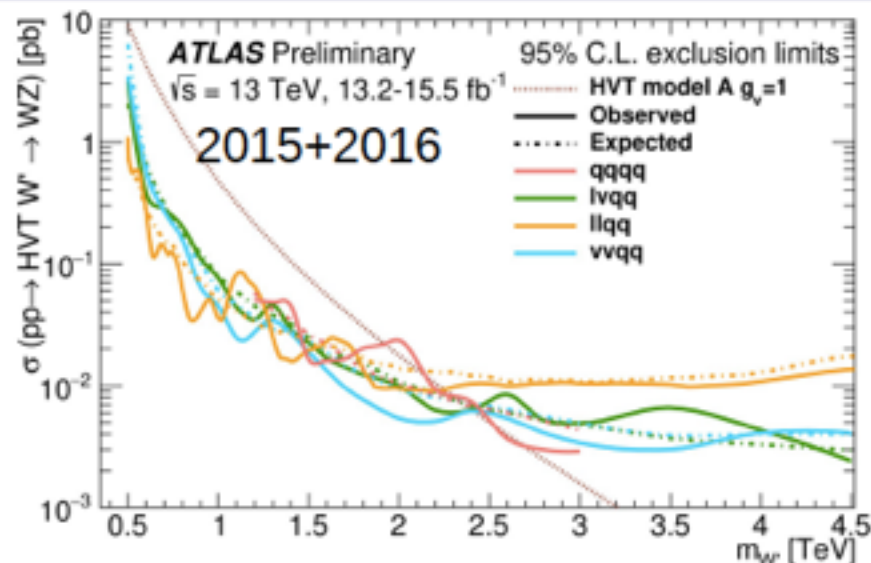
Diboson comparison of channels

- Full combination only with 2015 data currently
- Combine channels with lepton-multiplicity from 0 to 4:

(qq)(qq) (lv)(qq)
 (ll)(qq) (vv)(ll)

	vvqq	lvqq	llqq	qqqq (JJ)
WZ	X	X	X	X
WW		X		X
ZZ	X		X	X

- Run 1 excess is excluded by Run 2 (even with 2015 only)



SUSY is a class of models

SUSY breaking

Gravity, Gauge, Anomaly,
Dilaton/Moduli, Mirage,
Gaugino, D-term, Z-prime, ...

R-Parity conservation



SUSY models possibly with
extra matter/gauge bosons

Various forms of
SUSY spectra

LSP?

NMSSM, USSM, μ vSSM,
E6SSM, PQNMSSM,...

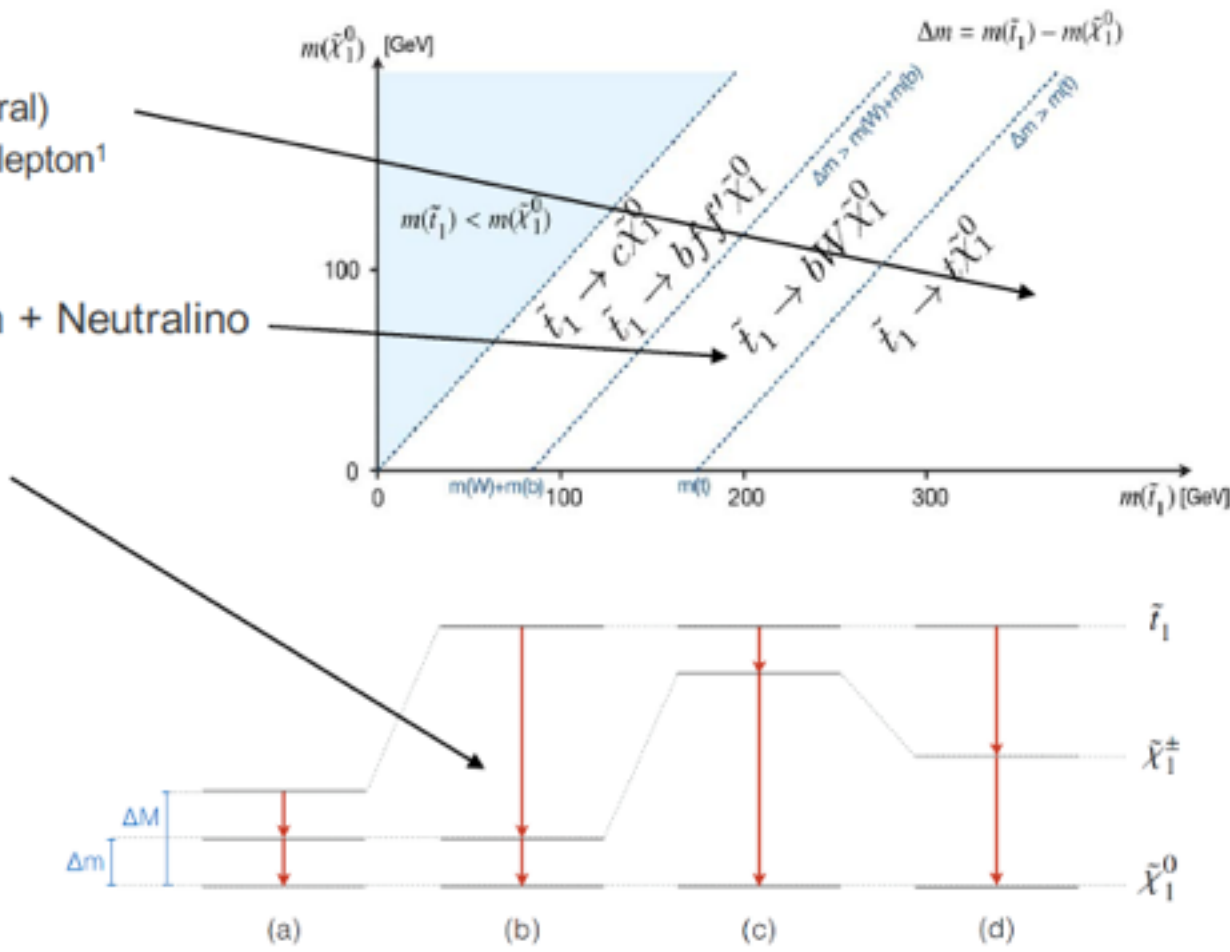
Natural, Split, Compressed, Stealth, ...

Neutralino, SM, Gravitino,
Axino, ...

- > Make blocks by “Naturalness”
 - 3rd generation
 - “Strong” gluino, squark
 - EW

Strategy and outline

- > stop -> top + Neutralino
 - Analysis approach (general)
 - Show a search with one lepton¹ in the final state
- > stop -> bottom + W-boson + Neutralino
- > stop -> bottom + Chargino
- > stop2 production
- > stop GMSB / stau
- > Stop RPV
- > sbottom production



Using this analysis also to explain the general analysis procedure

stop \rightarrow top + neutralino (stop 1L)

ATLAS-CONF-2016-050 (13.3 fb⁻¹)

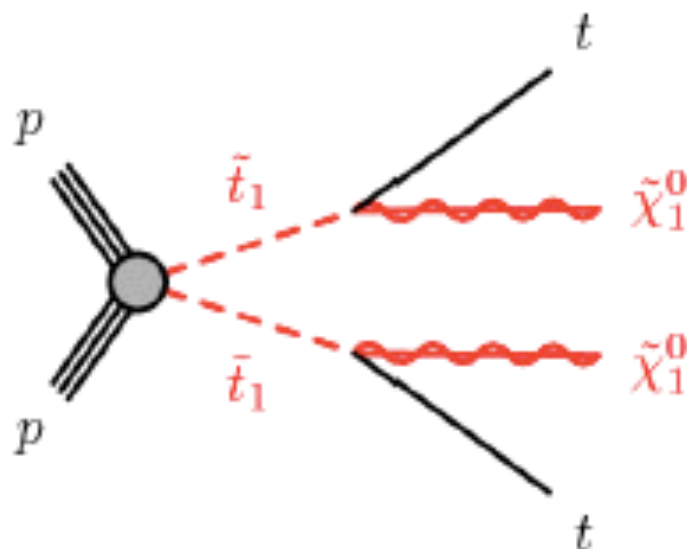
arxiv:1606.03903 with 3.2 fb⁻¹

Using this analysis also to explain the general analysis procedure

stop \rightarrow top + neutralino (stop 1L)

ATLAS-CONF-2016-050 (13.3 fb $^{-1}$)

arxiv:1606.03903 with 3.2 fb $^{-1}$



Using this analysis also to explain the general analysis procedure

stop \rightarrow top + neutralino (stop 1L)

ATLAS-CONF-2016-050 (13.3 fb⁻¹)

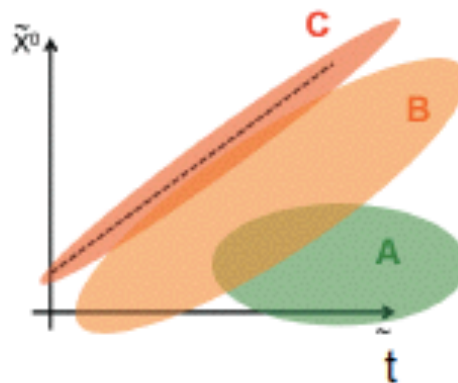
arxiv:1606.03903 with 3.2 fb⁻¹

Using this analysis also to explain the general analysis procedure

stop \rightarrow top \dagger

ATLAS-CONF-2016-0

Specific regions in the $m(\tilde{N}_1)$ - $m(\text{stop})$ plane result in different experimental signatures and topologies (different phenomenology).



e.g.:

B: Small mass splitting, top decay products are resolved.

Design specific “signal regions” (SR1 in this case) to target this region.

A: Large mass splitting, resulting in boosted top quarks (“ $t\tilde{N}_{\text{high}}$ ” signal region)

Using this analysis also to explain the general analysis procedure

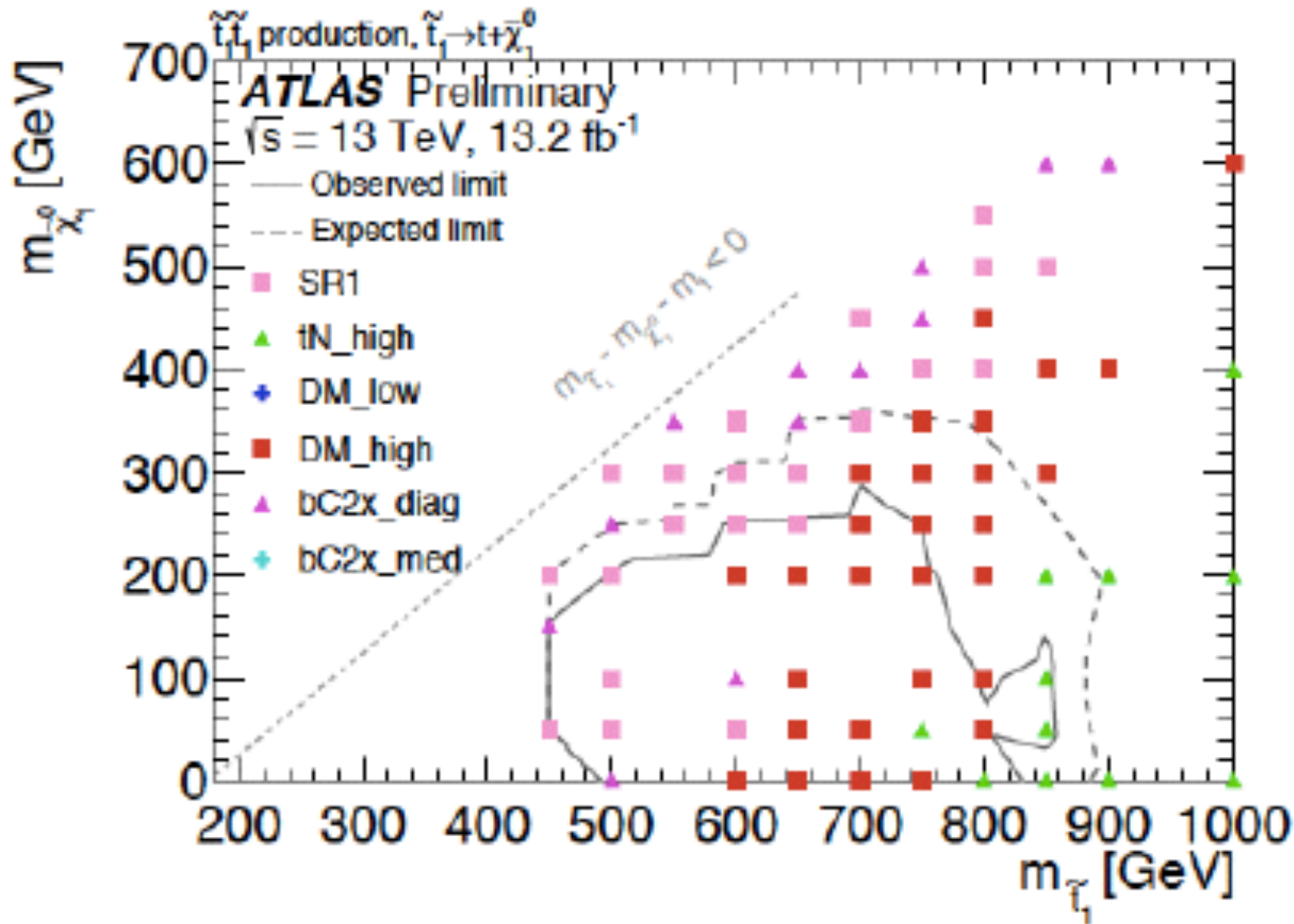
stop \rightarrow top + neutralino (stop 1L)

ATLAS-CONF-2016-050 (13.3 fb⁻¹)

arxiv:1606.03903 with 3.2 fb⁻¹

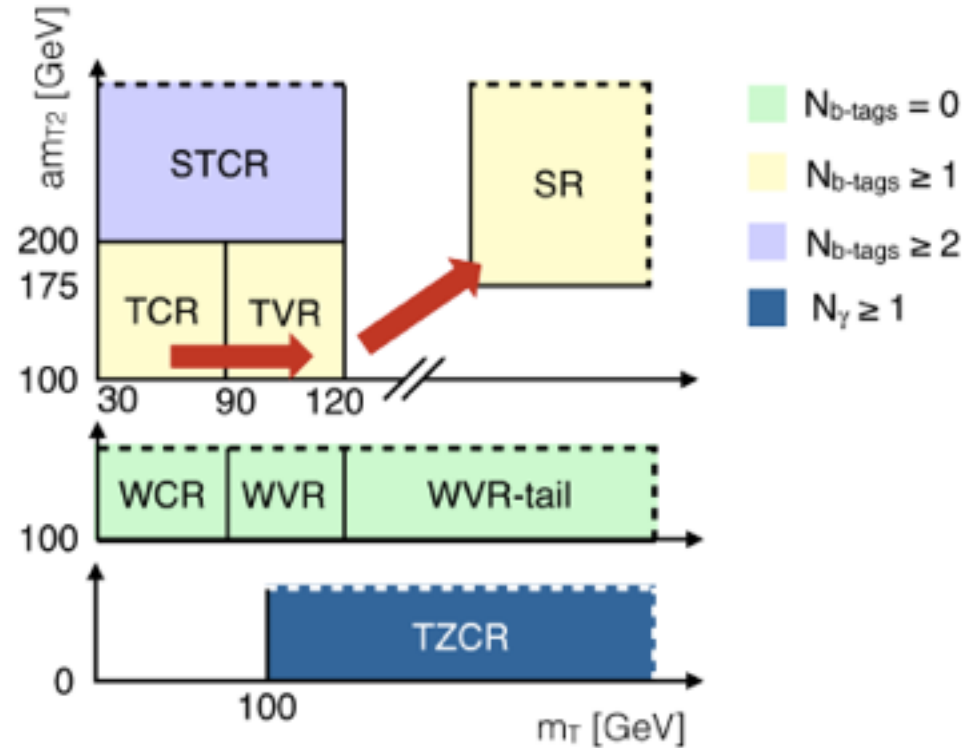
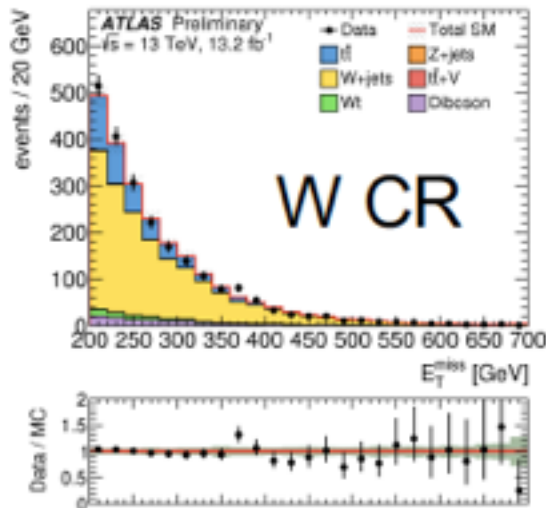
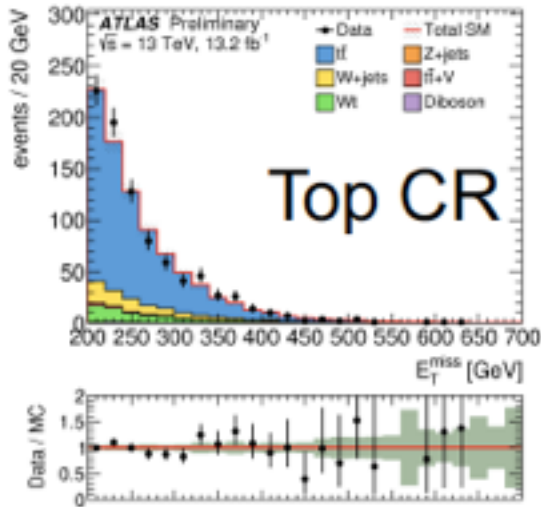
Using this analysis also to explain the general analysis procedure

st
ATL



Control regions (normalization)

Data driven background

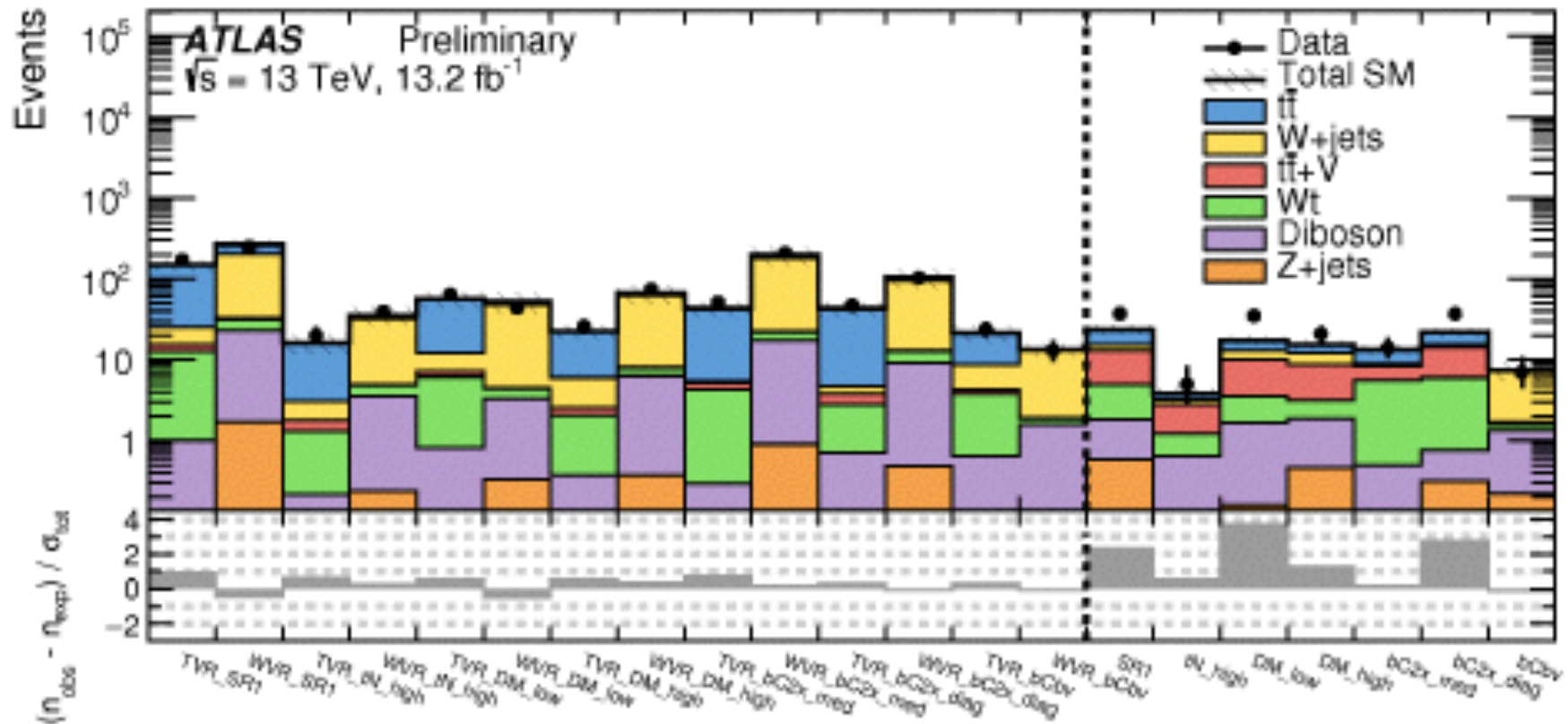


$$M_{T2}^2 \equiv \min_{\vec{p}_1 + \vec{p}_2 = \vec{p}_T} \left[\max \{ m_T^2(p_{Tl-}, \vec{p}_1), m_T^2(p_{Tl+}, \vec{p}_2) \} \right]$$

$$m_T^2 = 2p_T^{\text{lep}} E_T^{\text{miss}} [1 - \cos(\Delta\phi)]$$

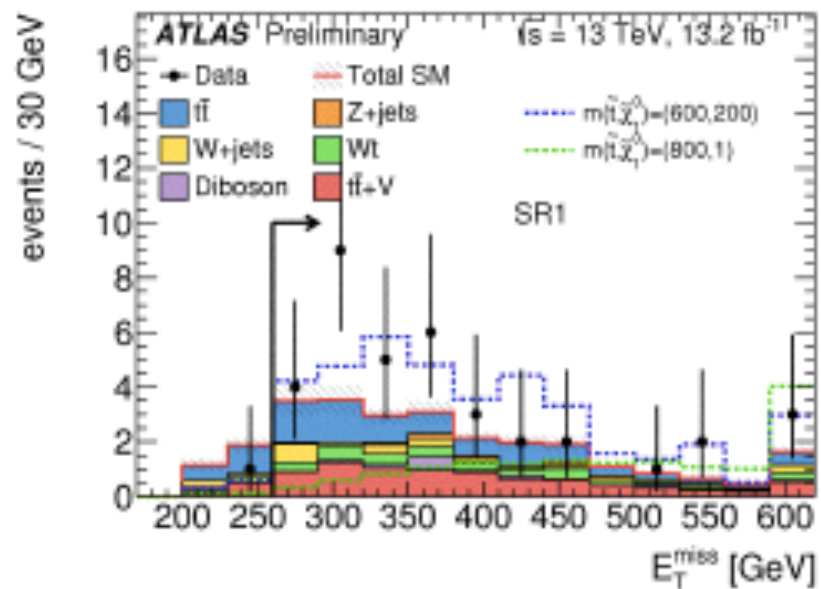
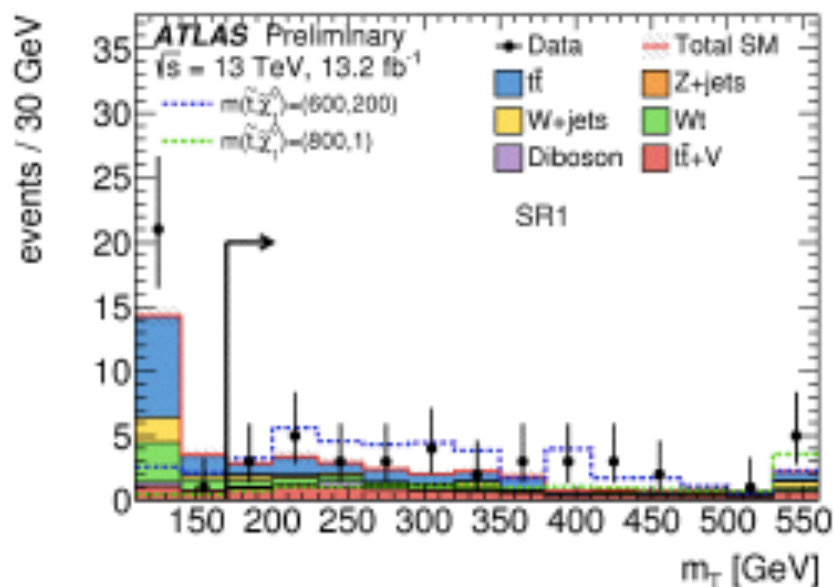
Signal region and limits stop 1L

Signal region and limits stop 1L

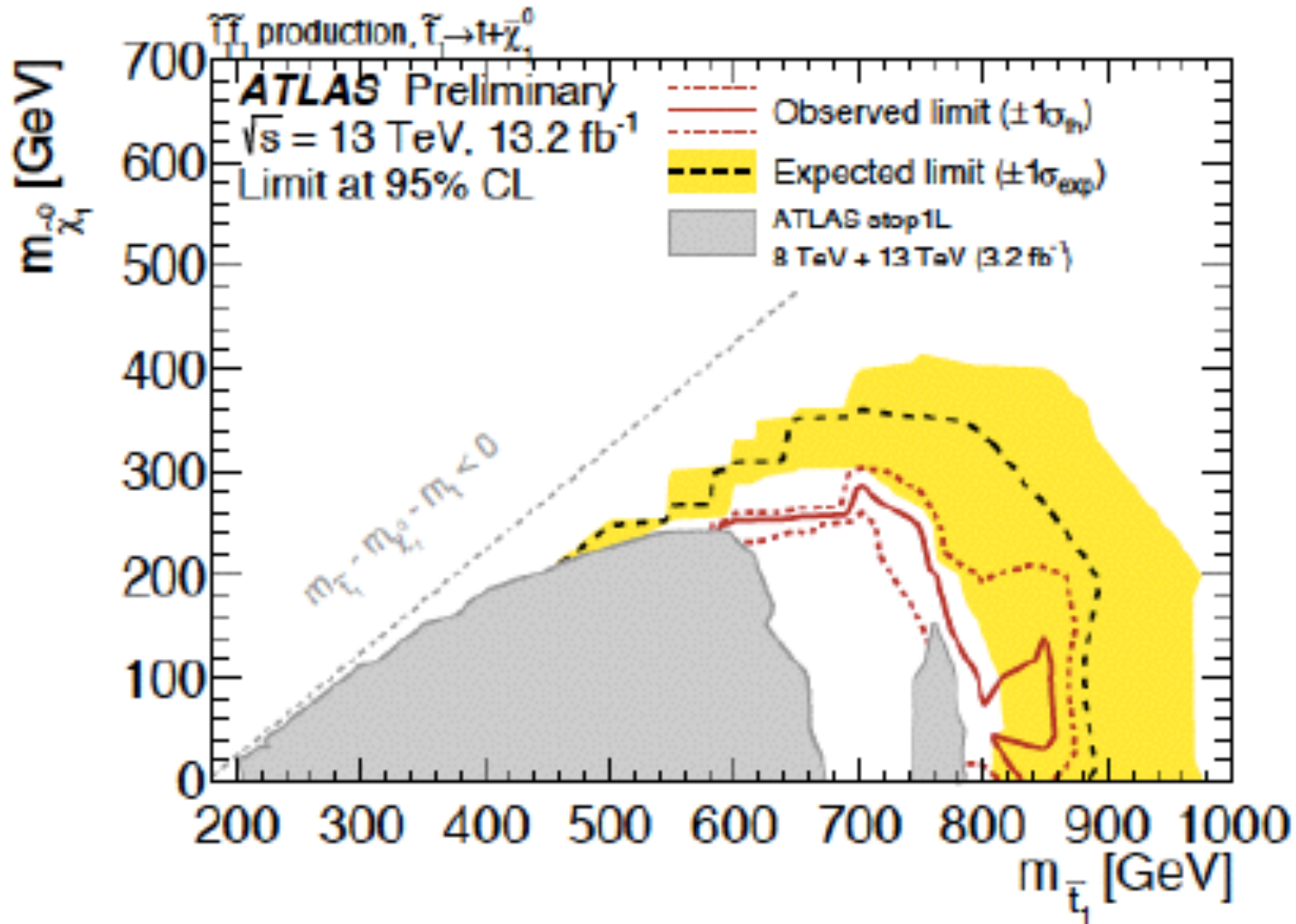


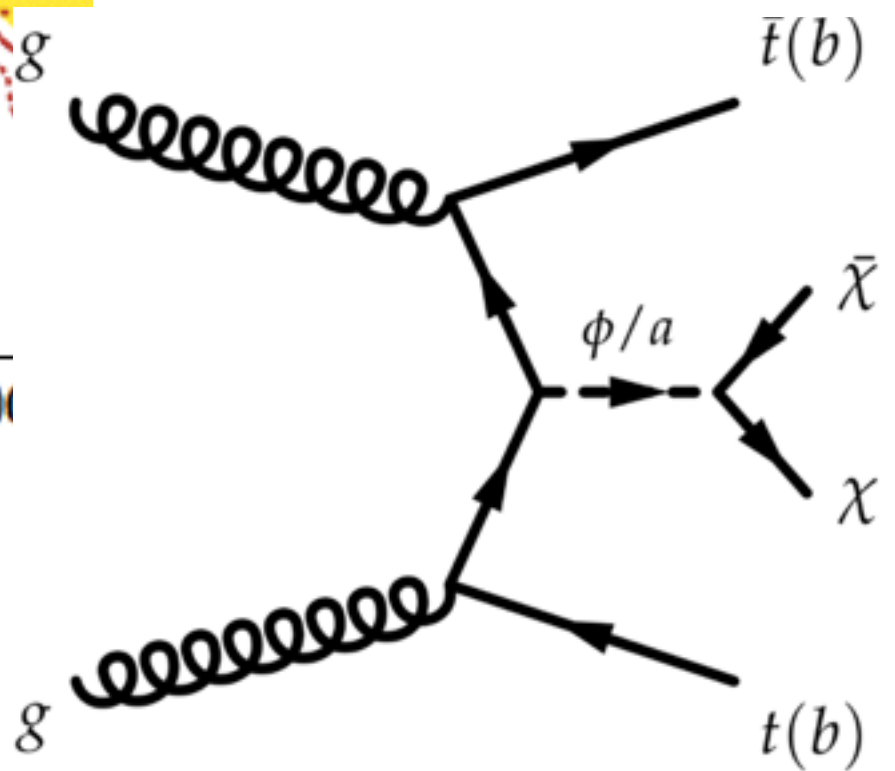
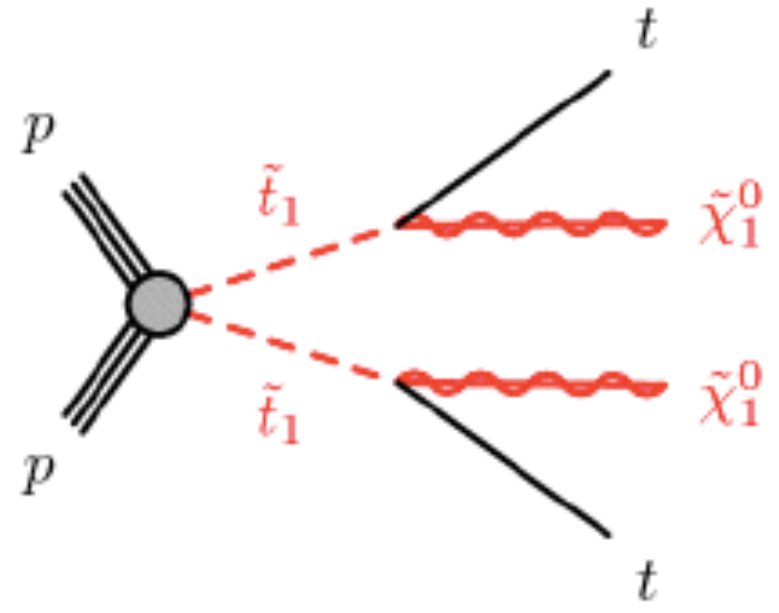
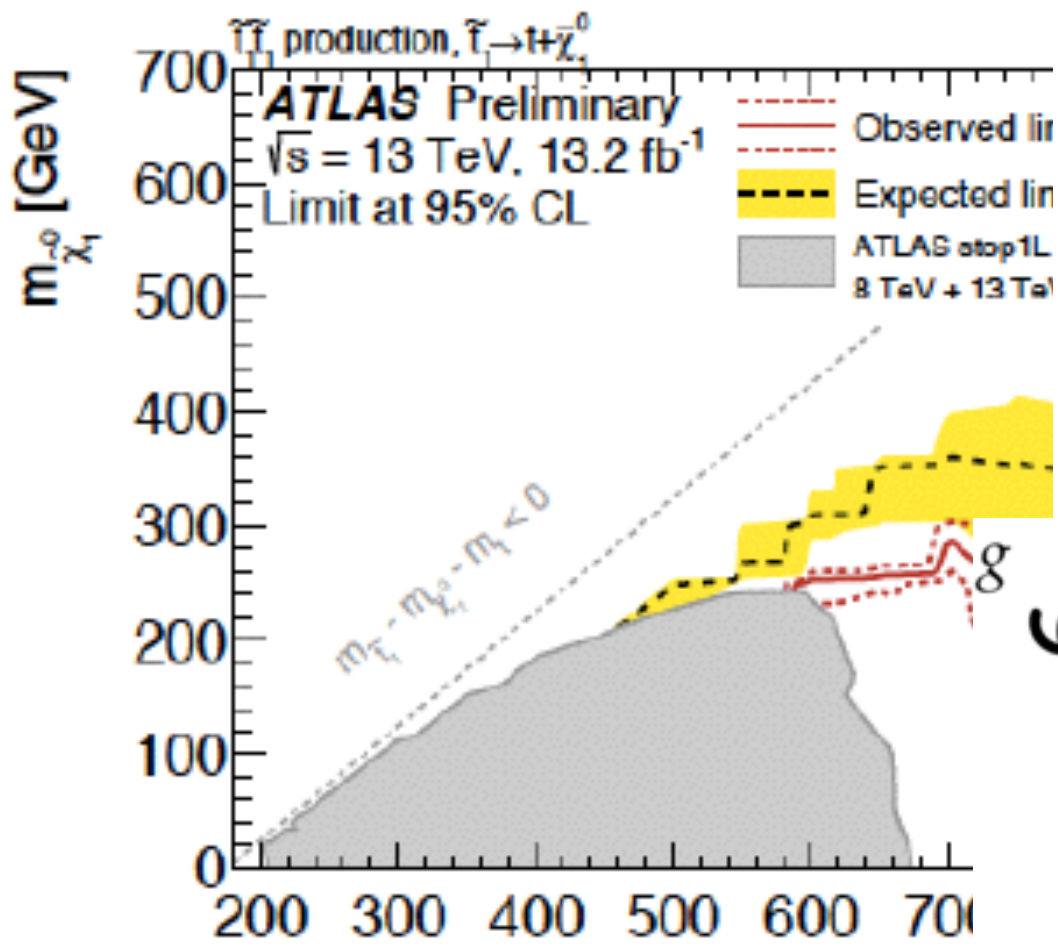
Signal region and limits stop 1L

Signal region and limits stop 1L



Signal region and limits stop 1L





Why Supersymmetry?

- ☑ Hierarchy problem
- ☑ Gauge coupling unification
- ☑ WIMP dark matter

For 30 years, experiments testing these suggestive reasons were “right around the corner”, and SUSY became the dominant BSM paradigm.

The experiments are now here.





“Strong.
Light.
Cheap.

Pick two.”

An evening with
Keith Bontrager

September 28, 2013
7:30 p.m.

801 Ridge Road
Wilmette, Illinois
847.920.9360
velosmith.com

velosmith
bicycle studio

Famous truism in
the MTB industry:

Impossible to have
simultaneously
strong, light, and
cheap components.





“Naturalness.
Unification.
Dark matter.

Pick two.”

An evening with
Nathaniel Craig

September 1, 2016
11:15 a.m.

St. Catherine's College
Oxford, UK

SEARCH2016

Supersymmetry in light
of data:

Impossible to have a
simple theory that is
natural, unifies, and
gives WIMP DM.

Picking two is a useful
guide.



Where we are now: Higgsinos

“Natural SUSY”

5 TeV



\tilde{g}



\tilde{t}



$\tilde{t}_L, \tilde{t}_R, \tilde{b}_L$

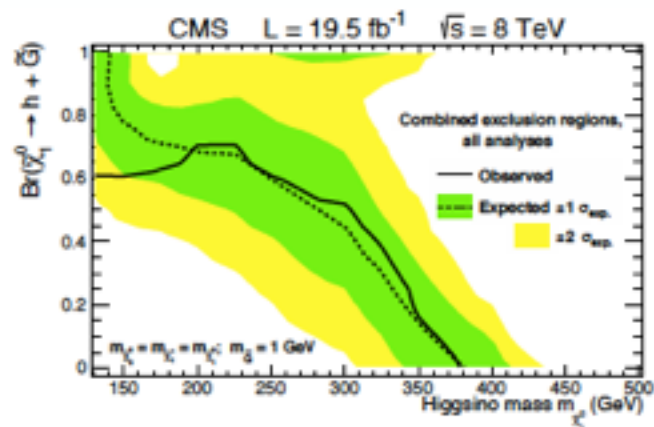
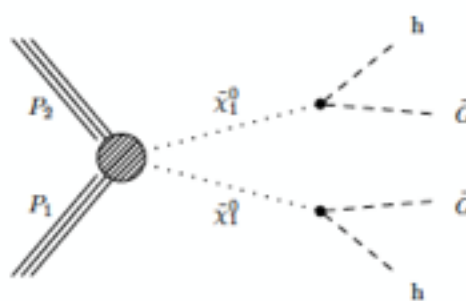


\tilde{h}



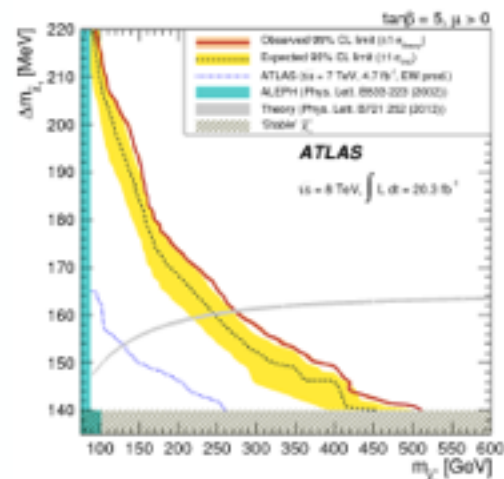
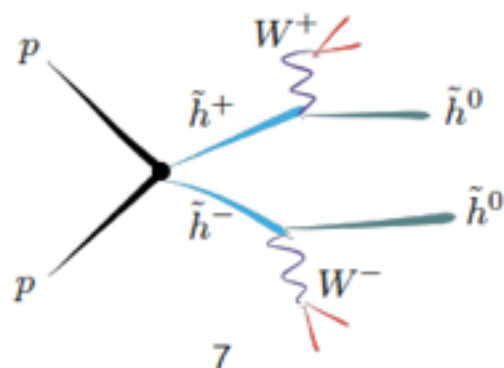
h

Lots of searches...



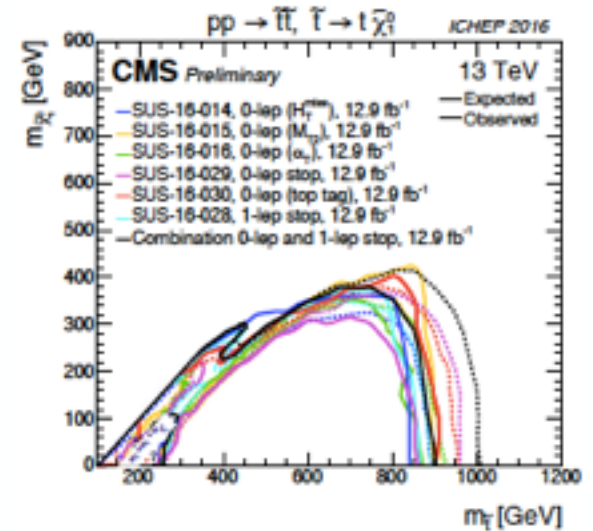
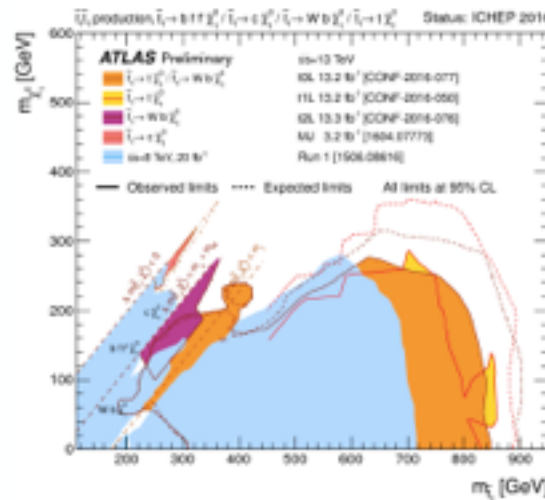
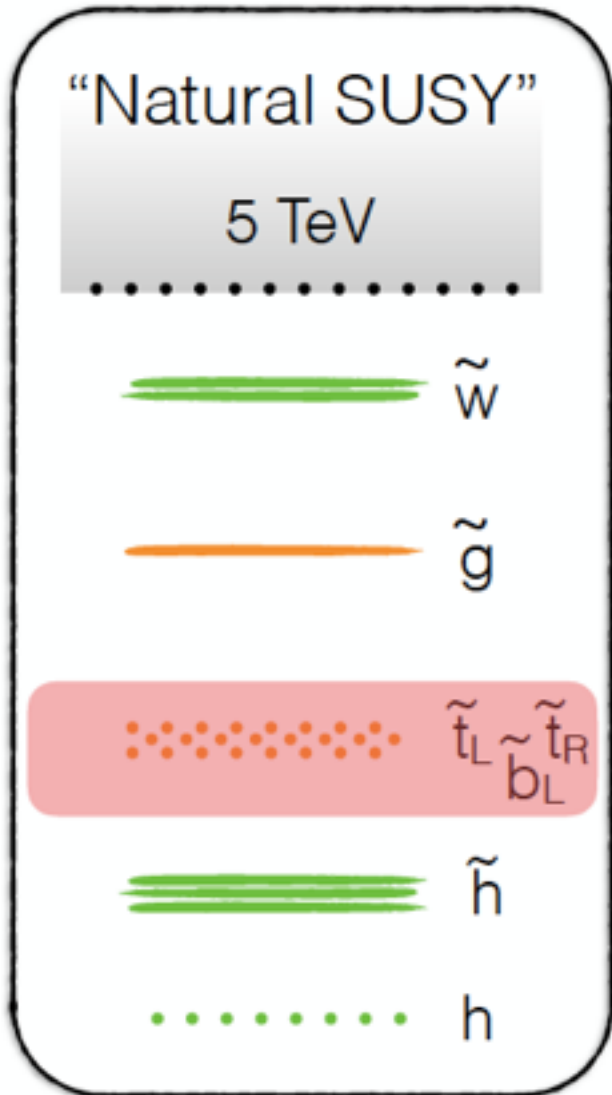
...but no irreducible limits

Chargino-neutralino splitting in pure higgsino multiplet: 355 MeV
[Thomas, Wells '98]



Where we are now: Stops

$$\delta m_H^2 \sim -\frac{3}{8\pi^2} y_t^2 (m_{Q_3}^2 + m_{u_3}^2 + |A_t|^2) \log(\Lambda/\text{TeV})$$

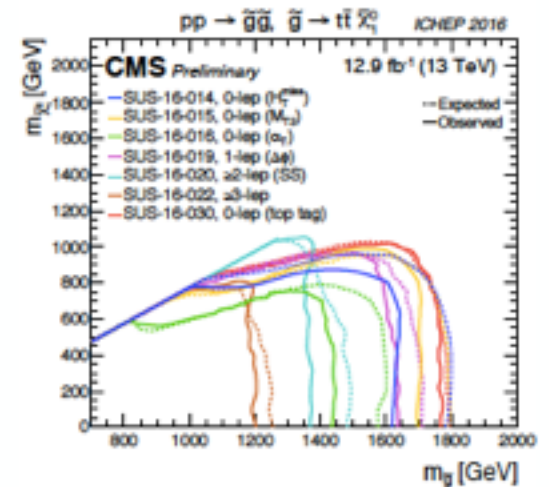
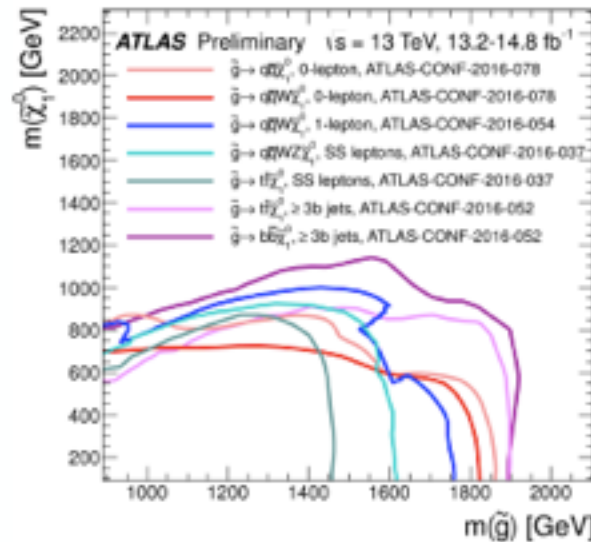
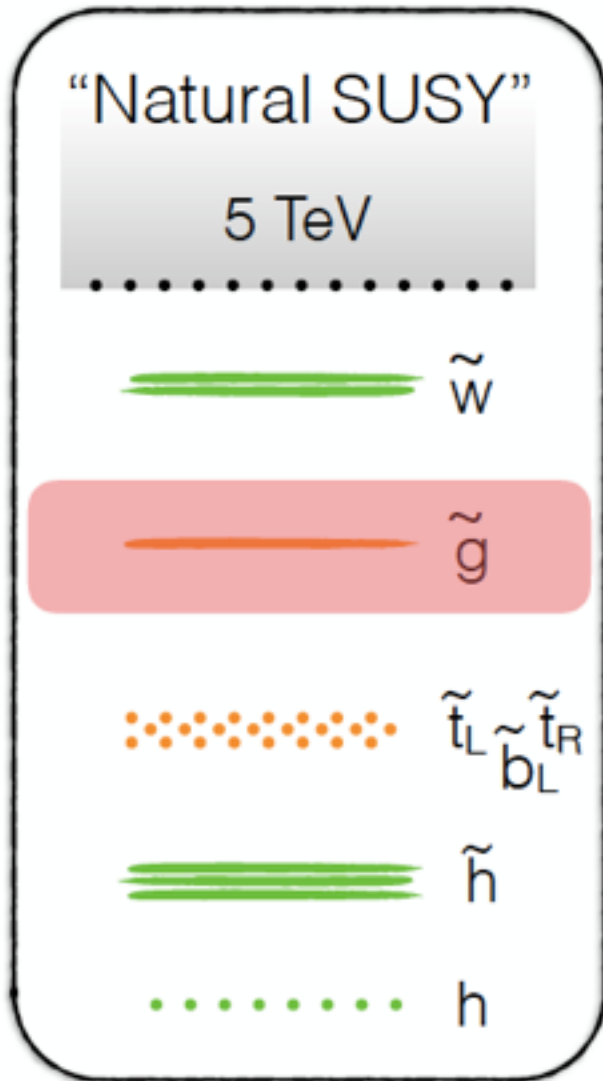


Generic limit > 800 GeV (both stops)
 → $\Delta \sim 90$ (1% tuning)
 ($\Lambda = 100 \text{ TeV}$)*

Where we are now: Gluinos

$$\delta m_H^2 \sim -\frac{2}{\pi^2} y_t^2 \left(\frac{\alpha_s}{\pi} \right) |M_3|^2 \log^2(\Lambda/\text{TeV})$$

Leads to “ $m_{\tilde{t}} \gtrsim M_3/2$ ”



Generic limit $> 1800 \text{ GeV}$
 $\rightarrow \Delta \sim 57$ (2% tuning)
 $(\Lambda = 100 \text{ TeV})$

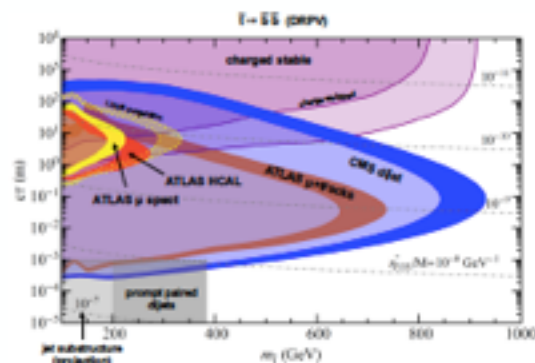
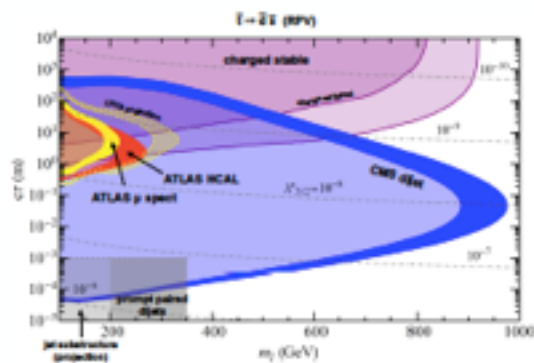
Naturalness & Unification

Violate the parity that gave you DM (RPV).



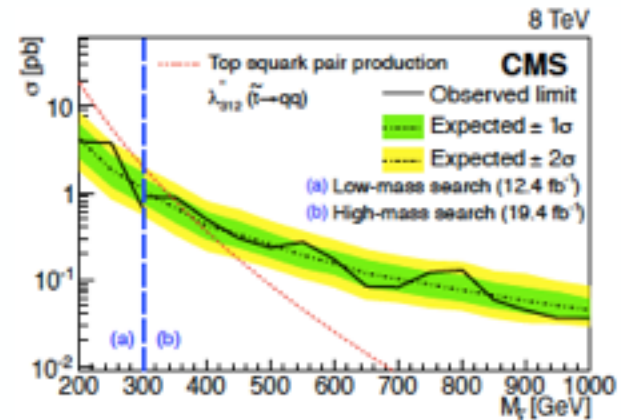
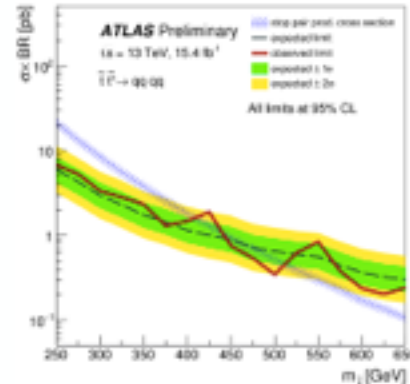
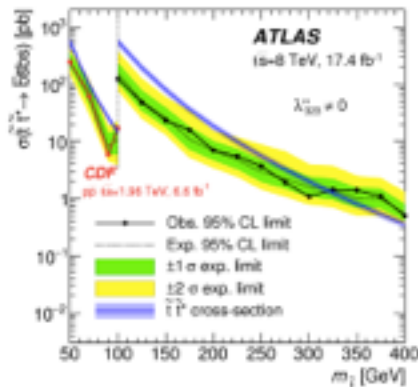
Lesson 1: Leptonic RPV kills you quickly.

RPV	LFV $pp \rightarrow \tau, \mu, X, \tau, \mu \rightarrow \tau \nu \tau \nu$	$g_{\tau \mu \tau}$	-	-	\sqrt{s}	$\lambda_{\tau \tau \tau} = 0.11, \lambda_{\tau \tau \mu} = 0.07$
Minimal RPV CMSSM	$2 c_{\tau, \mu}$ (SB)	0-3 b	Yes	20.3	1.8 TeV	$m_0 = 0, m_{1/2} = 1 \text{ TeV}$
$X^0_1 \tau^+ \tau^- \rightarrow W^+ W^- \tau^+ \tau^- \rightarrow \tau \nu \tau \nu$	$4 c_{\tau, \mu}$	-	Yes	13.3	1.14 TeV	$m_0 = 400 \text{ GeV}, A_0 = 0, \mu = 1.2$
$X^0_1 \tau^+ \tau^- \rightarrow W^+ W^- \tau^+ \tau^- \rightarrow \tau \nu \tau \nu$	$3 c_{\tau, \mu} + \epsilon$	-	Yes	20.3	450 GeV	$m_0 = 0, m_{1/2} = 1 \text{ TeV}, A_0 = 0$
$RR: \tau \rightarrow qW$	0	4-5 large- k jets	-	14.8	1.88 TeV	$m_0 = 0, m_{1/2} = 1 \text{ TeV}, A_0 = 0$
$RR: \tau \rightarrow qW, X^0_1 \tau^+ \tau^- \rightarrow qW$	0	4-5 large- k jets	-	14.8	1.88 TeV	$m_0 = 800 \text{ GeV}$
$RR: \tau \rightarrow qW, X^0_1 \tau^+ \tau^- \rightarrow qW$	$2 c_{\tau, \mu}$ (SB)	0-3 b	Yes	13.2	1.3 TeV	$m_0 = 450 \text{ GeV}$
$X^0_1 \tau^+ \tau^- \rightarrow b\bar{b}$	0	2 jets + 2 b	-	15.4	410 GeV	450-510 GeV
$X^0_1 \tau^+ \tau^- \rightarrow b\bar{b}$	$2 c_{\tau, \mu}$	2 b	-	20.3	0.4-1.0 TeV	870-1000 GeV



Lesson 2: Displacement kills you quickly. [Liu, Tweedie 1503.05923]

Lesson 3: Even UDD RPV no longer gets you out of jail for free.

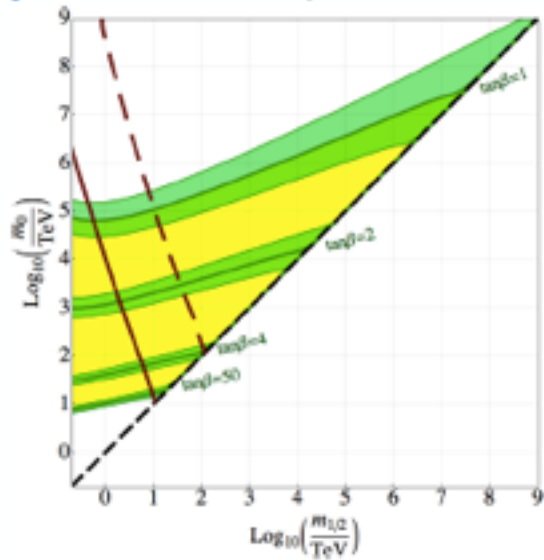


Unification & Dark Matter



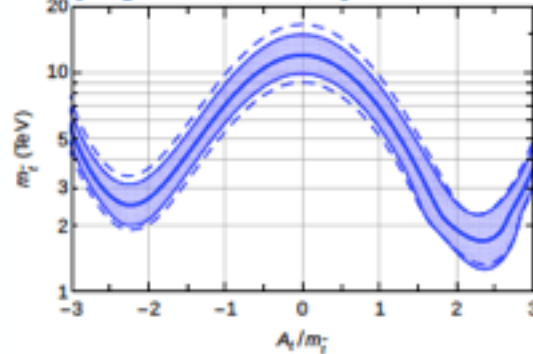
Forget the scalars, keep the fermions at a TeV.

[Arvanitaki, NC, Dimopoulos, Villadoro '12]



Use the MSSM Higgs mass as a guide; it was telling us that stops were above ~ 1 TeV all along.

[Vega, Villadoro '15]



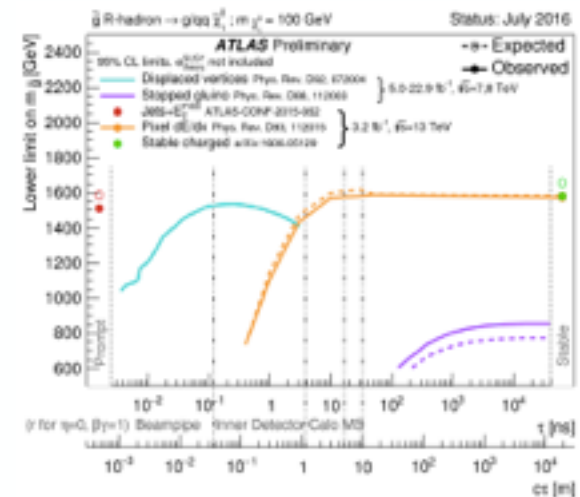
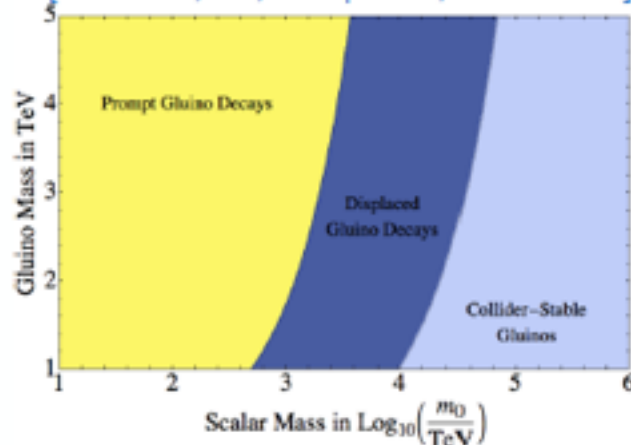
Pressure from unification & DM to keep higgsinos & gauginos beneath ~ 10 TeV

“Split / mini-split”

Canonical opportunity:
search for gluinos.
Opportunity, not
guarantee; not required to
be within kinematic reach

(*Request: what is the ct reach of “prompt” searches?)

[Arvanitaki, NC, Dimopoulos, Villadoro '12]



Naturalness & Dark Matter



Bring in new charged states at the TeV scale.

Dirac gauginos

$$m_{\tilde{\tau}} \neq M_3/2$$

SUSY broken by a D-term

$$\mathcal{D} \equiv \frac{1}{8} \langle D^2 \bar{D}^2 V' \rangle > 0$$

$$W \supset \frac{W'_\alpha W_j^\alpha A_j}{M} \rightarrow \mathcal{L} \supset \frac{\mathcal{D}}{M} \lambda \tilde{a}$$

Scalar masses radiative

$$\tilde{m}_i^2 \sim \frac{\alpha_i}{\pi} m_D^2 \log(m_a^2/m_D^2)$$

$$\text{Minimally } m_a \sim 2m_D$$

$$\text{so } m_{\tilde{\tau}} \sim M_3/5$$

Decouple gluinos! Predict new adjoint scalars

Global symmetry for Higgsinos

$$m_H^2 \neq \mu^2$$

[NC, Howe; Cohen, Kearney, Luty]

SUSY Higgs is a pNGB associated w/ spontaneously broken global symmetry

$$\mathcal{G} \rightarrow \mathcal{H}$$

μ term an invariant of

$$\mathcal{G}$$

doesn't contribute to Higgs potential

No problem w/ higgsinos @ TeV, but predict new states associated w/ global symmetry.

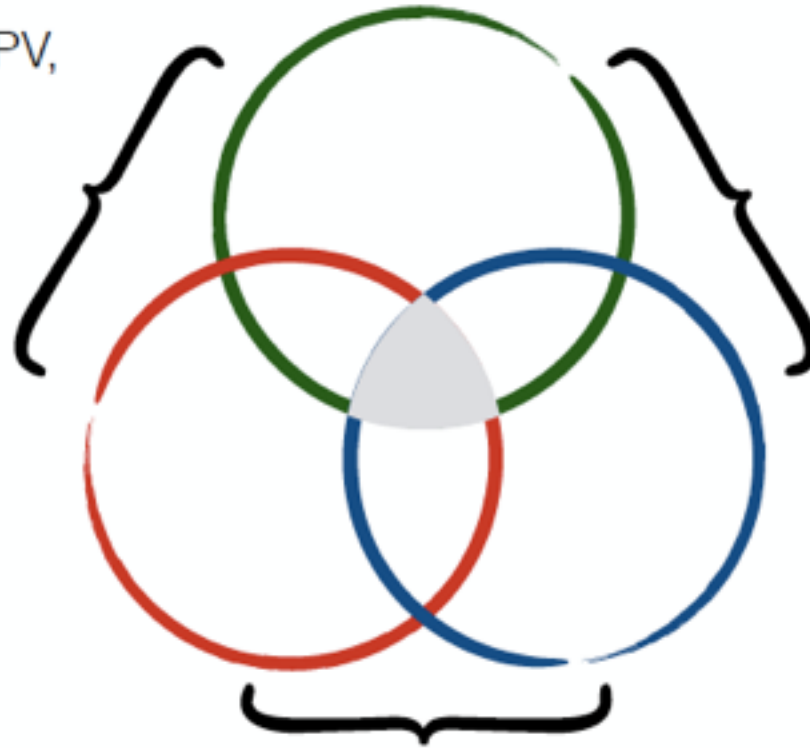
No local 4D SUSY

- E.g. 5D SUSY on S_1/Z_2 , SUSY broken by BCs.
- Spectrum finite, no large logs. (Often) dirac gauginos.
- Geography/localization can distinguish generations.
- Zero modes not supersymmetric ("hard breaking" for higgsino).
- Scale is $1/R \sim 5$ TeV
- Analogous models in 4D

Look for the new stuff. Often large cross sections or resonantly produced.

Naturalness & Unification

- Light-flavor UDD RPV, LQD w/ taus
- RPV Higgsino
- Higgs properties
- <Your idea here>



Naturalness & Dark Matter

- Additional states near weak scale (sgluon, KK resonances, ...)
- Higgs properties
- <Your idea here>

Unification & Dark Matter

- Conventional split SUSY searches
- Pure wino, higgsino LSP
- Extended Higgs sector?
- <Your idea here>

Getting to know the new particle

ATLAS-CONF-2015-044, CMS-PAS-HIG-15-002, PRL 114 (2015) 191803, EPJC 75 (2015) 212, Phys. Lett. B 726 (2013), pp. 120-144

- SM is highly **predictive** for the Higgs boson: Only free parameter is the mass
- Measure the **mass** and **width**
- Measure the **production** rate
- Measure **spin** and **parity** (only elementary scalar):
- Measure **couplings** (including self-coupling)

$$J^{PC} = 0^{++}$$

H^0

$$J = 0$$

In the following H^0 refers to the signal that has been discovered in the Higgs searches. Whereas the observed signal is labeled as a spin 0 particle and is called a Higgs Boson, the detailed properties of H^0 and its role in the context of electroweak symmetry breaking need to be further clarified. These issues are addressed by the measurements listed below.

Concerning mass limits and cross section limits that have been obtained in the searches for neutral and charged Higgs bosons, see the sections “Searches for Neutral Higgs Bosons” and “Searches for Charged Higgs Bosons (H^\pm and $H^{\pm\pm}$)”, respectively.

H^0 MASS

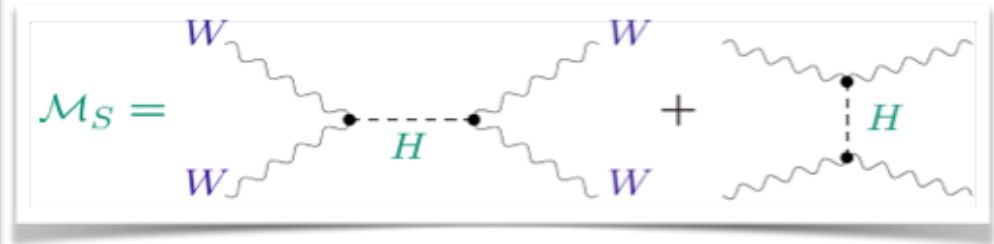
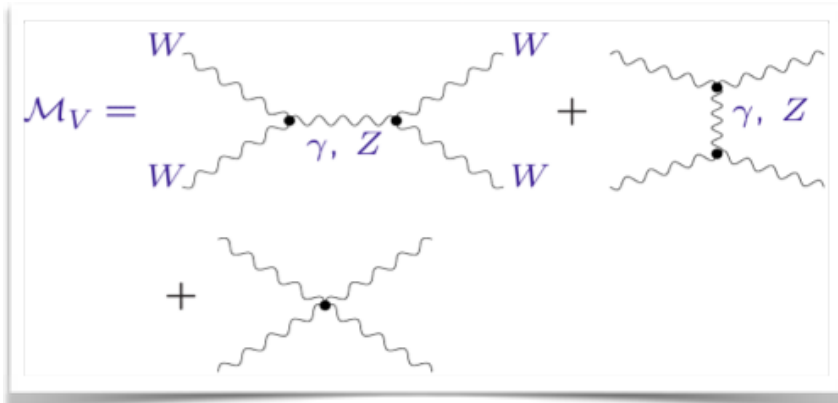
VALUE (GeV)	DOCUMENT ID	TECN	COMMENT
$125.09 \pm 0.21 \pm 0.11$	^{1,2} AAD	15B LHC	pp , 7, 8 TeV

• • • We do not use the following data for averages, fits, limits, etc. • • •



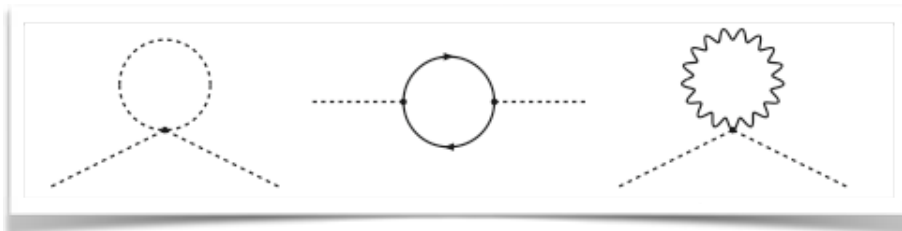
The Hierarchy Problem

- WW scattering violates unitarity above ~ 1 TeV
- New diagrams needed to regulate the cross section
- Adding diagrams with a scalar solves the problem



- BUT, the loop corrections to the Higgs boson mass are quadratically divergent, so although we've solved one problem, we're left with another

Planck scale



$$\Delta m^2 \propto \int^{\Lambda} \frac{d^4k}{(2\pi)^4} \frac{1}{k^2} \sim \frac{\Lambda^2}{16\pi^2}$$

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 LEP/LIBRARY
 LEP Note 440
 11.4.1993
 SCAN-0008106

PRELIMINARY PERFORMANCE ESTIMATES FOR A LEP PROTON COLLIDER

S. Myers and M. Schnell

1. Introduction

This analysis was stimulated by news from the United States where very large $p\bar{p}$ and pp colliders are actively being studied at the moment. Indeed, a first look at the basic performance limitations of possible $p\bar{p}$ or pp rings in the LEP tunnel seems overdue, however far off in the future a possible start of such a p-LEP project may yet be in time. What we shall discuss is, in fact, rather obvious, but such a discussion has, to the best of our knowledge, not been presented so far.

We shall not address any detailed design questions but shall give basic equations and make a few plausible assumptions for the purpose of illustration. Thus, we shall assume throughout that the maximum energy per beam is 8 TeV (corresponding to a little over 9 T bending field in very advanced superconducting magnets) and that injection is at 0.4 TeV. The ring circumference is, of course that of LEP, namely 26,659 m. It should be clear from this requirement of "Ten Tesla Magnets" alone that such a project is not for the near future and that it should not be attempted before the technology is ready.

Duration of projects /planning stability:
 First LHC workshop 1984 !

Higgs Overview

- Main analysis is a Multi-Variate Analysis (MVA)
 - MVA for photon ID and event classification
 - Fit mass distribution in 4 event classes based on a-photon MVA output + 2 $p\bar{p}$ categories
 - Improvement in expected limit ~15% over cut-based analysis
 - Cross-checked with an alternative background modelled extraction
 - Fit output of a $p\bar{p}$ MVA combining photon MVA and $m_{\tau\tau}$ using data on mass relations to constrain the background model
- Also cross-checked with a cut based analysis
 - Simple and robust
 - Cut based photon ID and event classification
 - Fit mass distribution in 4 event classes in a cleaner shape $m_{\tau\tau}$ categories and different signal over Background (S/B) + 2 $p\bar{p}$ categories
 - Published for ATLAS data
 - Phys Lett. B:66 (2012) 197-210, arXiv:1112.3547



H^0

$$J = 0$$

In the following H^0 refers to the signal that has been discovered in the Higgs searches. Whereas the observed signal is labeled as a spin 0 particle and is called a Higgs Boson, the detailed properties of H^0 and its role in the context of electroweak symmetry breaking need to be further clarified. These issues are addressed by the measurements listed below.

Concerning mass limits and cross section limits that have been obtained in the searches for neutral and charged Higgs bosons, see the sections “Searches for Neutral Higgs Bosons” and “Searches for Charged Higgs Bosons (H^\pm and $H^{\pm\pm}$)”, respectively.

 H^0 MASS

<i>VALUE</i> (GeV)	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
125.09 ± 0.21 ± 0.11	^{1,2} AAD	15B LHC	<i>pp</i> , 7, 8 TeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

It's real ... it's in the PDG

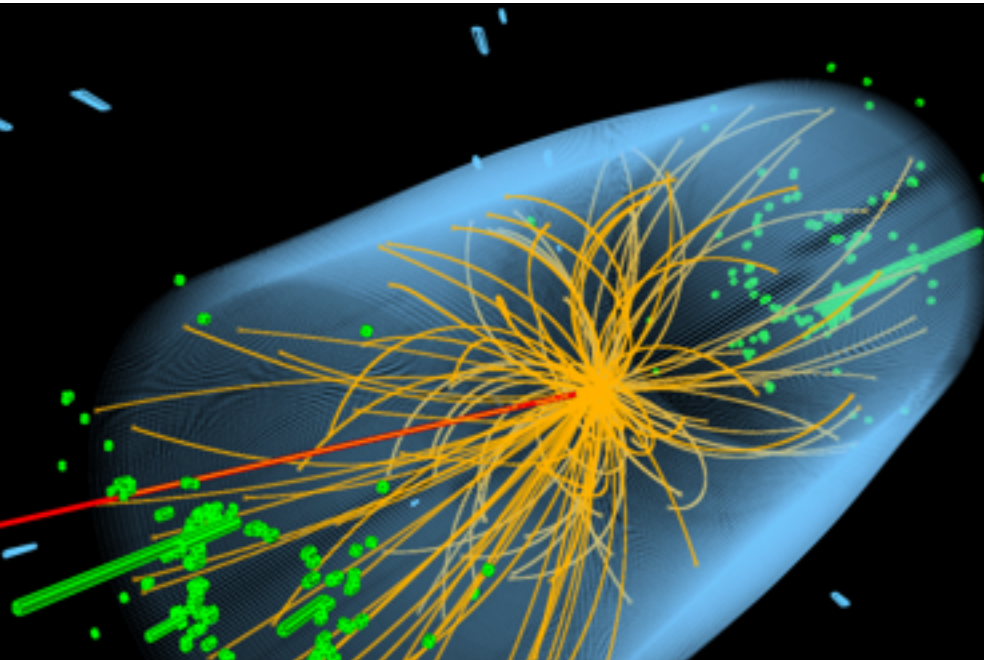
How many Higgs' ?



CMS Experiment at the LHC, CERN

Data recorded: 2012-Jun-05 09:58:43.400262 GMT[11:58:43 CEST]

Run / Event: 195592 / 61758463



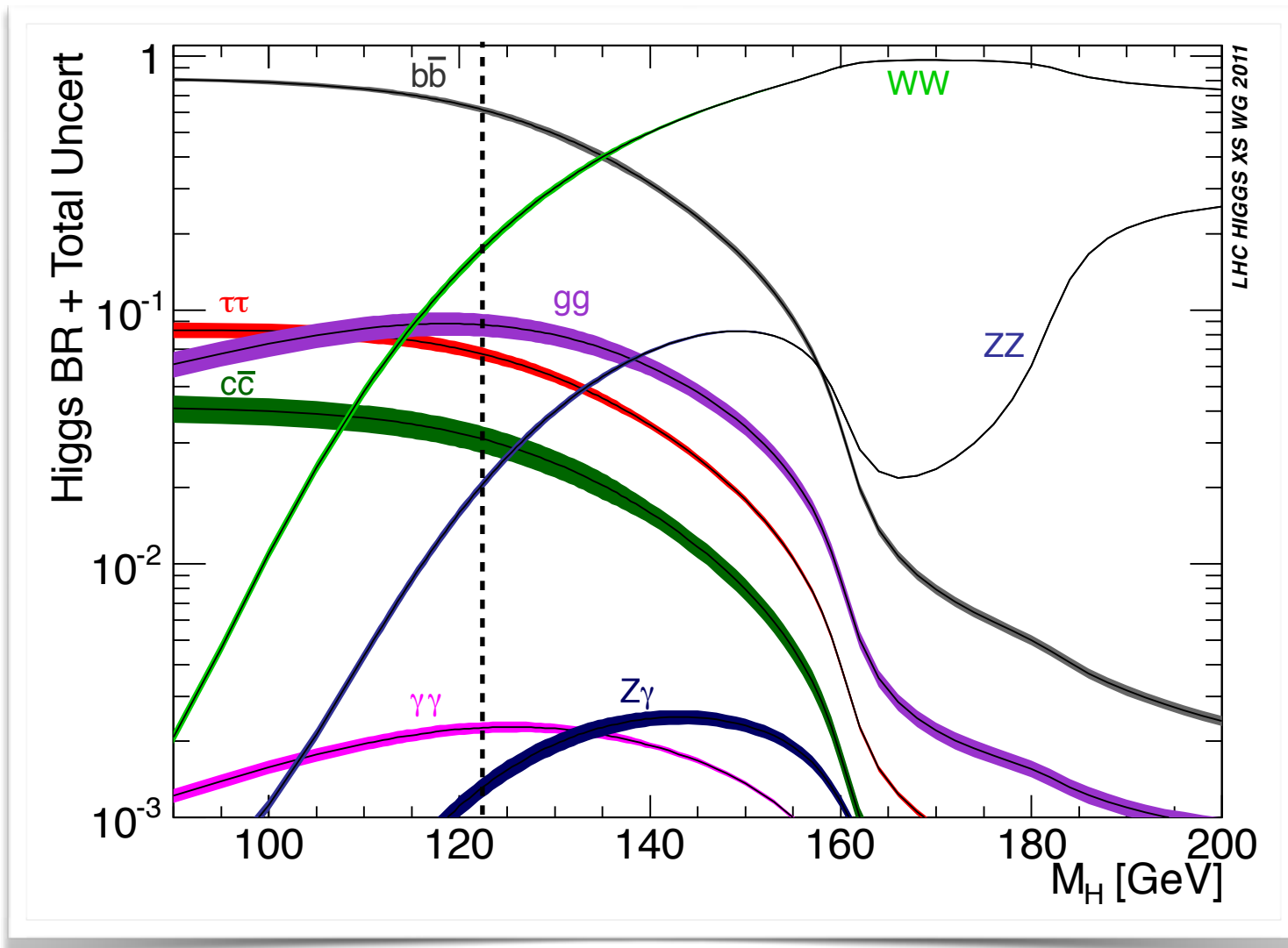
	Energy	Luminosity	Dates	Number of Higgs bosons
Run-1	7 TeV	$\sim 5 \text{ fb}^{-1}$	2010-2011	$\sim 80\text{k}$
	8 TeV	$\sim 20 \text{ fb}^{-1}$	2012-2014	$\sim 450\text{k}$
Run-2	13 TeV	$\sim 9 \text{ fb}^{-1}$	2015-	$\sim 400\text{k}$

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Medium-Angle

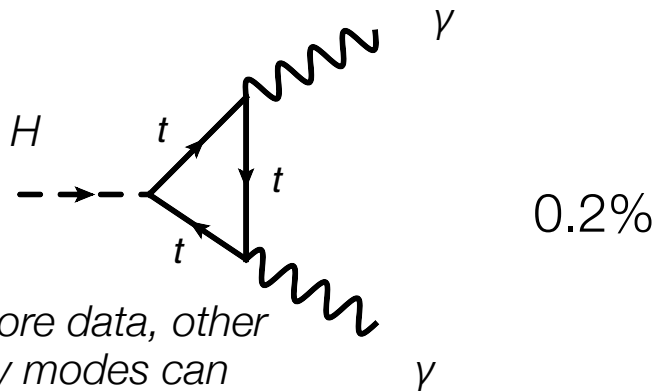
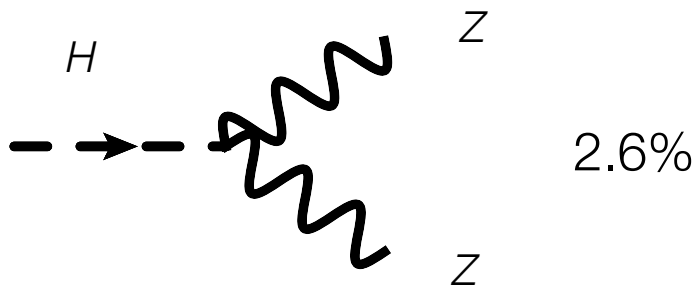
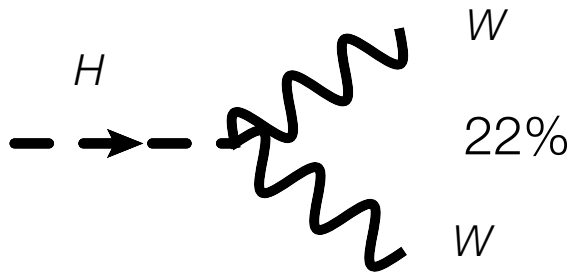
How hard can it be to find 500k particles?

Incredibly Lucky?

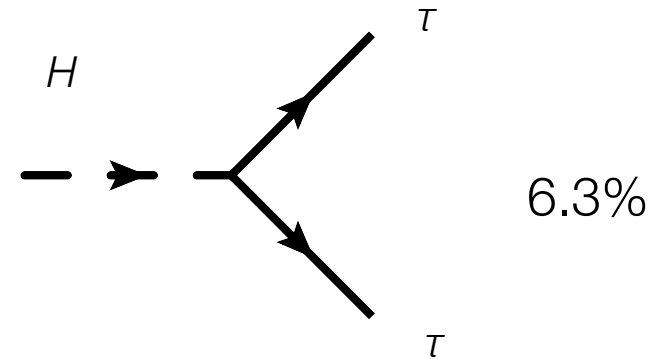
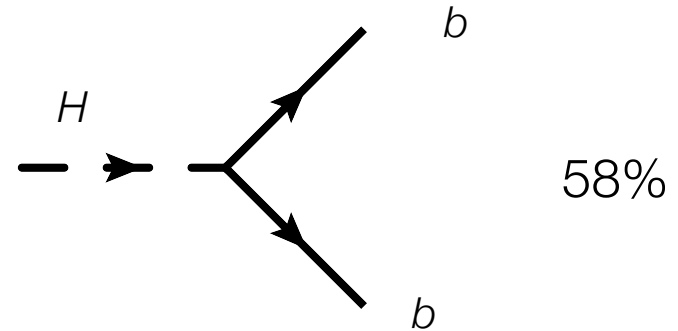


Nature just happened to choose a Higgs mass for which almost all experimental channels are accessible

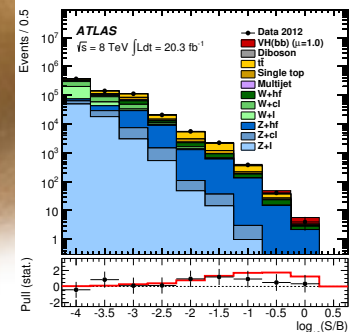
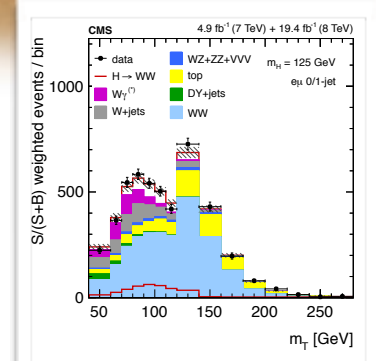
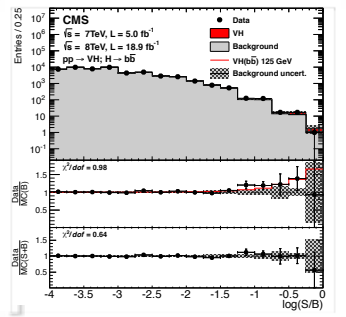
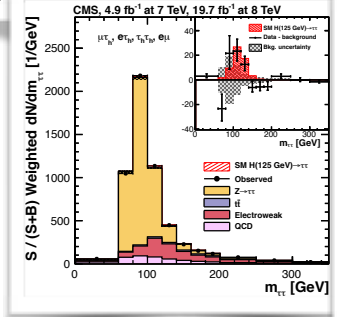
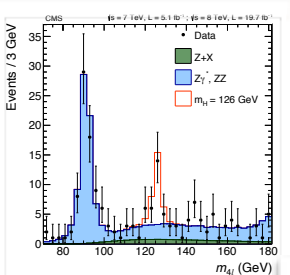
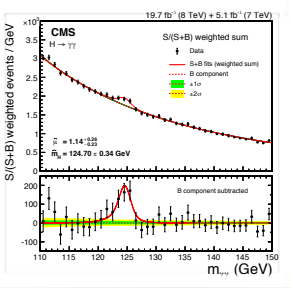
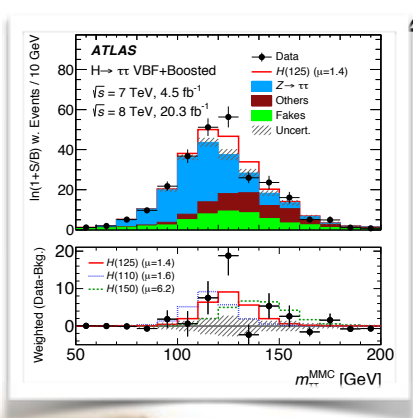
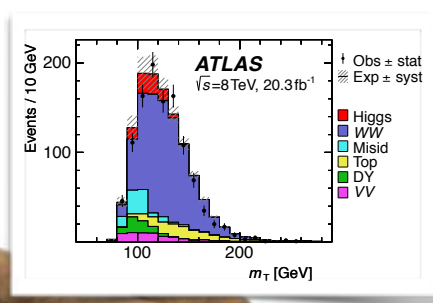
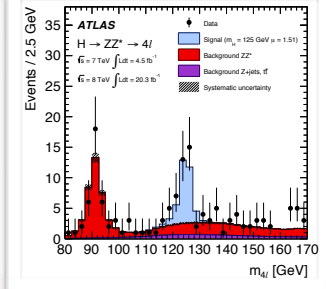
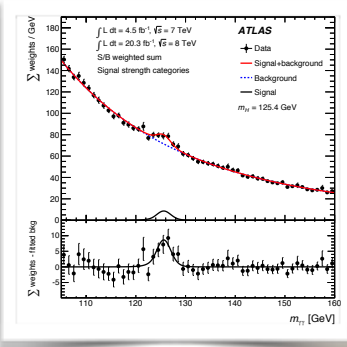
Higgs Decays: The Big 5



With more data, other decay modes can also be studied



Decay channel	Mass resolution
$H \rightarrow \gamma\gamma$	1-2%
$H \rightarrow ZZ \rightarrow \ell^+\ell^-\ell'^+\ell'^-$	1-2%
$H \rightarrow W^+W^- \rightarrow \ell^+\nu_\ell\ell'^-\bar{\nu}_{\ell'}$	20%
$H \rightarrow b\bar{b}$	10%
$H \rightarrow \tau^+\tau^-$	15%



Studied Higgs Modes

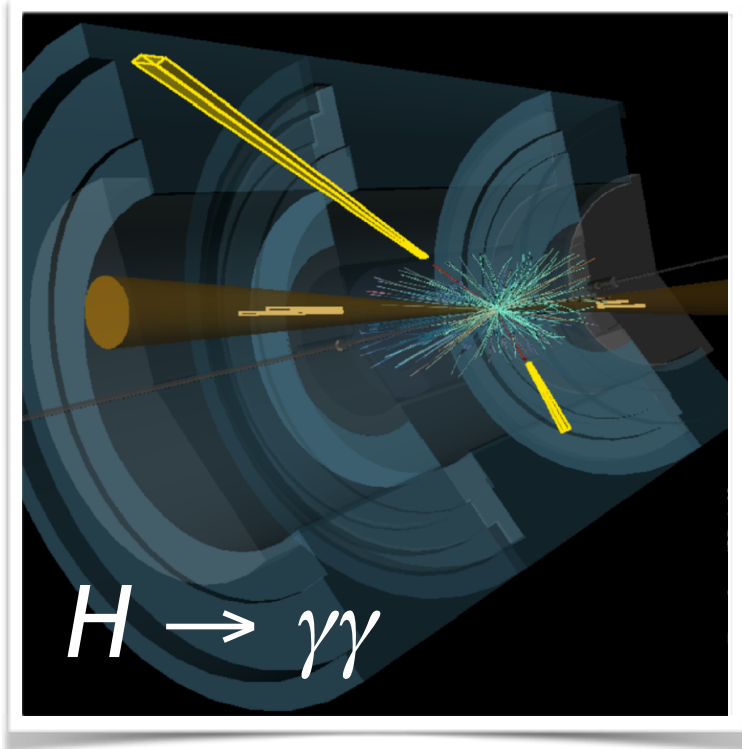
	mass range [GeV]	Branching Ratio [%]	Mass Resolution [%]	ggF	VBF	VH	ttH
bb	110-135	58	10				
WW	110-600	22	20				
$\tau\tau$	110-145	6,3	15				
ZZ	110-1000	2,6	1-2				
$\gamma\gamma$	110-150	0,2	1-2				

Searches for almost every decay modes and production channels*

*of course, there are some more exotic production and decay modes

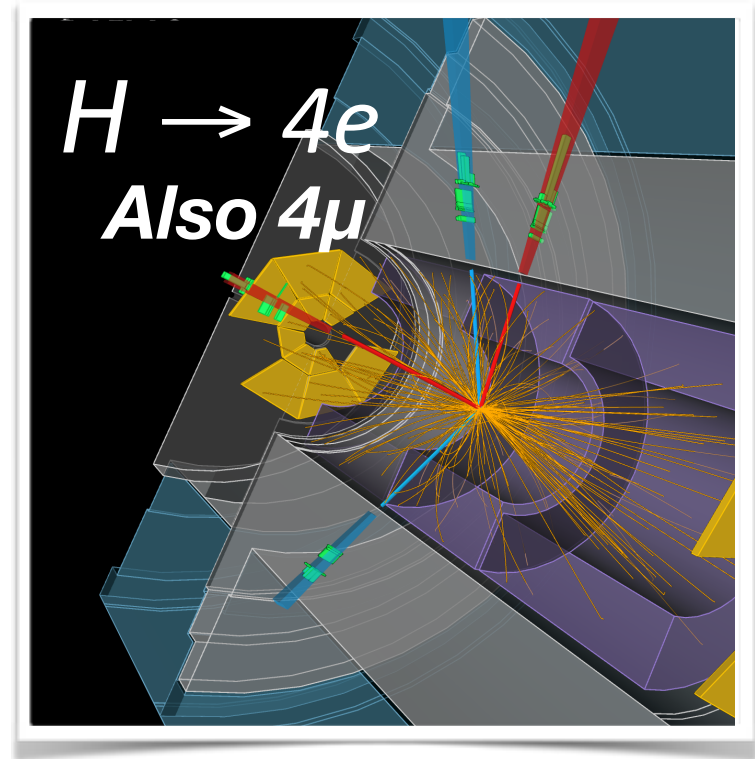
Discovery Channels

An excellent channel for
 $m_H = 125 \text{ GeV}$



$N_s \sim 500$

Golden channel over a
wide mass range

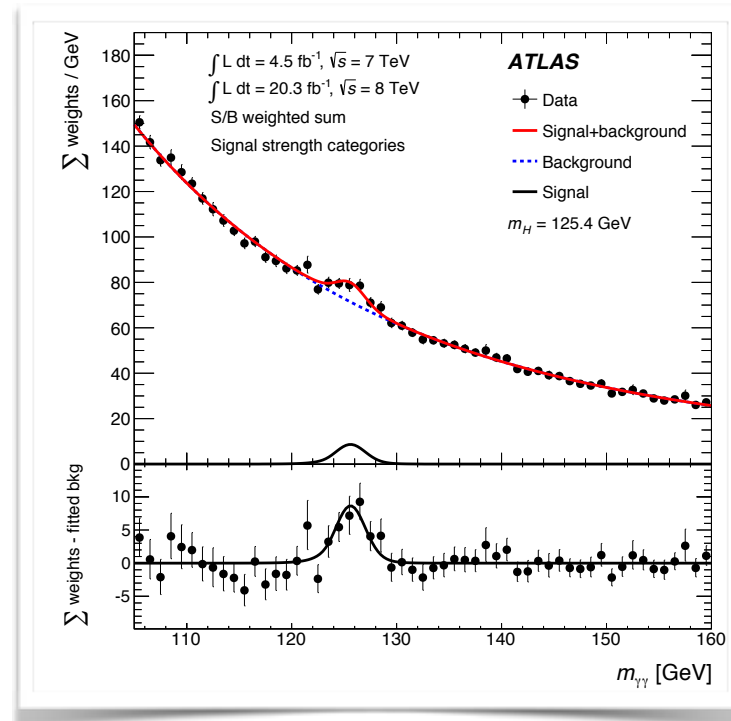
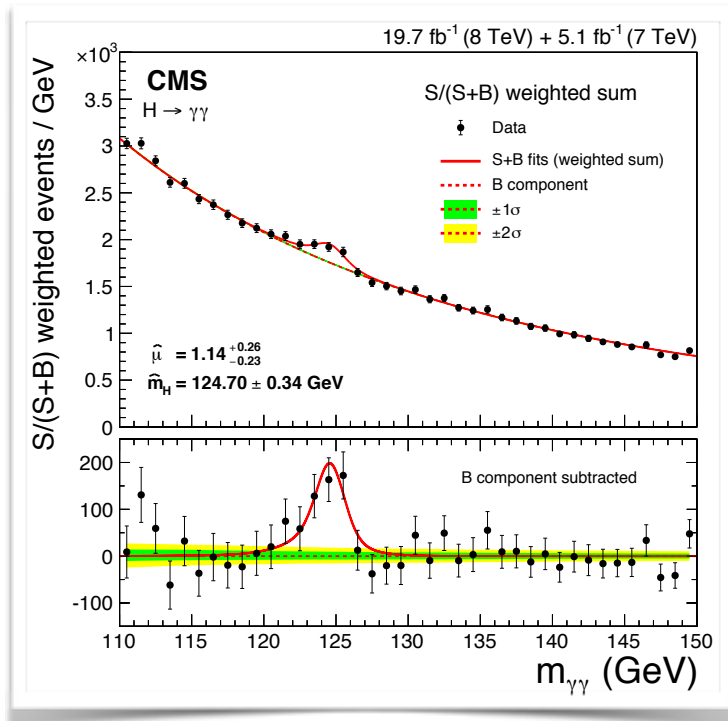


$N_s \sim 20-30$

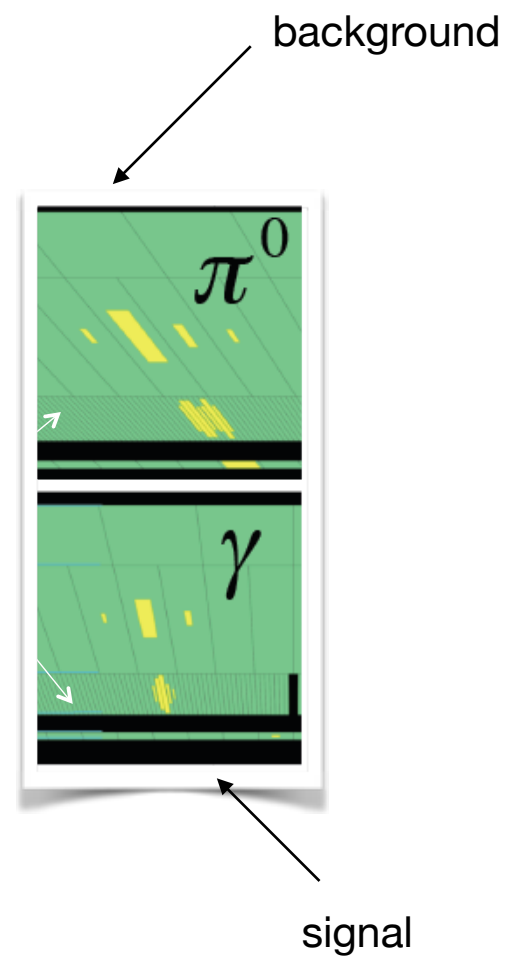
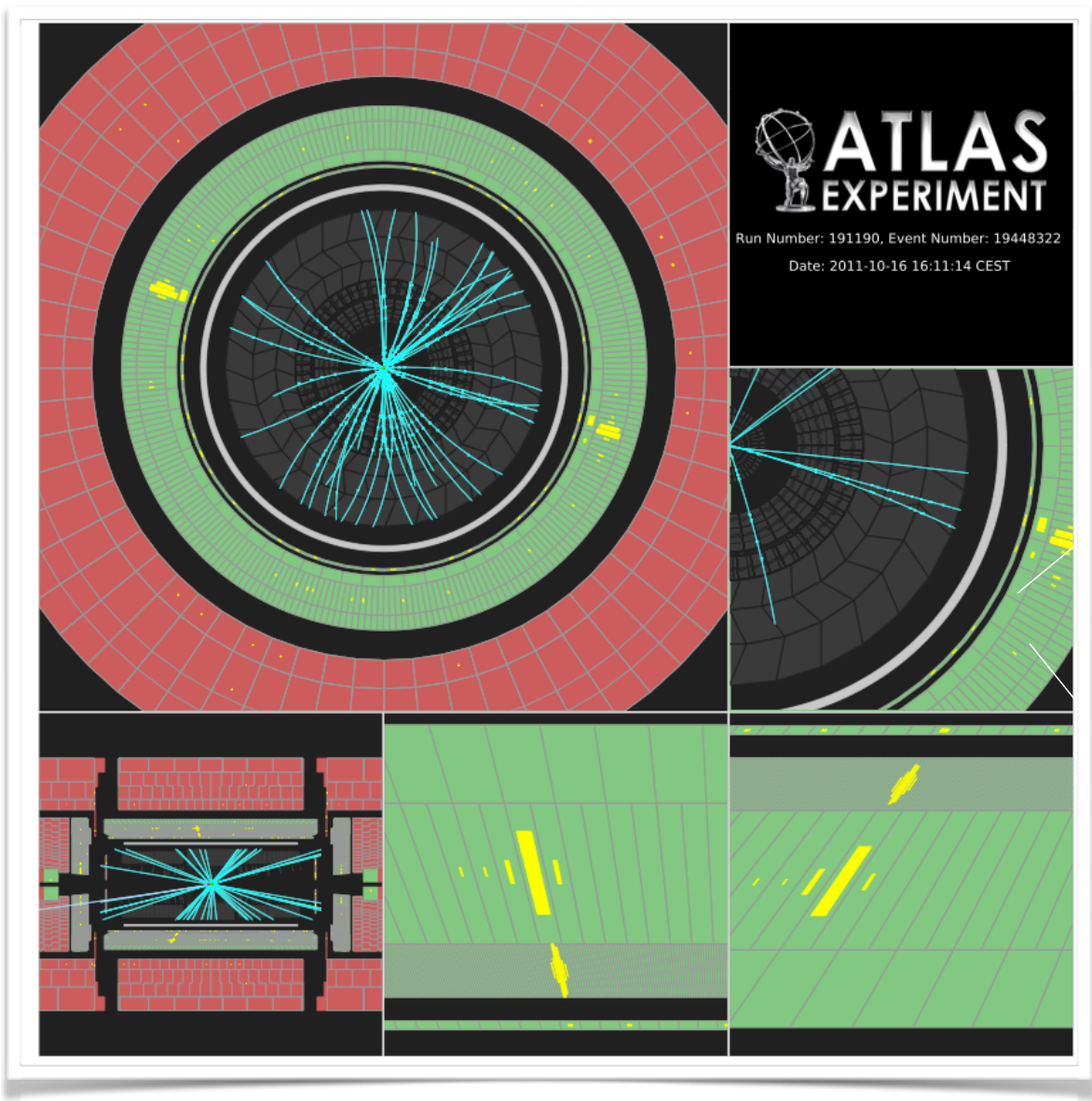
Simple channels with excellent mass resolution

Higgs to two photons ($H \rightarrow \gamma\gamma$)

- A good **discovery** final state
 - Resonance on top of a **smooth background**
 - Excellent Higgs mass **resolution**
- Large backgrounds: need good **photon identification**
 - Key consideration in calorimeter design

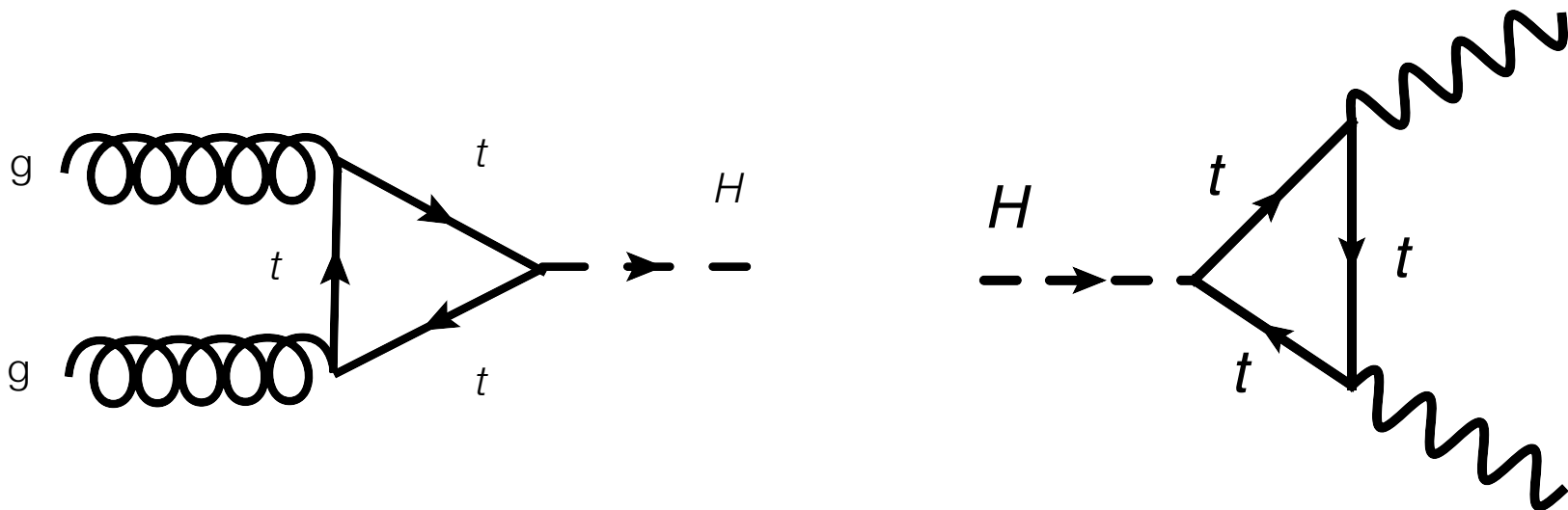


H → $\gamma\gamma$: Signal and Background



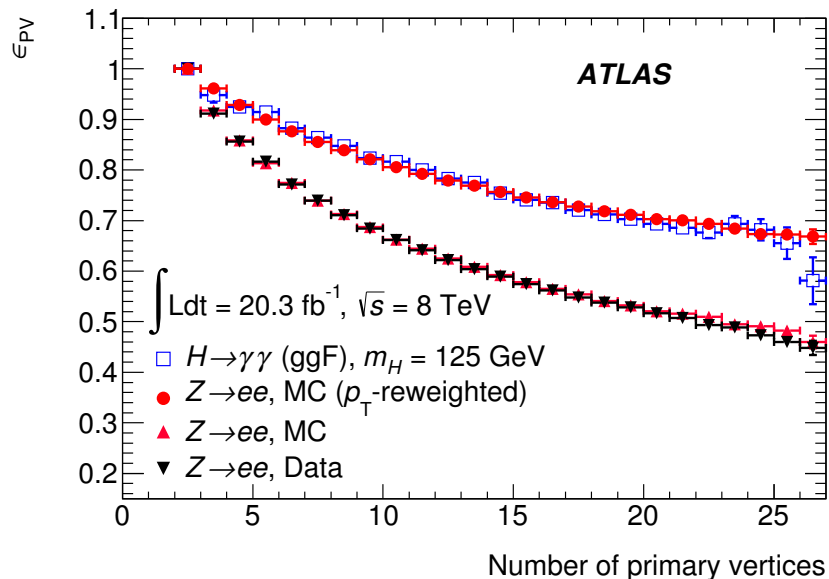
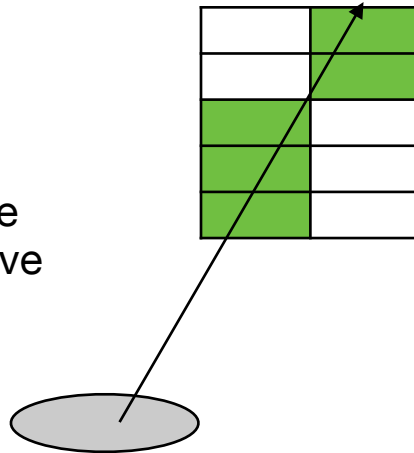
Production and Decay

- No coupling of Higgs to gluons
 - Main **production** through a loop containing top (and bottom) quarks
 - Cross-section depends on Higgs coupling to **top**
- **No coupling of the Higgs to photons**
 - Decay through loops containing tops and W bosons
 - Decay depends on coupling to **top** and **W** boson
 - Small **branching ratio** (0.2%)

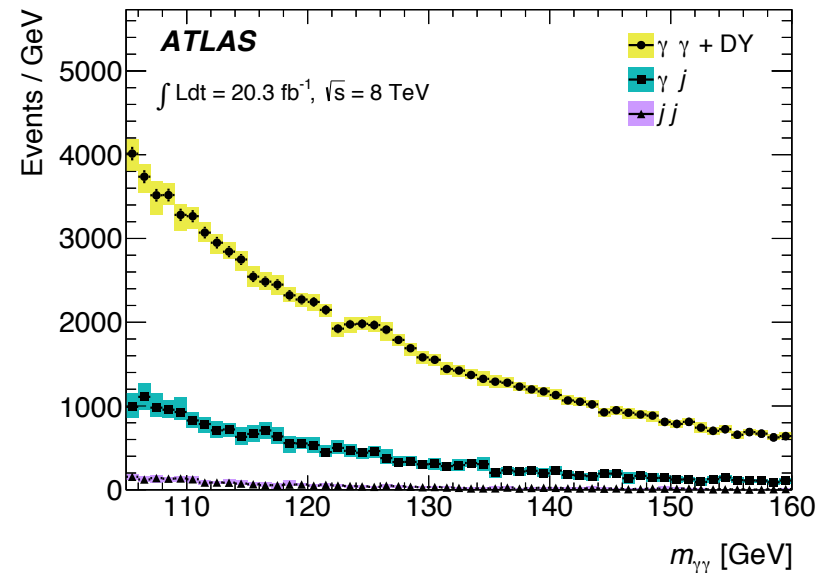


Backgrounds: pile up and jets

ATLAS uses longitudinal segmentation of the calorimeter to improve vertex resolution

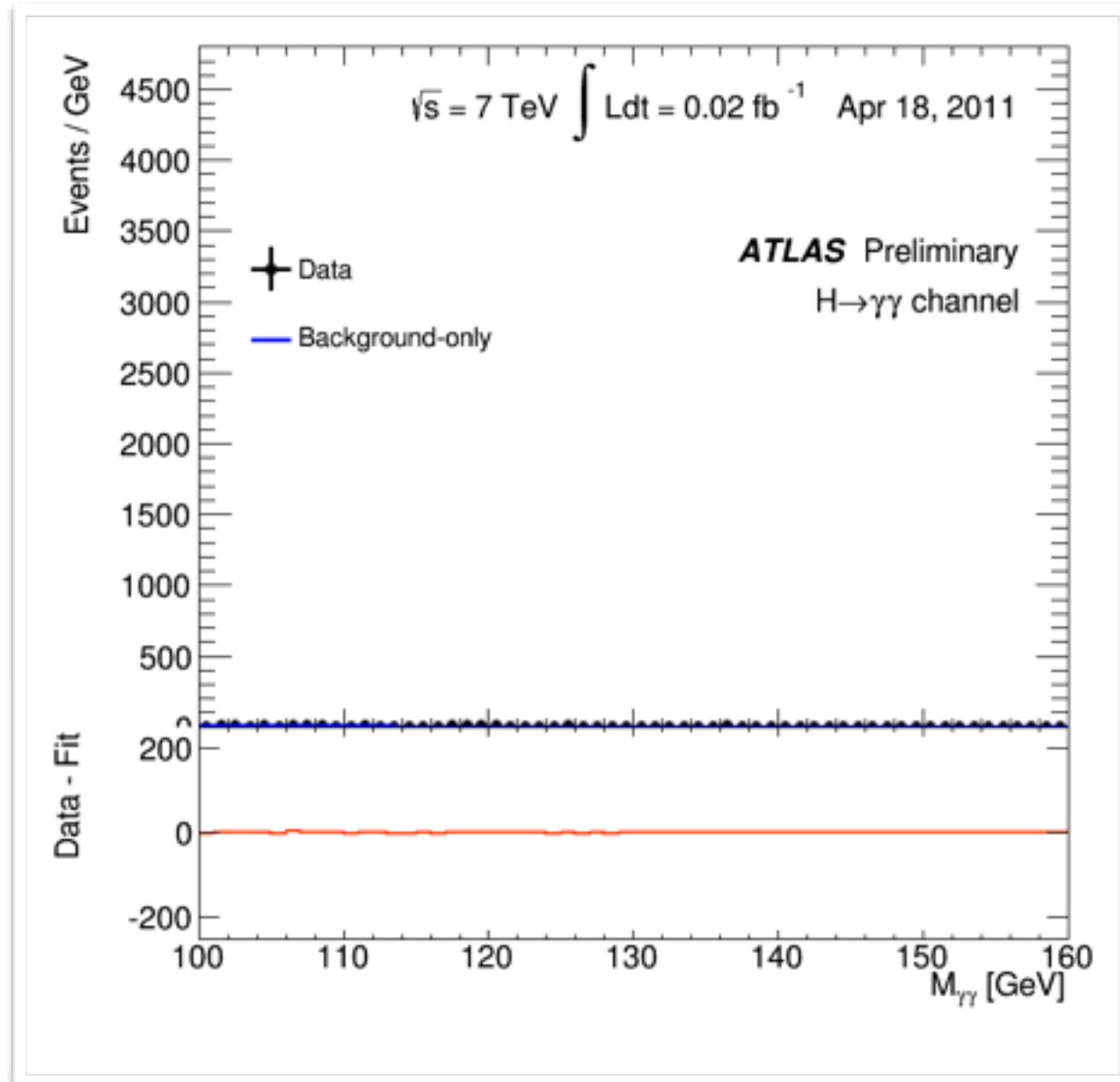


$$m_{\gamma\gamma} = \sqrt{2E_1 E_2 (1 - \cos \alpha)}$$

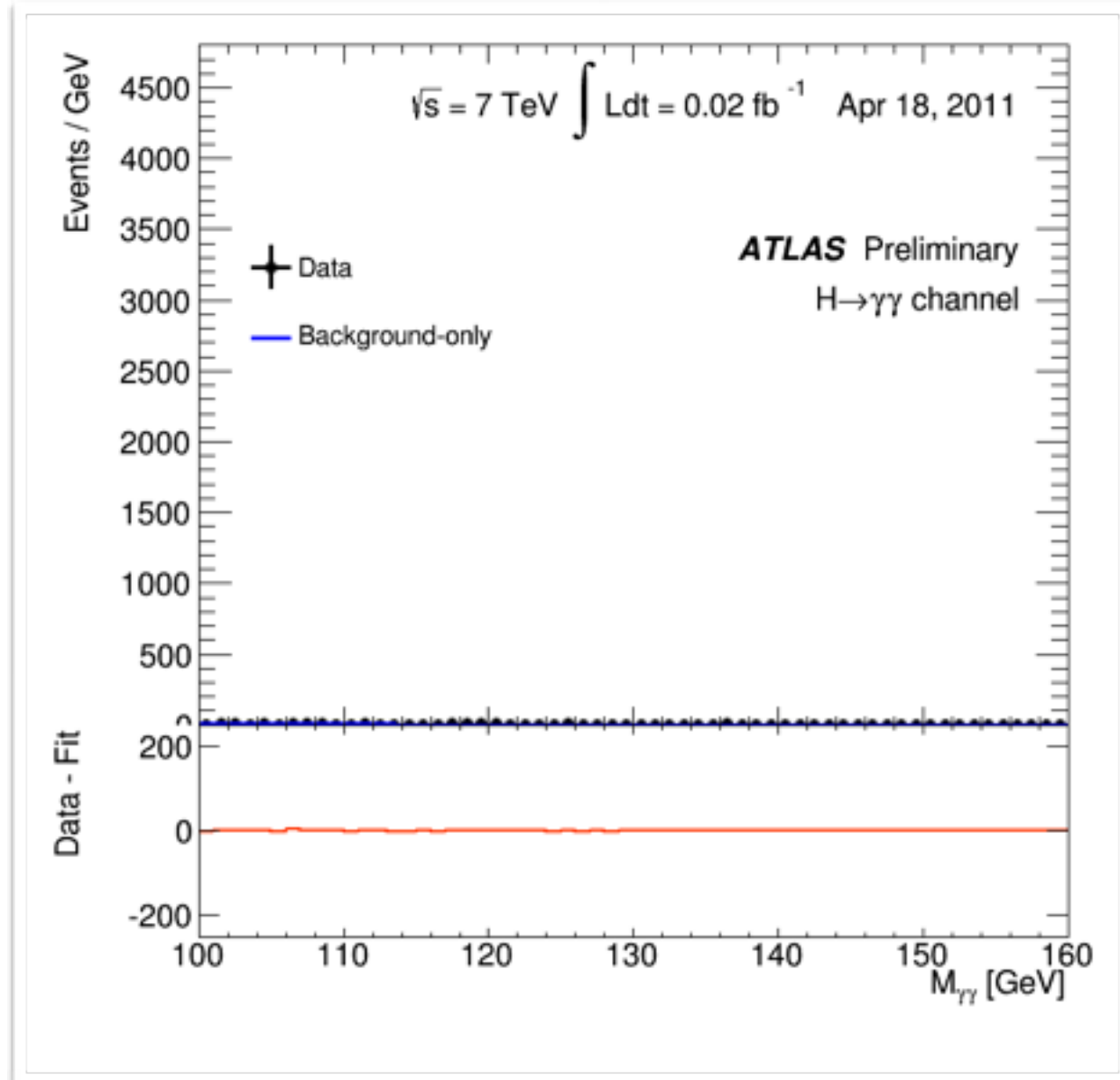


Main background from real photons
Small contribution from jets faking photons

Higgs to two photons ($H \rightarrow \gamma\gamma$)



Higgs to two photons ($H \rightarrow \gamma\gamma$)

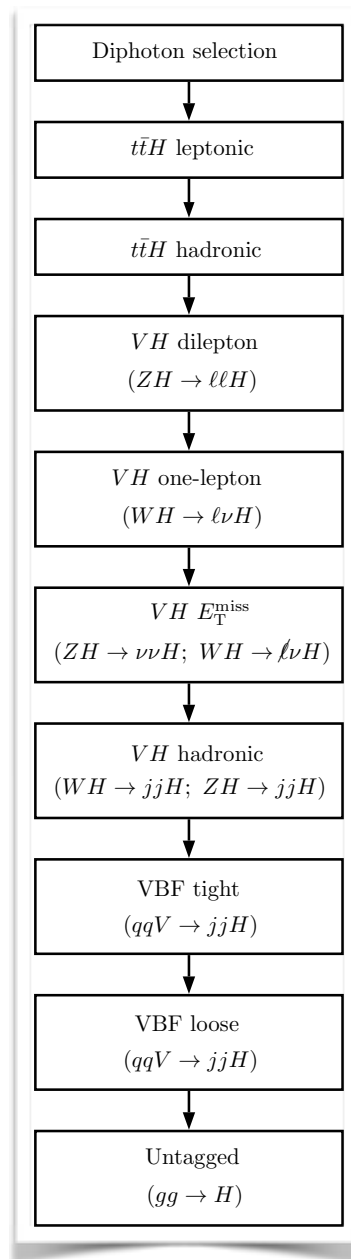


Why Categories ?

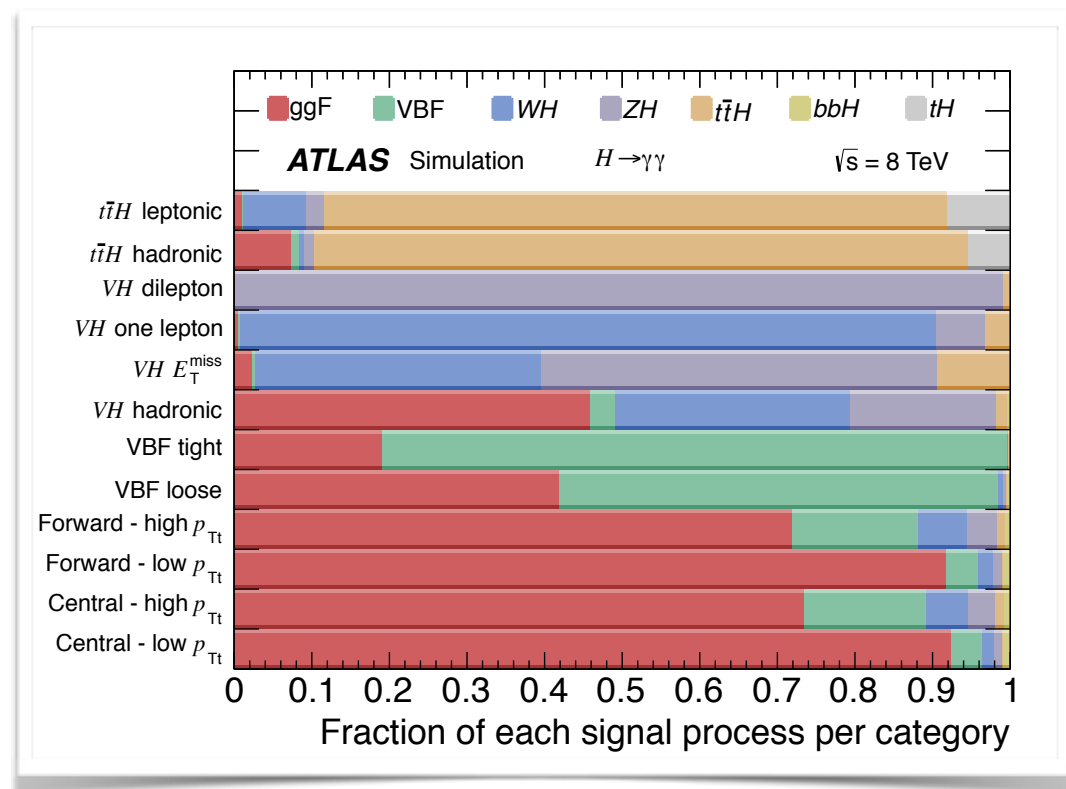
- Most LHC Higgs analyses use categories to improve sensitivity
- Main strategy is to separate events with different significance $\sim S/\sqrt{B}$
- Differences can depend on resolution, background type or size, signal production mechanism or systematic uncertainties
- Take a simple example with two categories:
 - C1: $s=12$ and $b=60$
 - C2: $s=18$ and $b=40$
- Inclusively we have
 - $s = 30$
 - $b = 100$
 - Significance of 3σ
- Now calculate for the two categories
 - C1: 2.85σ
 - C2: 1.55σ
 - Combined significance: 3.24

Improved
significance!

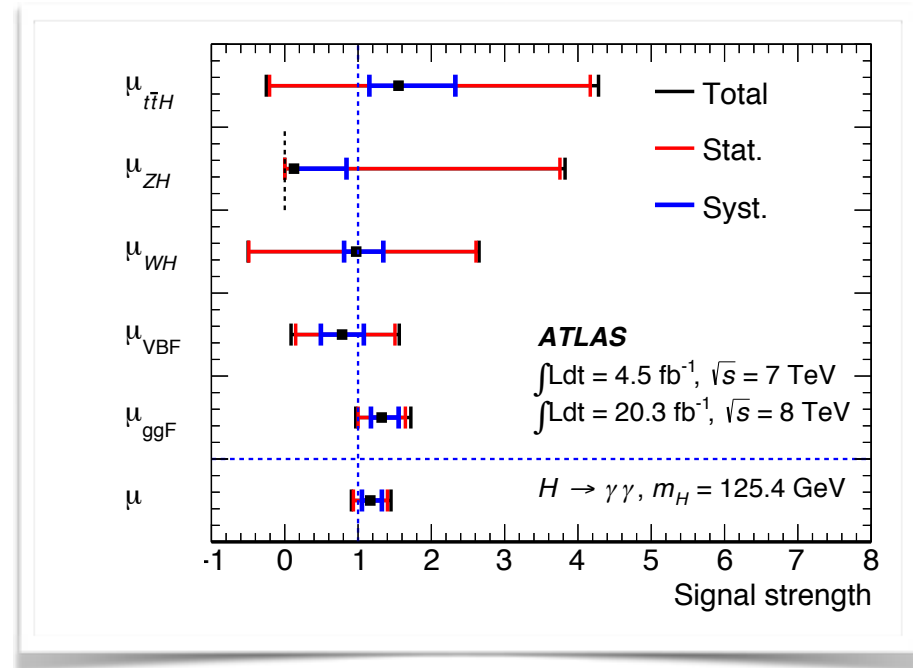
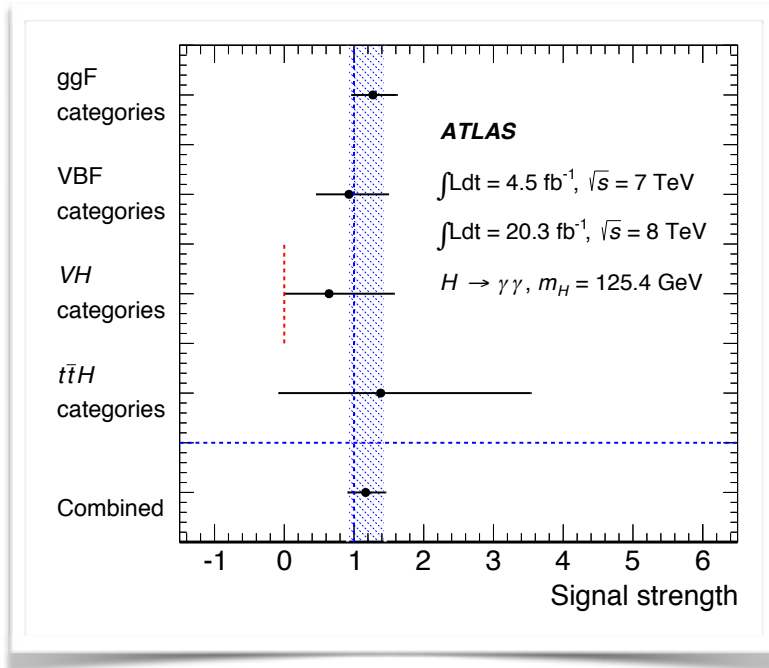
Improving Sensitivity: Categories



- For $\gamma\gamma$, categories are also used to improve sensitivity to the different production modes
- Define categories with higher or lower purity of a specific production mode



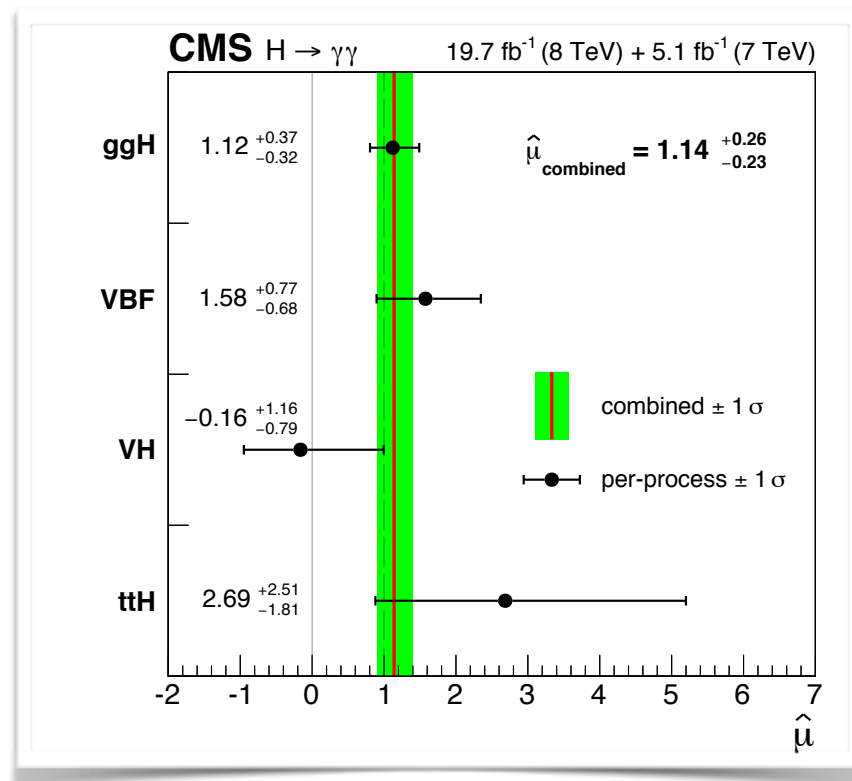
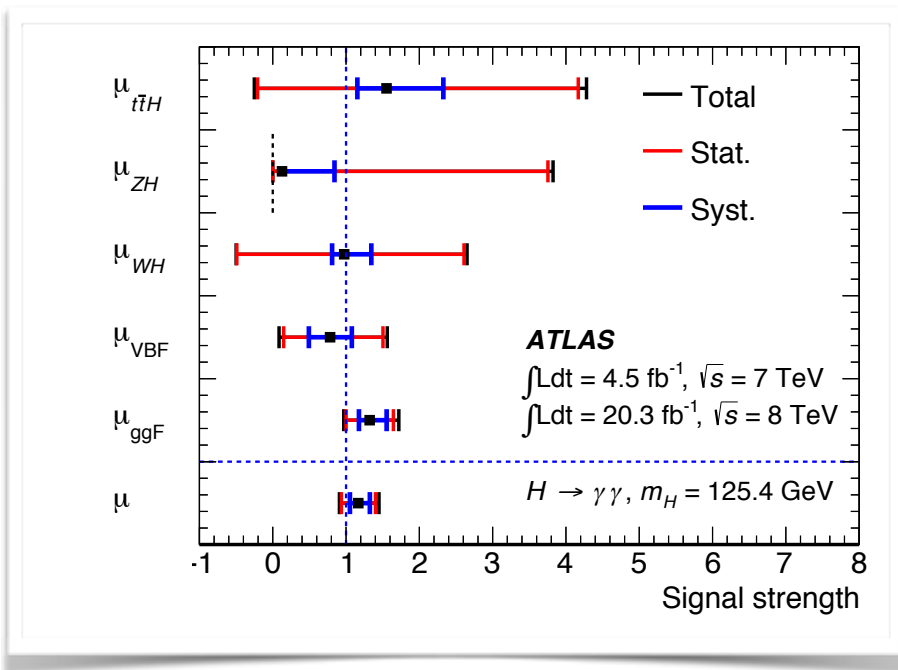
Why Categories? (2)



$$\begin{pmatrix} n_s^1 \\ \vdots \\ n_s^N \end{pmatrix} = \begin{pmatrix} \varepsilon_{ggF}^1 A_{ggF}^1 L^1 & \varepsilon_{VBF}^1 A_{VBF}^1 L^1 & \varepsilon_{VH}^1 A_{VH}^1 L^1 & \varepsilon_{t\bar{t}H}^1 A_{t\bar{t}H}^1 L^1 \\ \vdots & \vdots & \vdots & \vdots \\ \varepsilon_{ggF}^N A_{ggF}^N L^N & \varepsilon_{VBF}^N A_{VBF}^N L^N & \varepsilon_{VH}^N A_{VH}^N L^N & \varepsilon_{t\bar{t}H}^N A_{t\bar{t}H}^N L^N \end{pmatrix} \begin{pmatrix} \mu_{ggF} \sigma_{ggF}^{SM} \\ \mu_{VBF} \sigma_{VBF}^{SM} \\ \mu_{VH} \sigma_{VH}^{SM} \\ \mu_{t\bar{t}H} \sigma_{t\bar{t}H}^{SM} \end{pmatrix} B_{\gamma}^{SM} \quad \mu = \frac{\sigma}{\sigma_{tot}}$$

Extract measurements of all production modes

Final Run-1 $\gamma\gamma$ Results



$$\mu = 1.18 \pm 0.27$$

$$m_H = 125.4 \pm 0.27 \text{ GeV}$$

$$Z = 5.2(4.6)\sigma$$

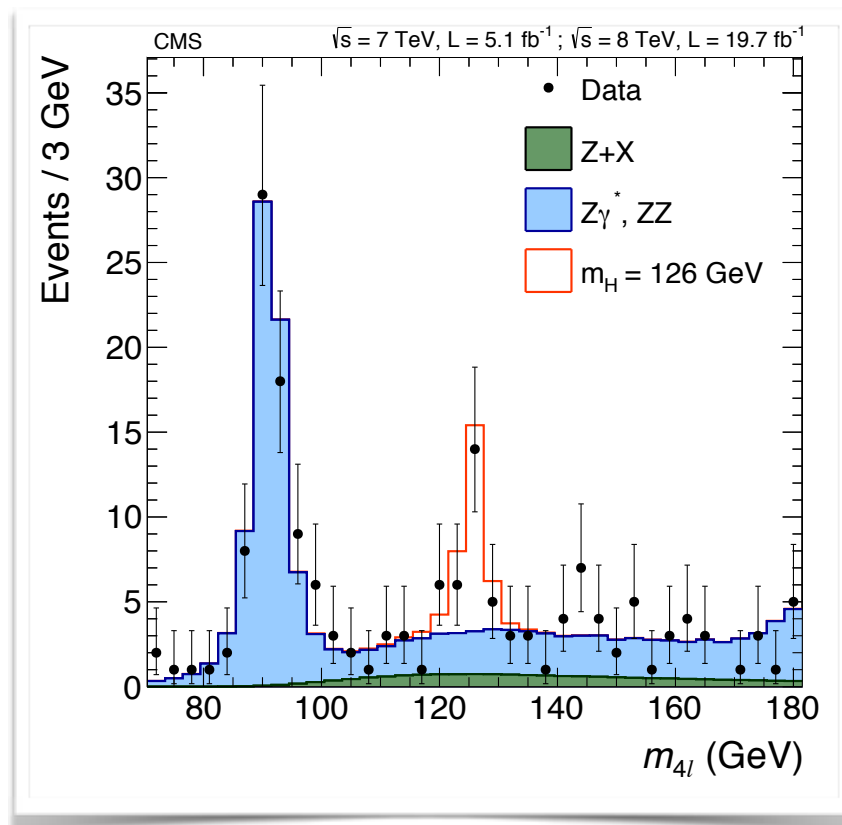
$$\mu = 1.14 \pm 0.26$$

$$m_H = 124.7 \pm 0.34 \text{ GeV}$$

$$Z = 5.7(5.2)\sigma$$

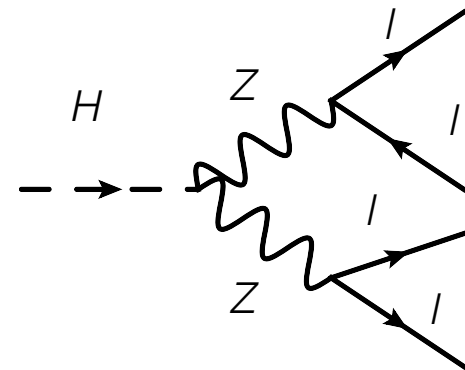
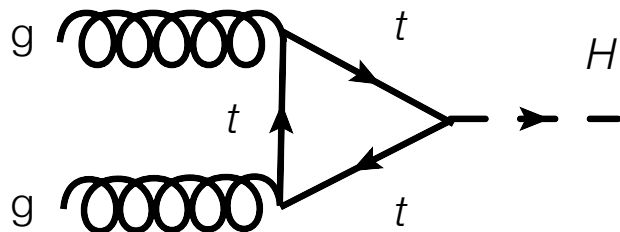
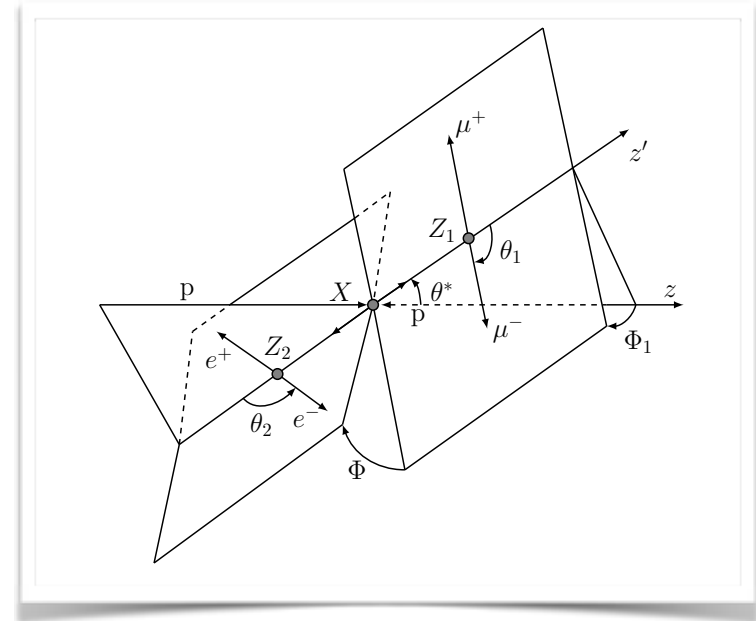
Higgs to 4 leptons ($H \rightarrow ZZ^* \rightarrow llll$)

- A good **discovery** final state
 - Low backgrounds
 - S/B: 1.5 - 10
 - Very good Higgs mass **resolution**
 - Requires good lepton reconstruction efficiencies
 - Muon spectrometers designed specifically for this channel
 - Clear and robust signal of Higgs coupling to weak bosons
- Select 4 reconstructed leptons
 - $4e, 4\mu, 2e2\mu$

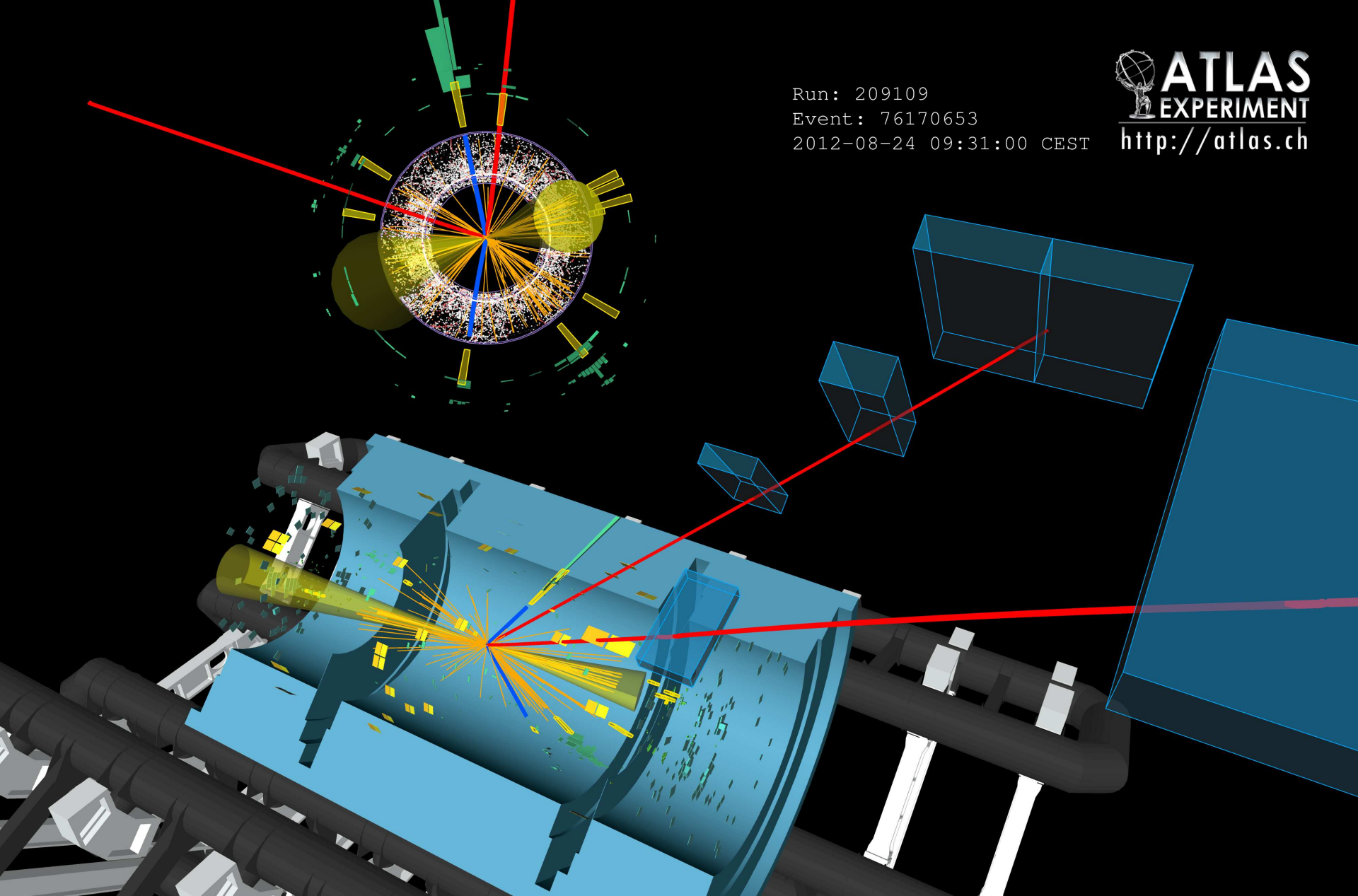


Production and Decay

- As for $\gamma\gamma$, production through top loop
- Decay depends only on coupling to Z boson
- Small branching fraction to the 4-lepton final state (2.6%)
- Improve sensitivity by using full event information (e.g. in MVA)
 - 2 production and 3 decay angles
 - Z_1 and Z_2 masses

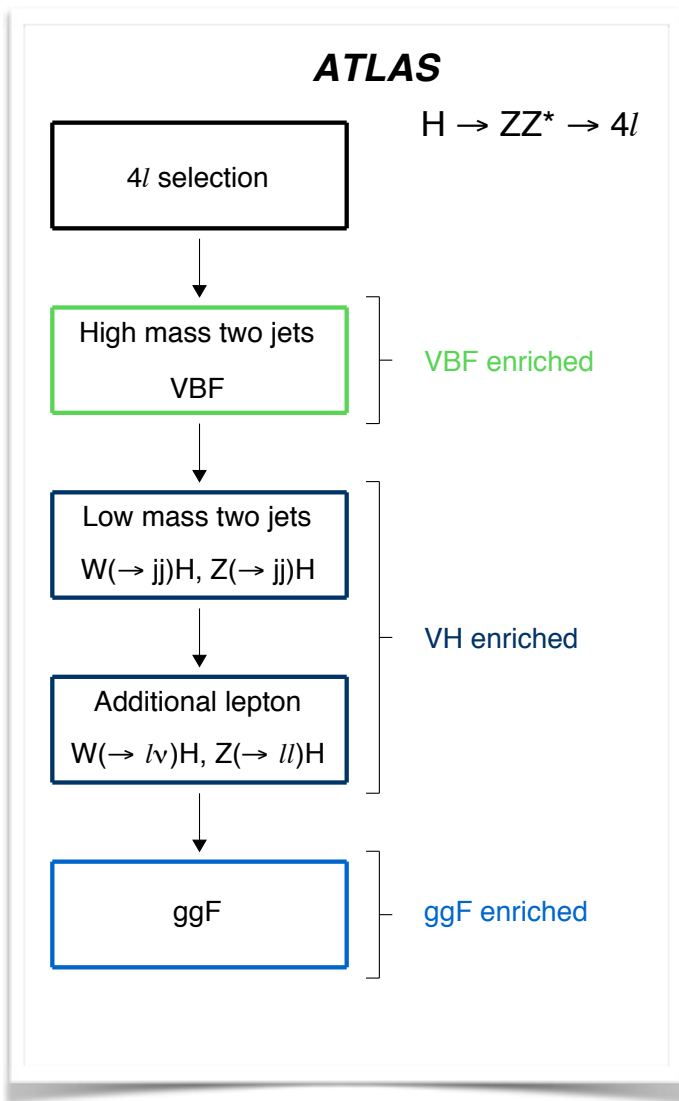


Run: 209109
Event: 76170653
2012-08-24 09:31:00 CEST

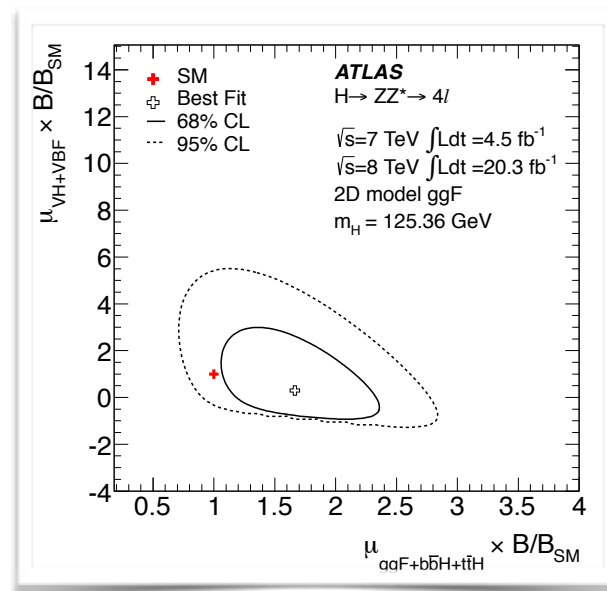
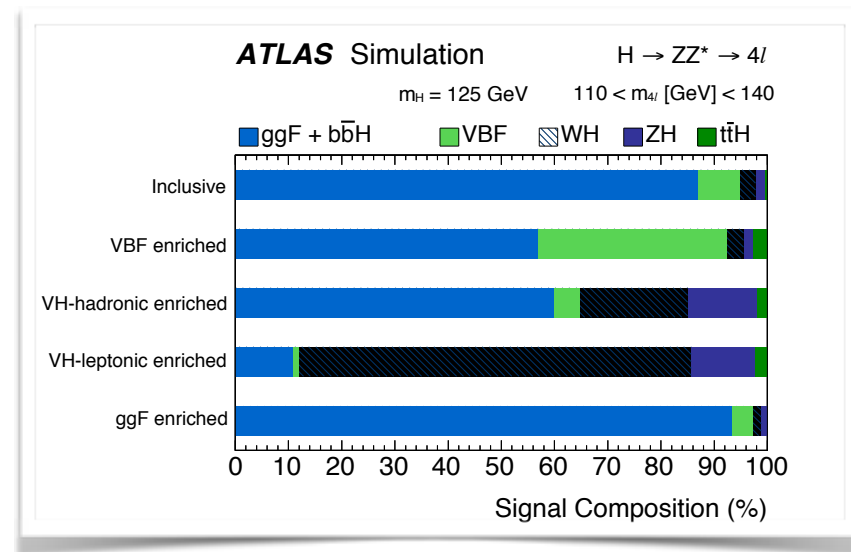


EventNumber : 76170653 RunNumber : 209109 $m_{4\ell}=123.4$ GeV. The BDT_{VBF} value is 0.7. Six jets in total, and the two leading jets have $p_T = 180$ and 150 GeV, $\Delta\eta_{jj} = 3.4$, and $p_{Tjj} = 200$ GeV. The missing $E_T = 40$ GeV.

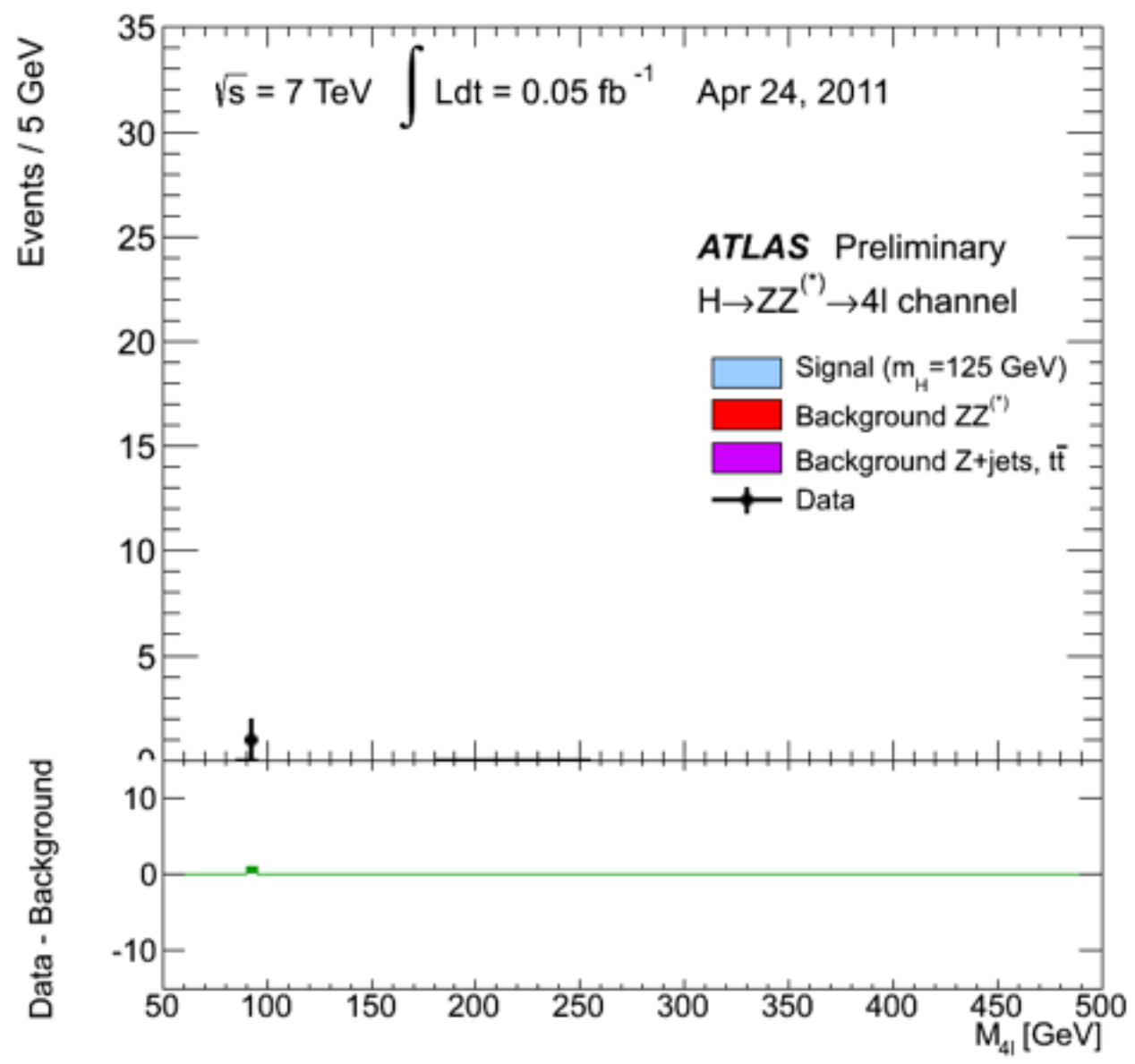
Categories



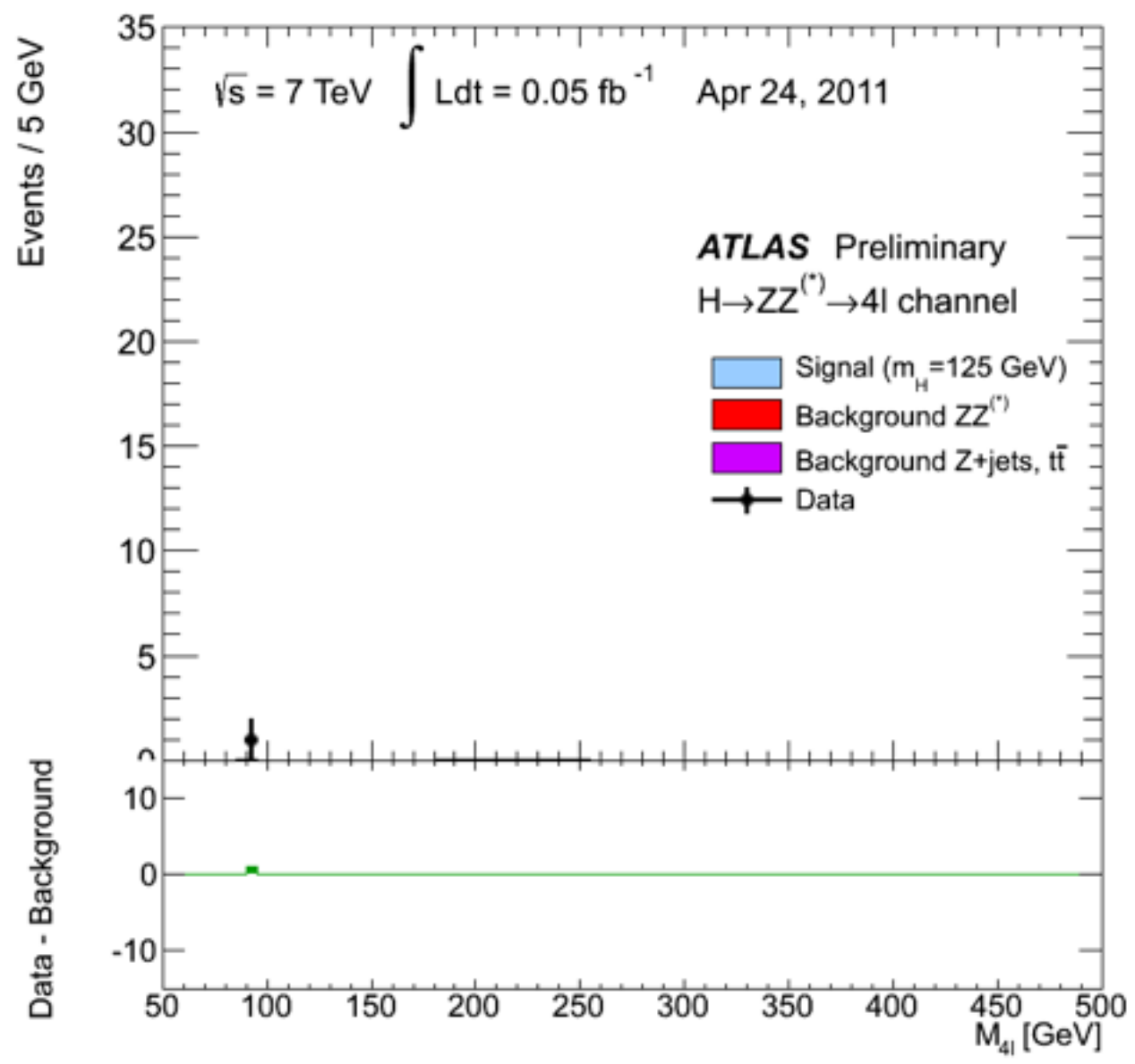
Limited by statistics so will improve quickly with more data



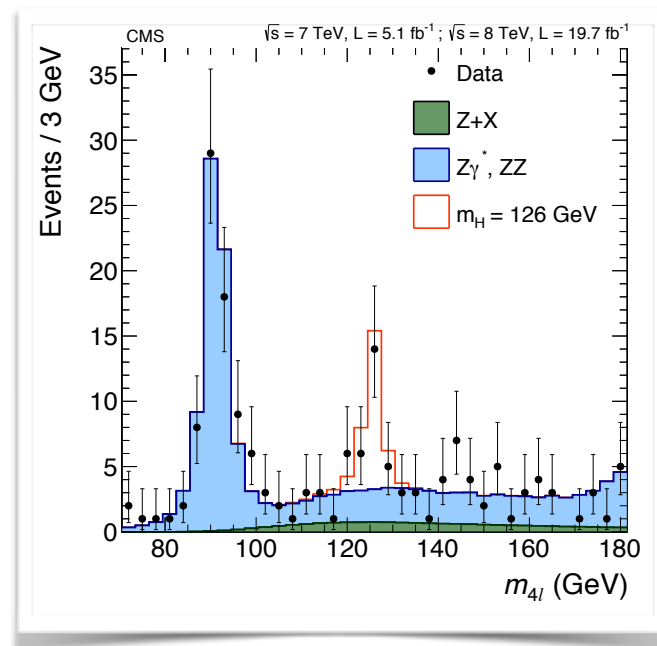
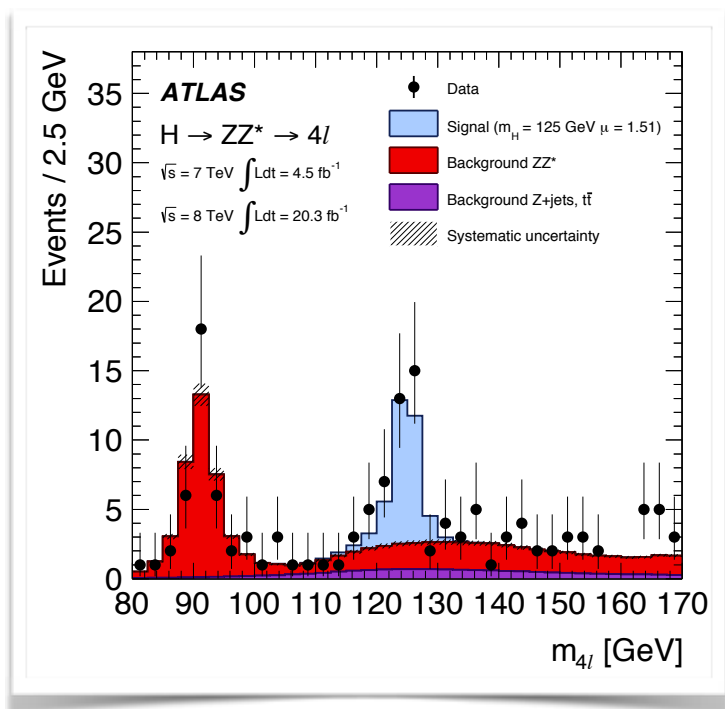
Higgs to 4 leptons ($H \rightarrow ZZ^* \rightarrow llll$)



Higgs to 4 leptons ($H \rightarrow ZZ^* \rightarrow llll$)



Final Run-1 ZZ Results



$$\mu = 1.44^{+0.40}_{-0.33}$$

$$m_H = 125.4 \pm 0.27 \text{ GeV}$$

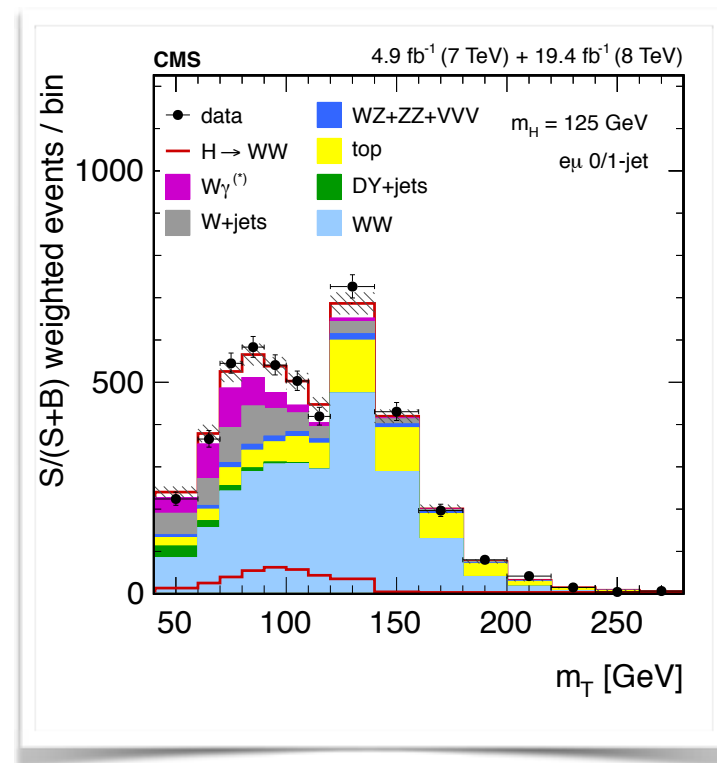
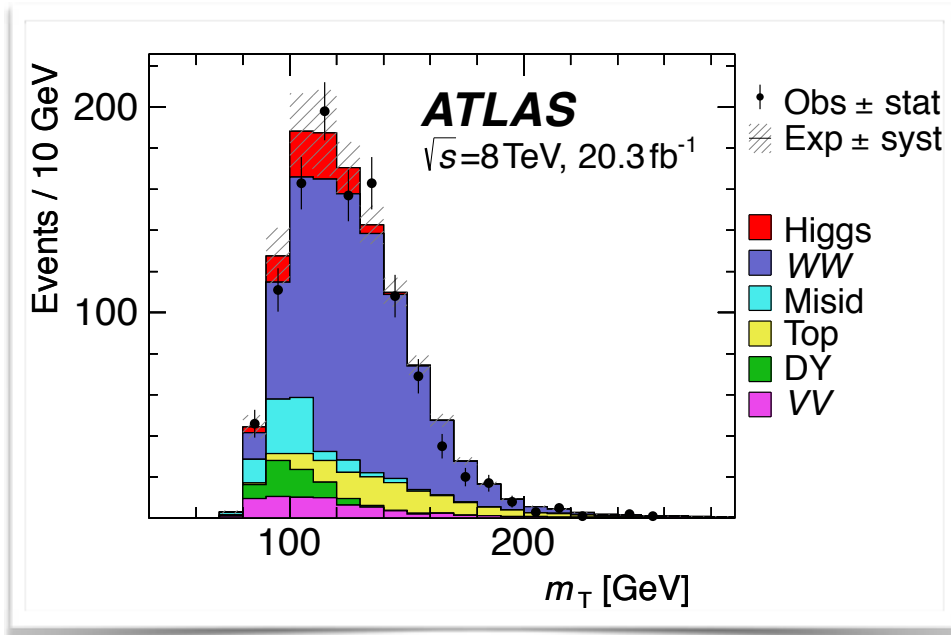
$$Z = 8.1(6.2)\sigma$$

$$\mu = 0.93 \pm 0.29$$

$$m_H = 125.6 \pm 0.45 \text{ GeV}$$

$$Z = 6.8(6.7)\sigma$$

H → WW

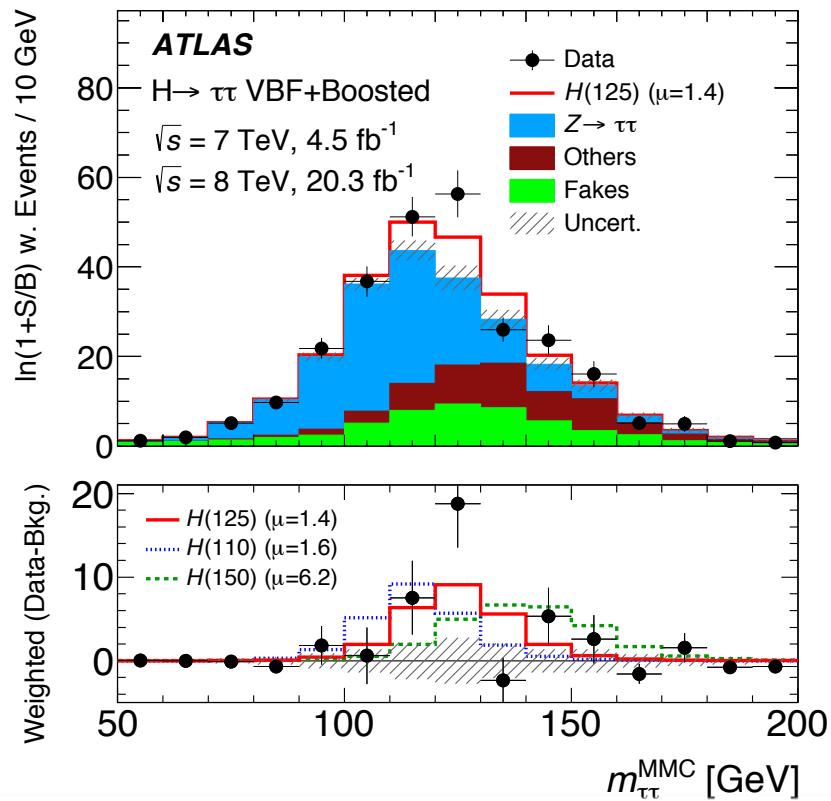


$$\mu = 1.22^{+0.23}_{-0.21}$$

$$Z = 6.8(5.8)\sigma$$

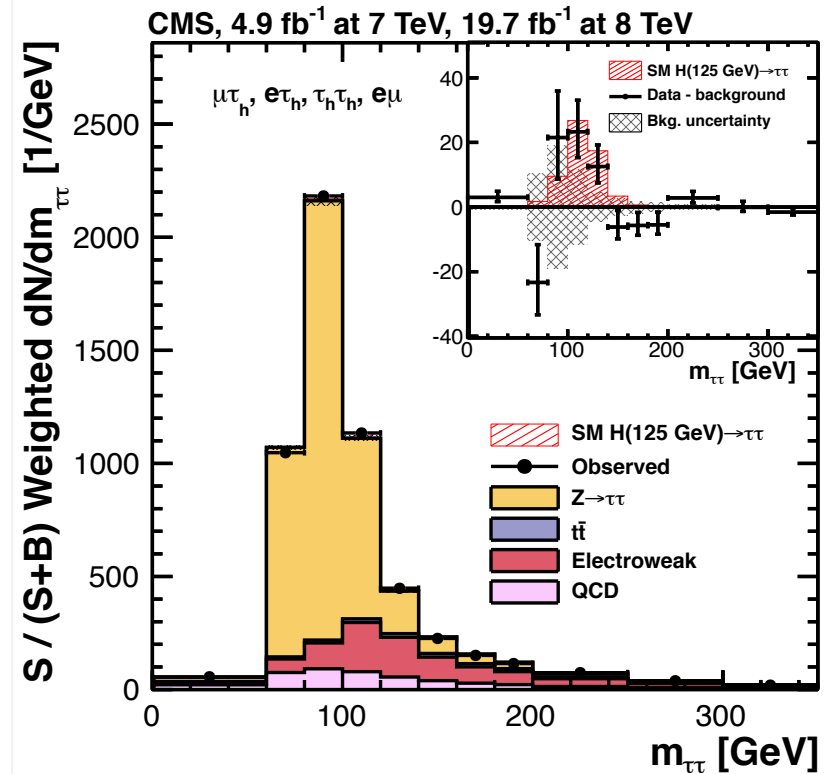
$$\mu = 0.90^{+0.23}_{-0.21}$$

$$Z = 4.8(5.6)\sigma$$

$H \rightarrow \tau\tau$ 

$$\mu = 1.41^{+0.40}_{-0.36}$$

$$Z = 4.4(3.3)\sigma$$

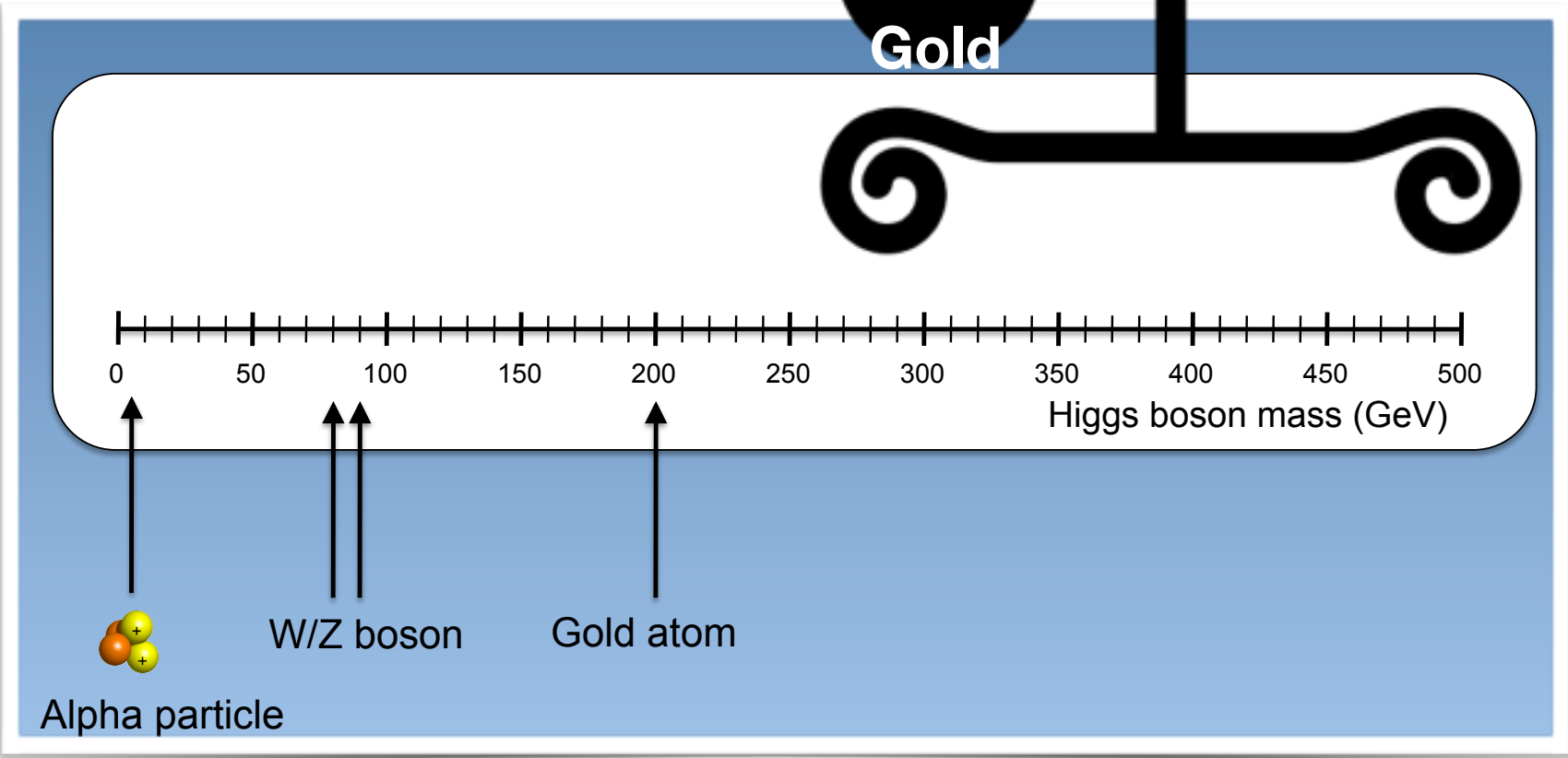
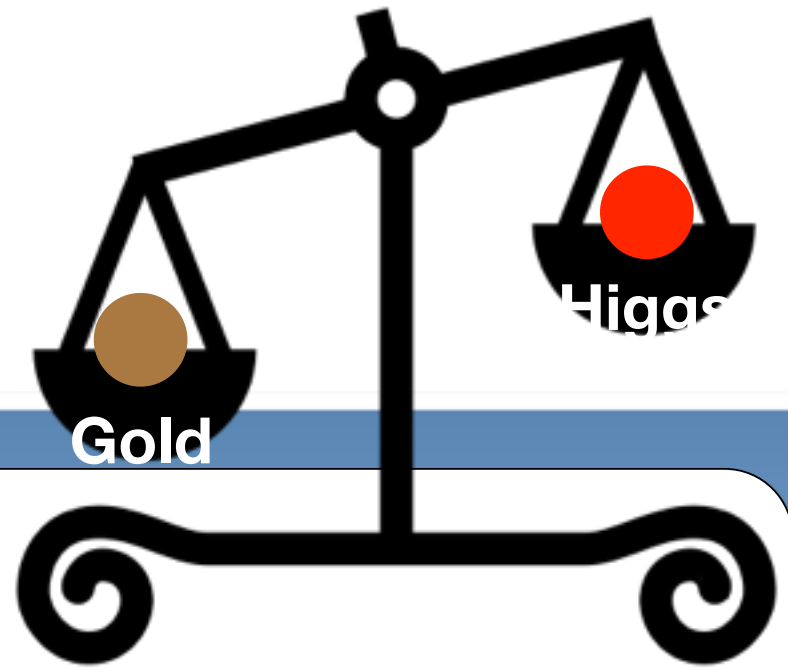


$$\mu = 0.88^{+0.30}_{-0.28}$$

$$Z = 3.4(3.7)\sigma$$

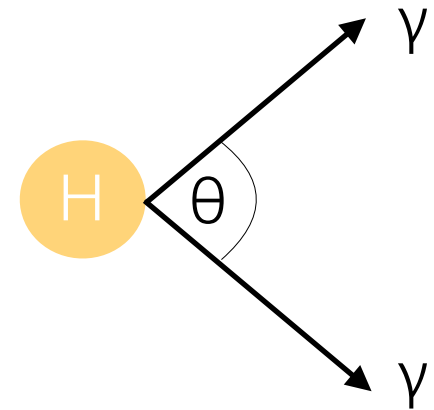
Mass

Higgs boson is almost as heavy as a gold atom



Measuring Particle Masses: An Example

- Higgs boson decays to photon 1 and photon 2
- Photons are reconstructed in the calorimeter
- Use the momentum (energy) of these photons to calculate the Higgs boson mass
- $(E, p_x, p_y, p_z) =$
 - $(|P_1|, p_{x1}, p_{y1}, p_{z1}) + (|P_2|, p_{x2}, p_{y2}, p_{z2})$
 - $= (|P_1| + |P_2|, p_{x1} + p_{x2}, p_{y1} + p_{y2}, p_{z1} + p_{z2})$
- Mass of Higgs boson: $m^2 = E^2 - p^2$
- Or $m^2 = 2|P^1||P^2| \cdot (1 - \cos\theta)$



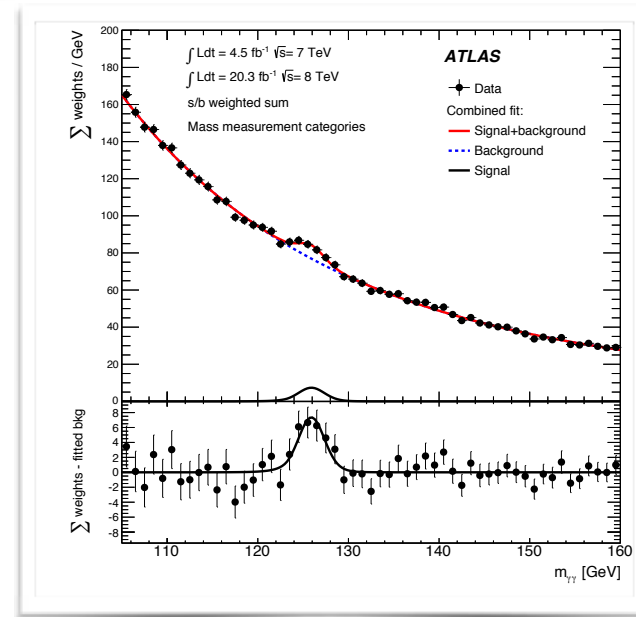
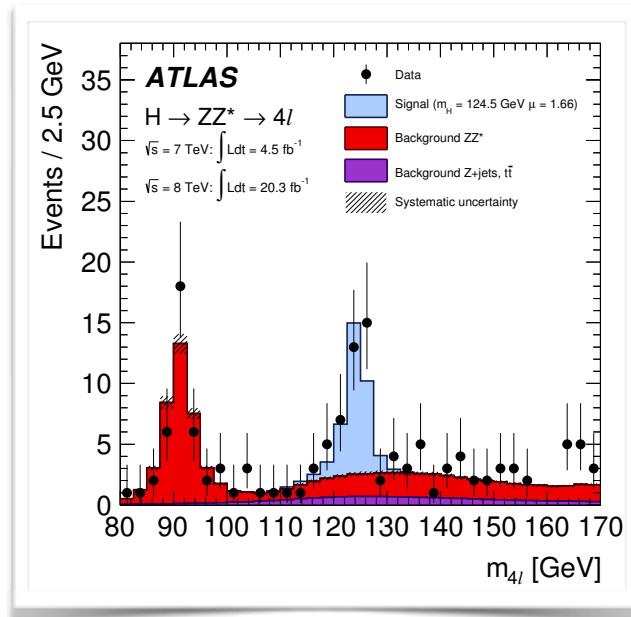
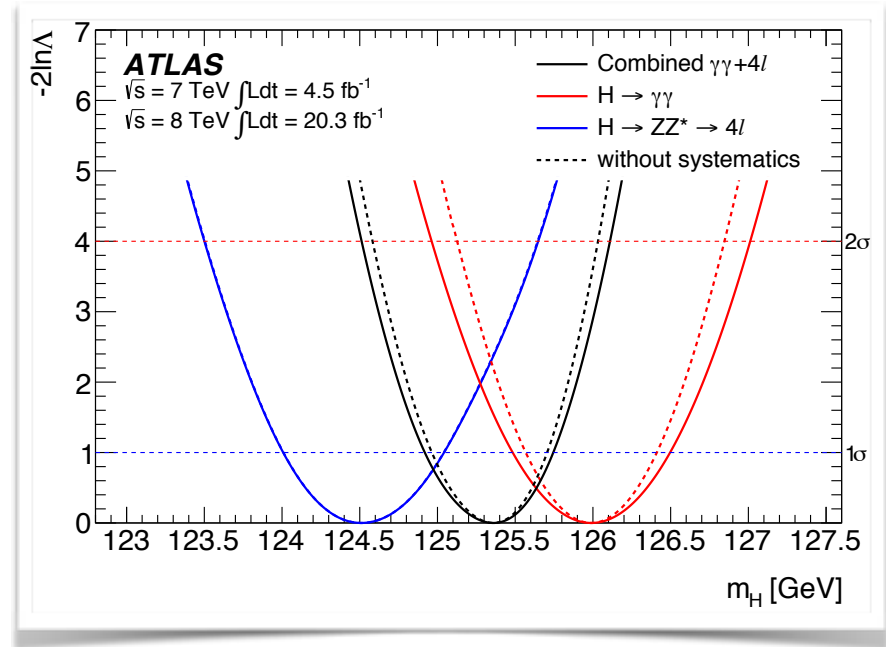
ATLAS Mass Measurement

Phys. Rev. D. 90, 052004 (2014)

ZZ: 124.51 ± 0.52 (stat) ± 0.06 (syst) GeV

$\gamma\gamma$: 125.98 ± 0.42 (stat) ± 0.28 (syst) GeV

Combined: 125.36 ± 0.37 (stat) ± 0.18 (syst) GeV



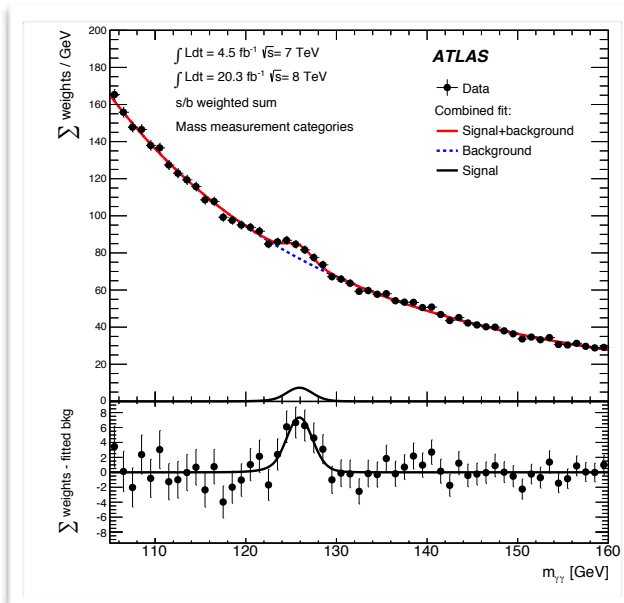
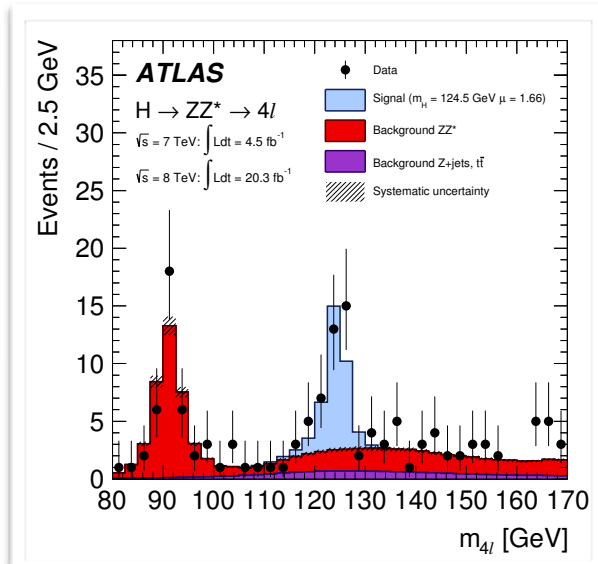
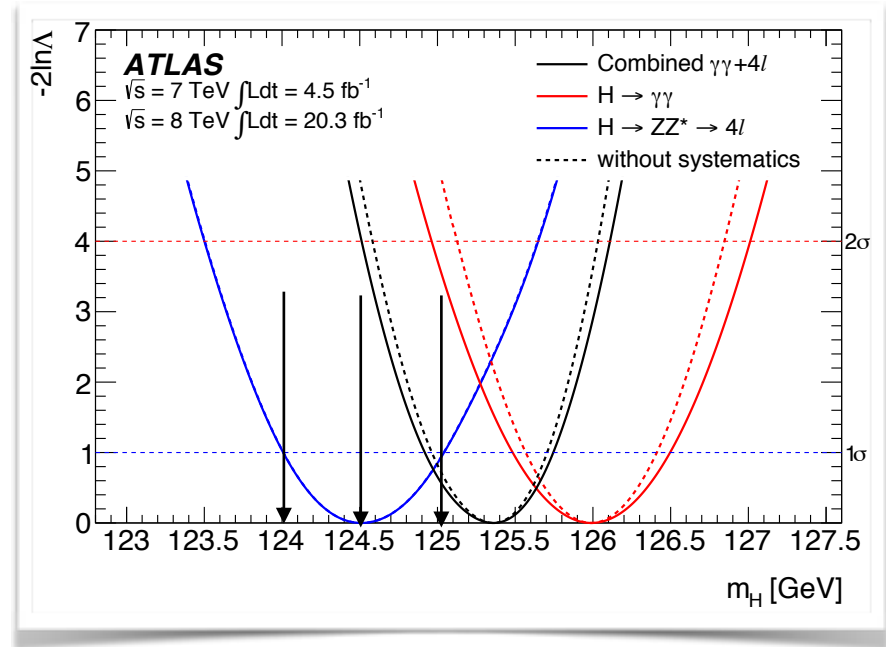
ATLAS Mass Measurement

Phys. Rev. D. 90, 052004 (2014)

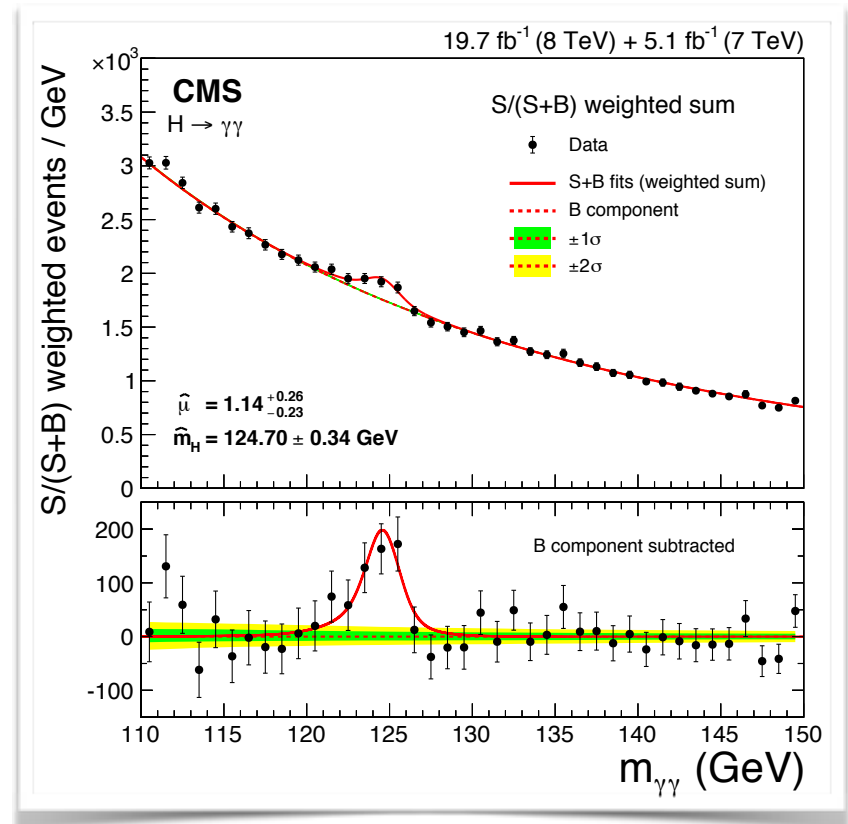
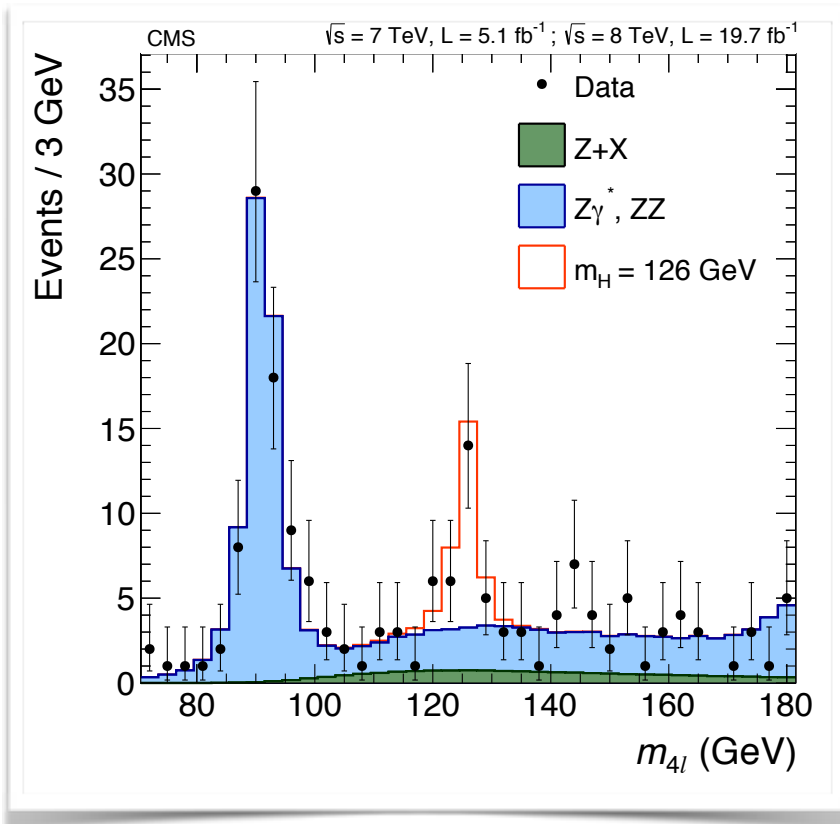
ZZ: 124.51 ± 0.52 (stat) ± 0.06 (syst) GeV

$\gamma\gamma$: 125.98 ± 0.42 (stat) ± 0.28 (syst) GeV

Combined: 125.36 ± 0.37 (stat) ± 0.18 (syst) GeV



CMS Mass Measurement

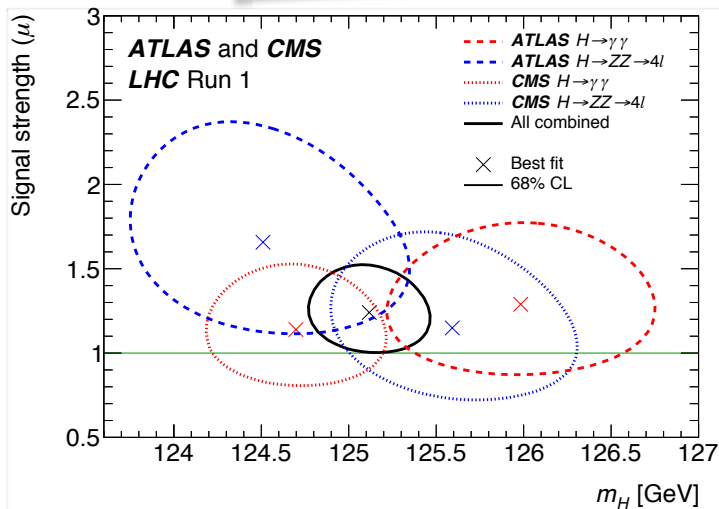
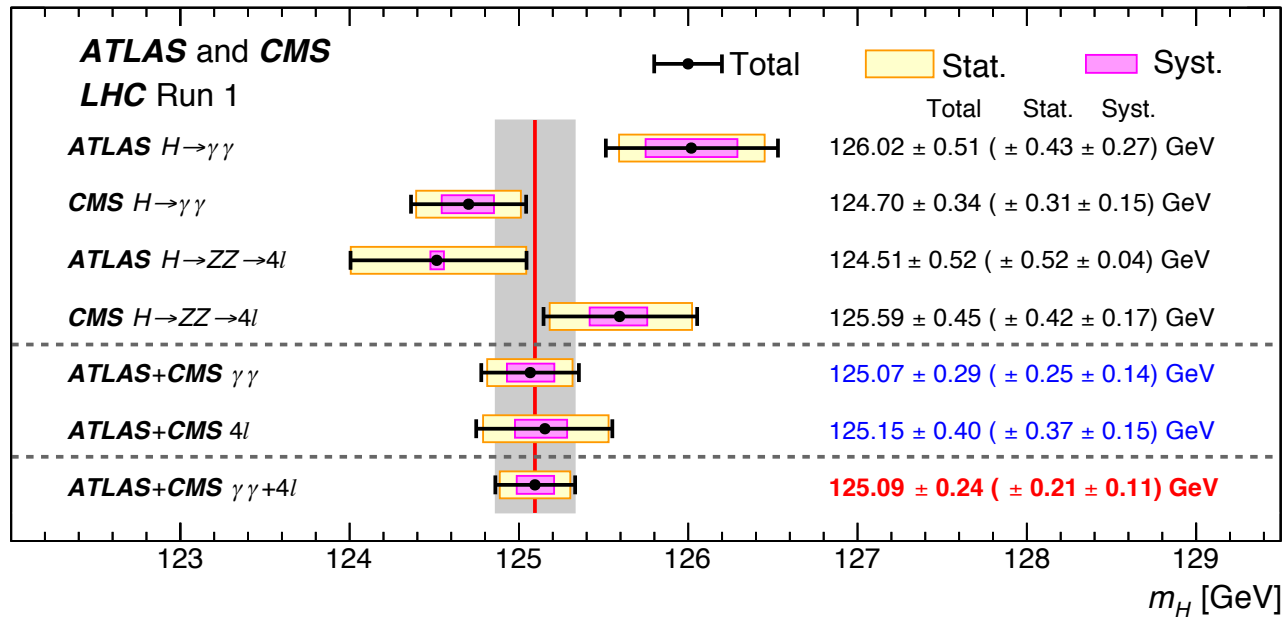


ZZ: $125.59 \pm 0.42 \text{ (stat)} \pm 0.17 \text{ (syst)} \text{ GeV}$

$\gamma\gamma$: $124.70 \pm 0.31 \text{ (stat)} \pm 0.15 \text{ (syst)} \text{ GeV}$

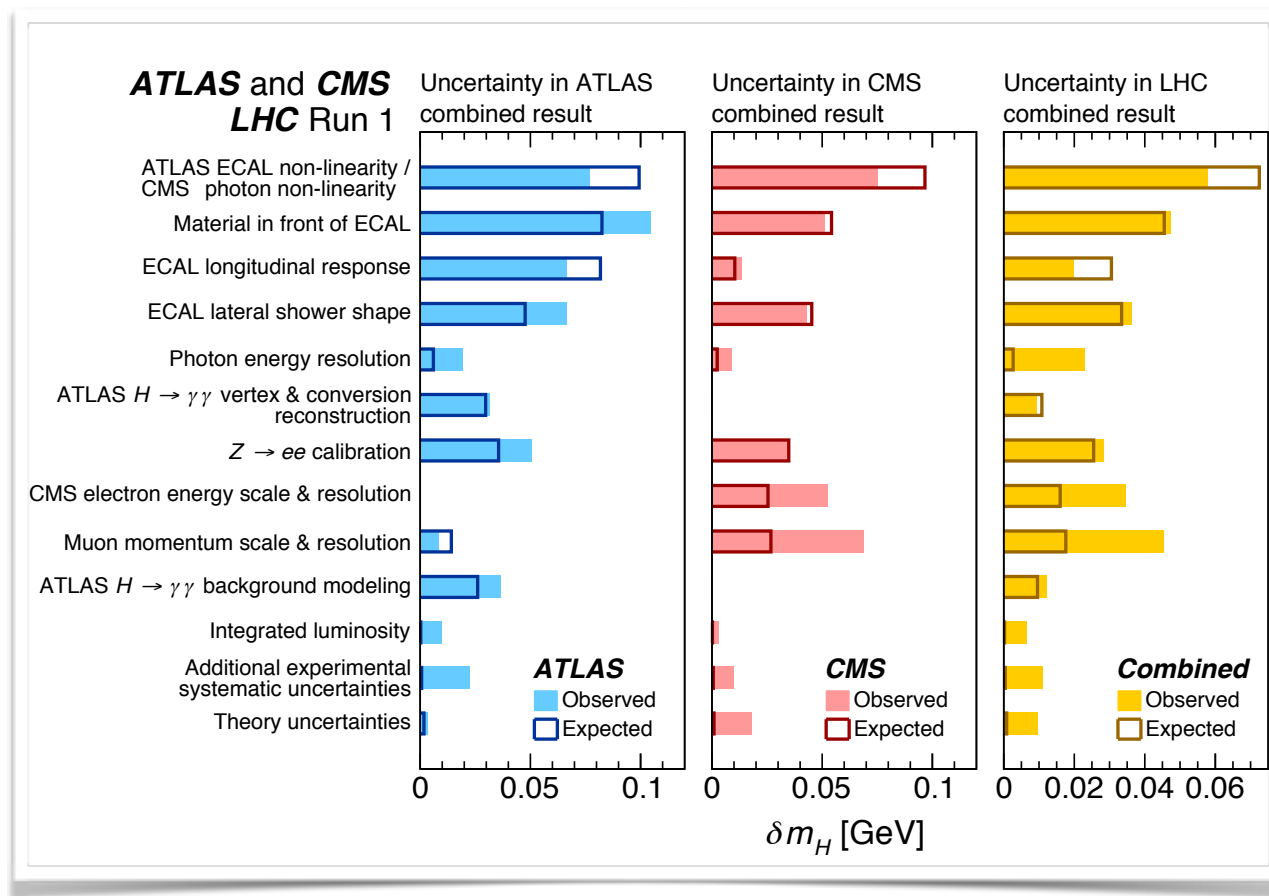
Combined Higgs Mass Measurement

HIGG-2014-14



- Measurement dominated by statistics
- The compatibility of the four measurements is to within 10%
- Tension between ATLAS $4l$ and $\gamma\gamma$ $\sim 2\sigma$

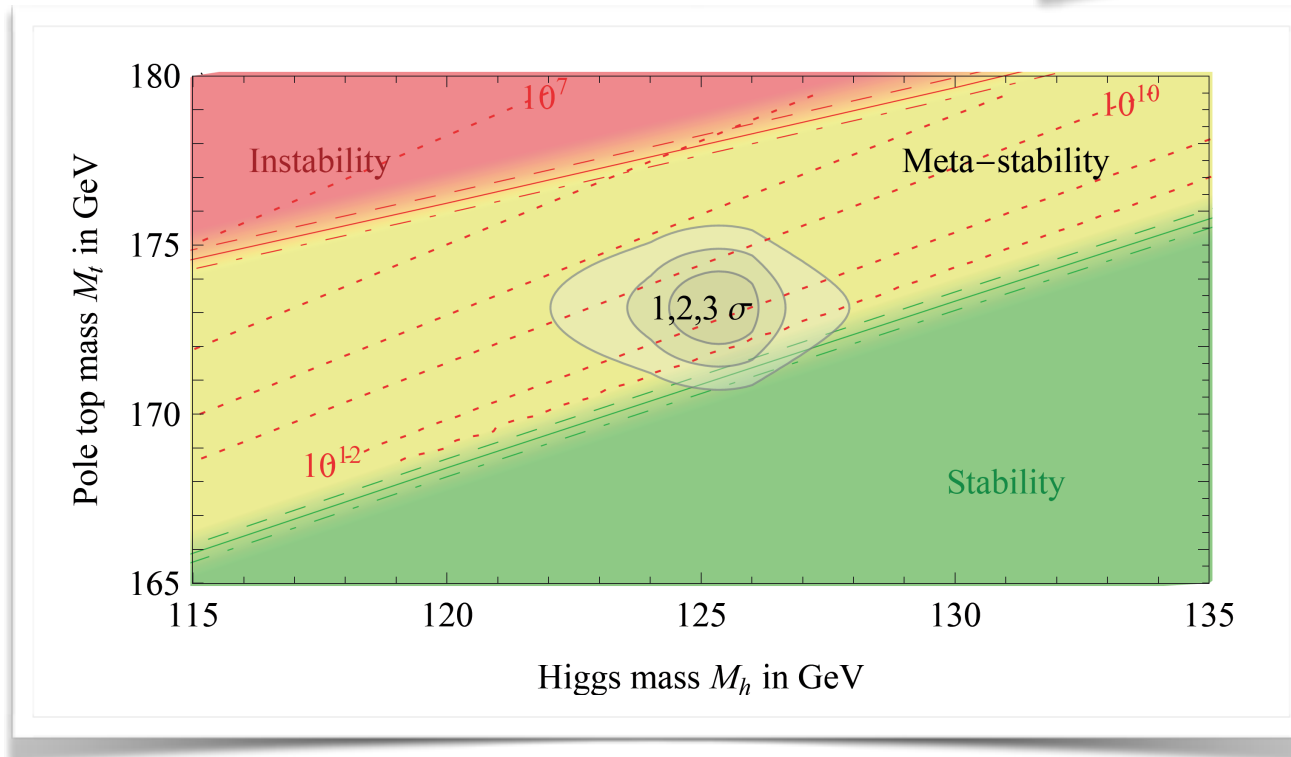
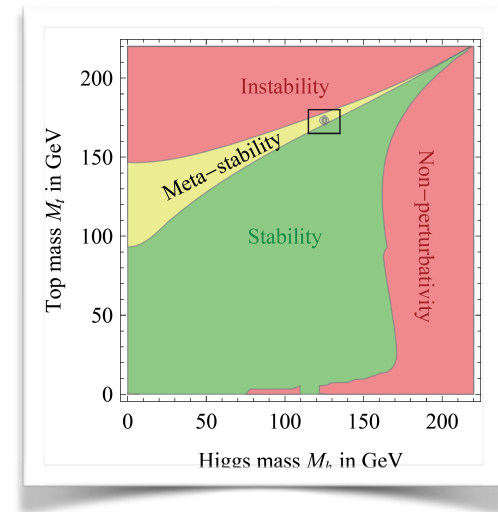
Uncertainties on the Mass Measurement



Largest systematic uncertainties are those from the calibration of the calorimeter

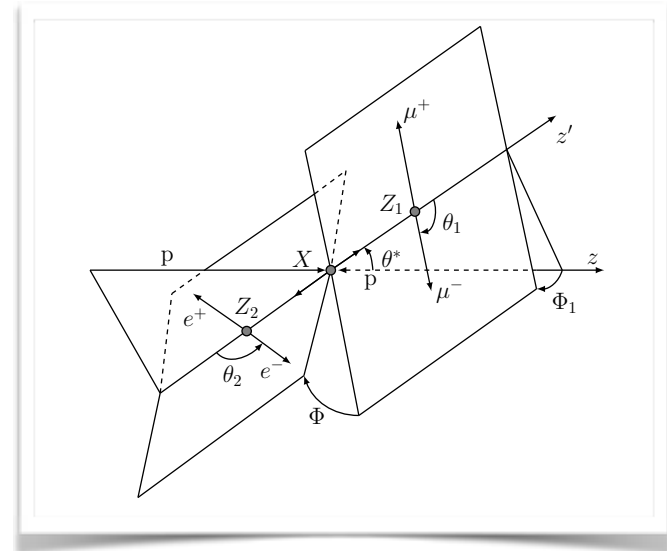
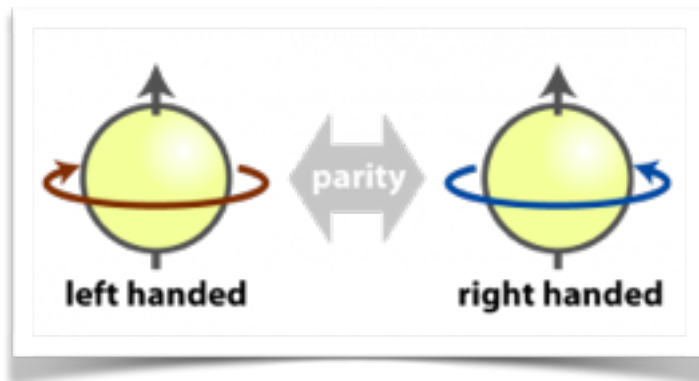
Mass Measurement Implications

- $m_H = 125 \text{ GeV} \rightarrow$ our universe lies on the boundary between instability and stability
- No need to panic: metastability means that the universe is unlikely to end tomorrow
 - But intriguing, nonetheless

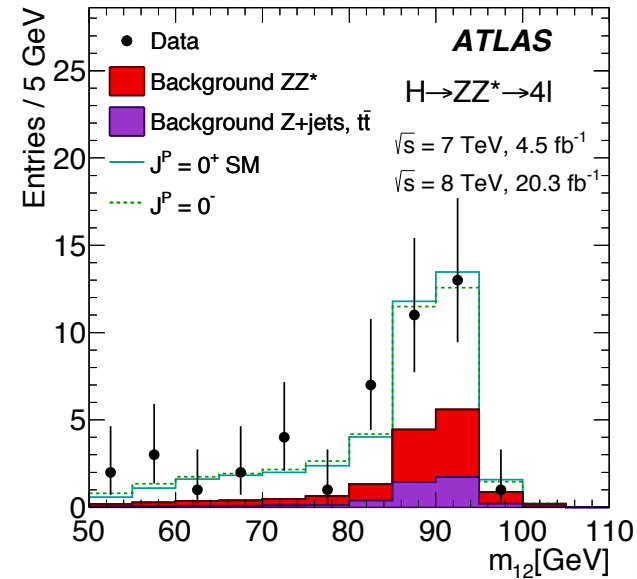
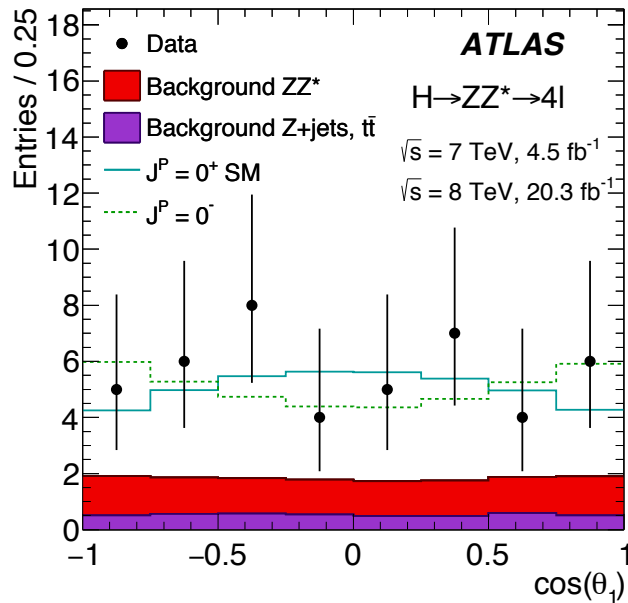


Spin/Parity

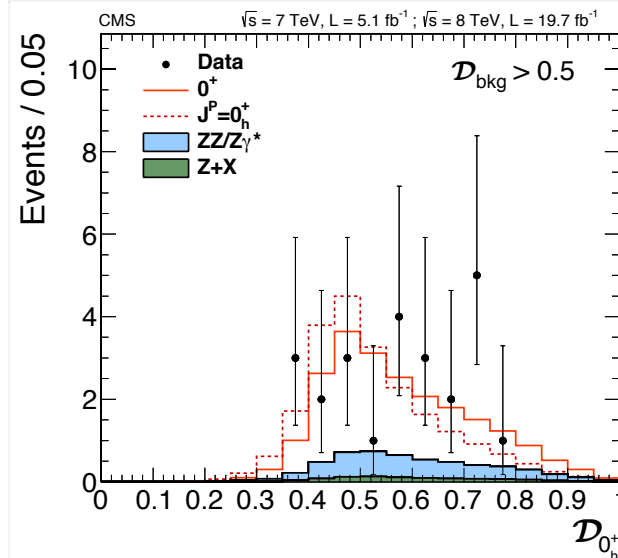
- Only elementary particle with spin-0
- Spin and parity determine angular distributions of decay products
 - Use $\gamma\gamma$, ZZ and WW
- Don't forget, though, that the $\gamma\gamma$ observation implies
 - does not originate from spin 1 : Landau-Yang theorem
 - charge conjugation is +1 (assuming C and P separately conserved)
 - WW/ZZ channels disfavour CP odd hypothesis (can occur through loops)



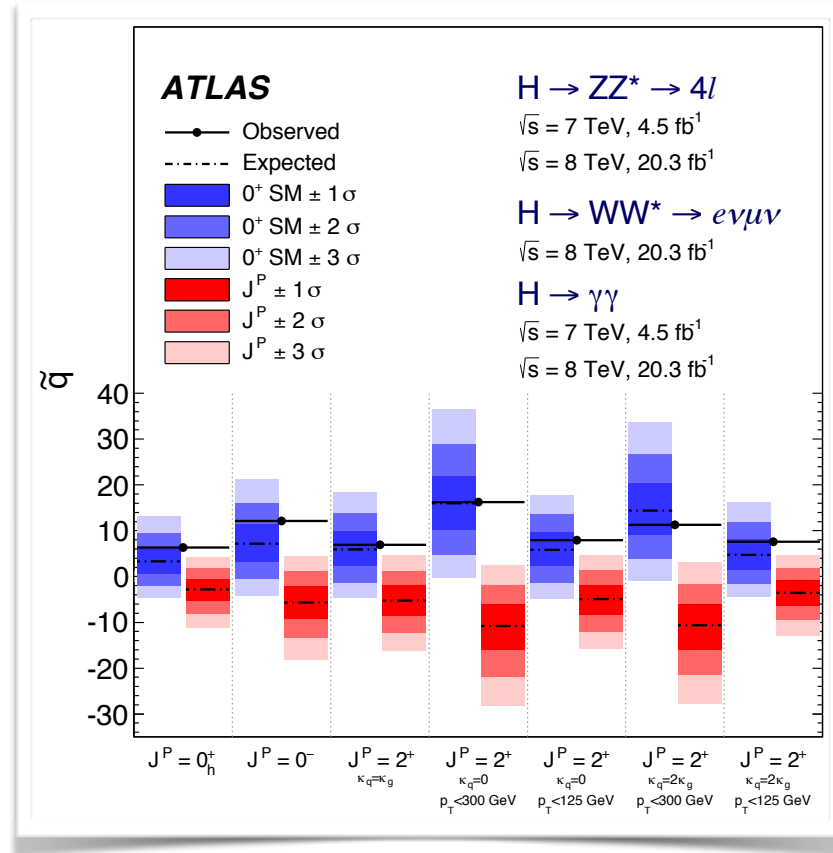
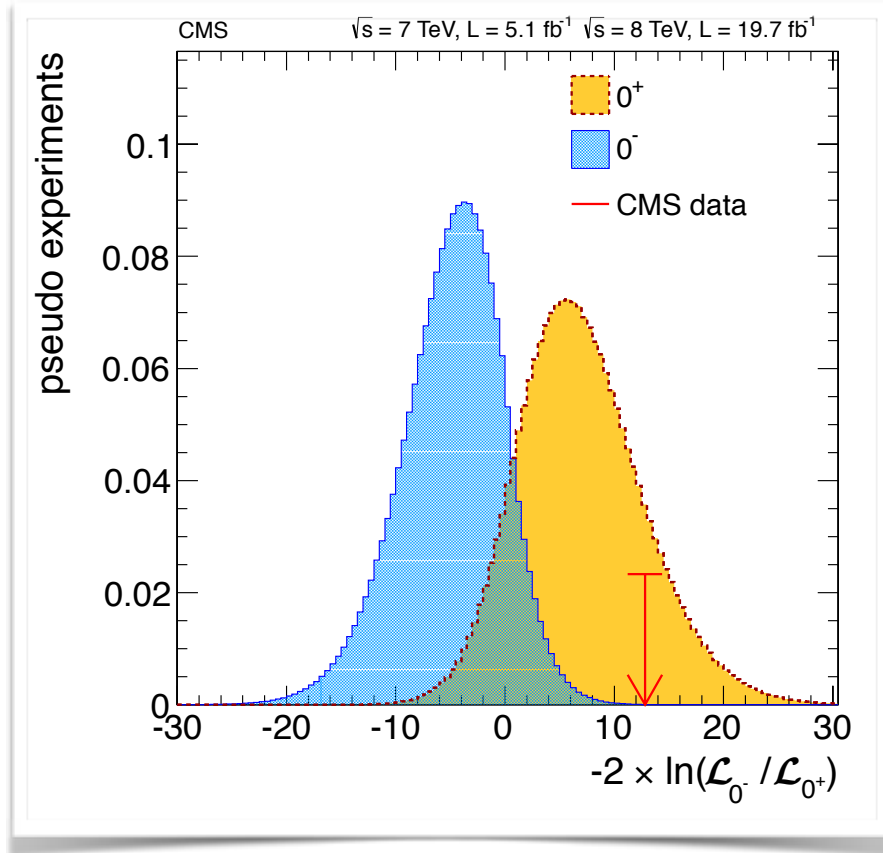
Variables sensitive to spin



- Study a number of angular variables sensitive to spin and parity
- Combine into a single discriminant using an MVA or theory-based matrix element technique

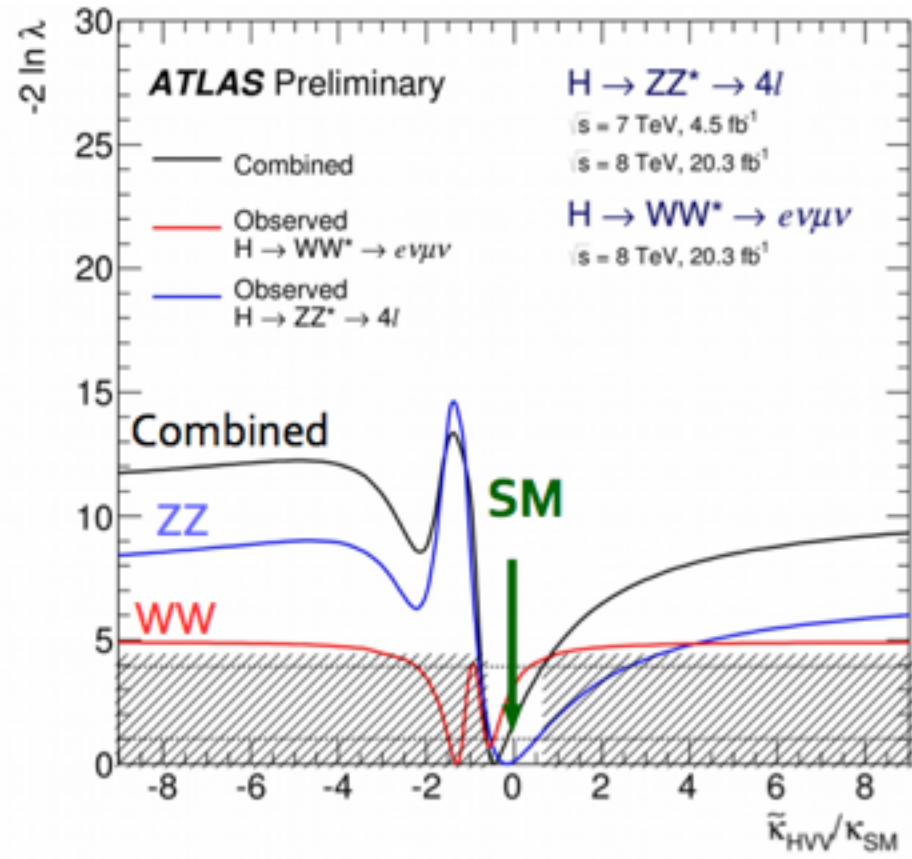
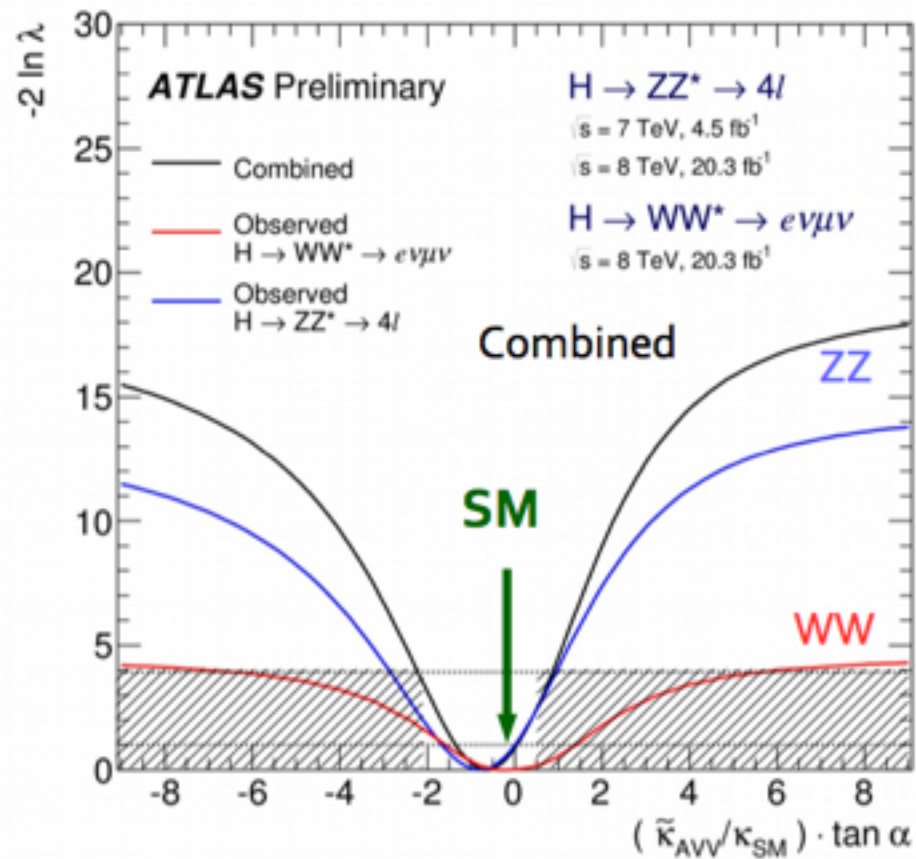


Spin/Parity Results



Strong evidence that the Higgs is 0^+ as predicted by the Standard Model

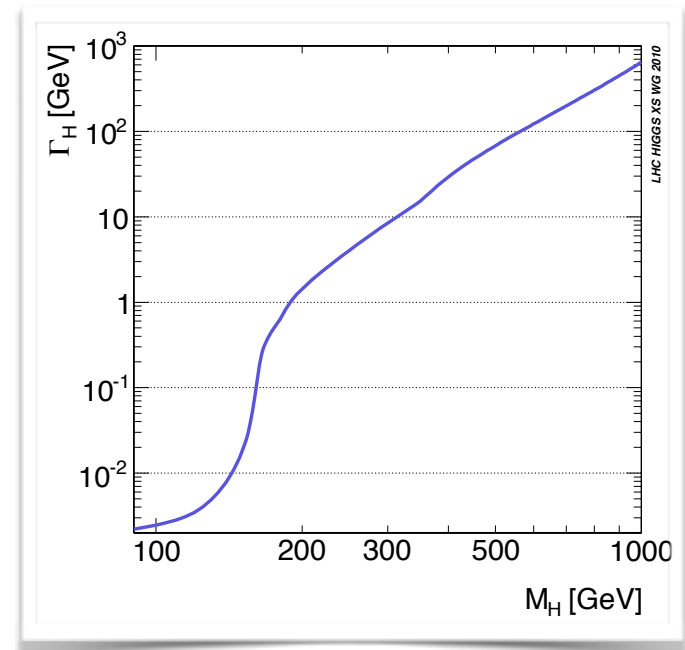
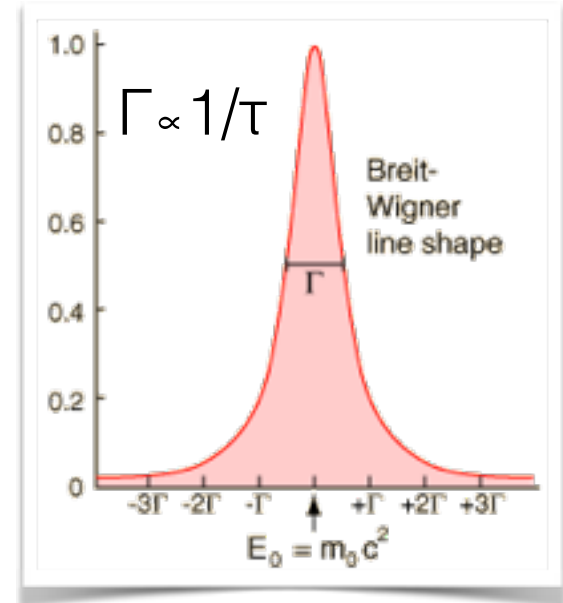
CP Mixing



Both ATLAS and CMS find that the observed Higgs boson is compatible with a standard CP-even

Width

- As an highly unstable elementary particle, the **lifetime** of the Higgs is **very short**
- For $m_H = 125 \text{ GeV}$
 - $\Gamma = 4.07 \times 10^{-3} \text{ GeV}$
- Direct experimental measurements probe widths **3 orders of magnitude larger** $\sim 1.6 \text{ GeV}$ (ATLAS, ZZ)
- Thought to be **impossible to measure the width** at a hadron collider



Off-shell Higgs Production

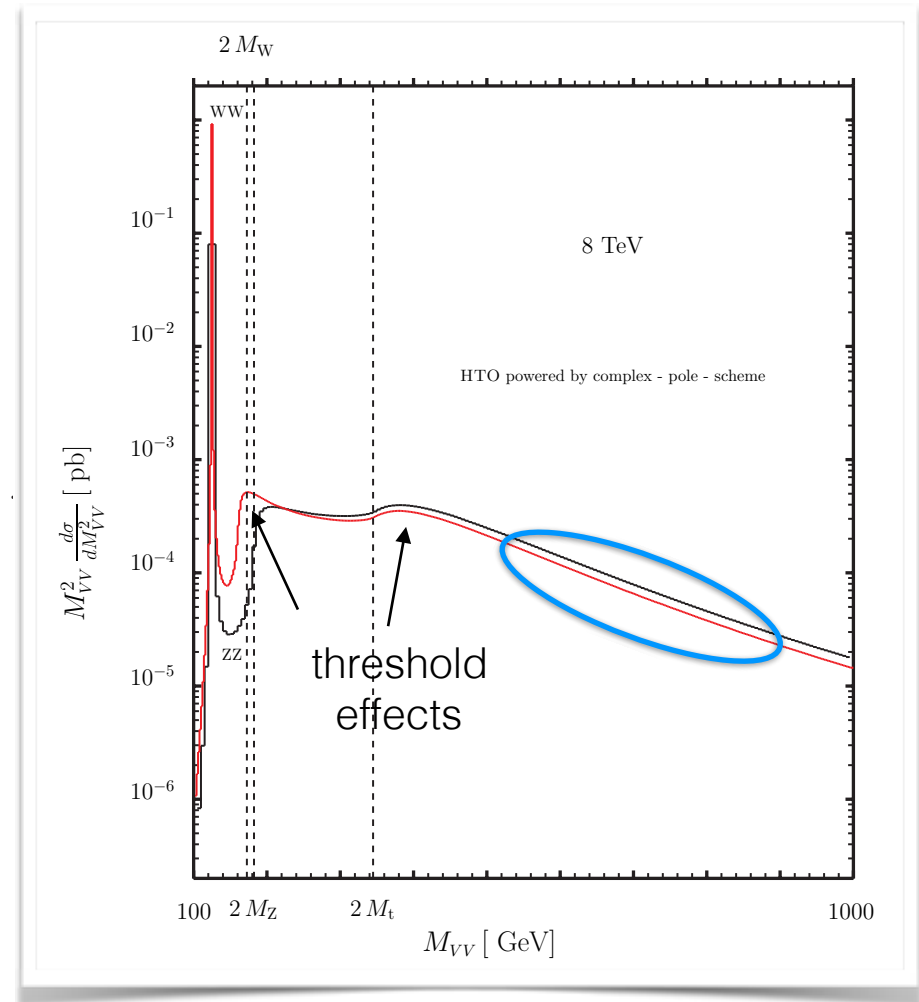
N. Kauer and G. Passarino, JHEP 08 (2012) 116

- A paper from Kauer and Passarino in 2012 pointed out a peculiar cancellation between the Breit-Wigner trend and the width as a function of m_{VV} enhances the cross-section at high mass

$$\left(\frac{d\sigma}{dM_{VV}}\right)_{ZWA} = \sigma_{H,ZWA} \frac{M_H \Gamma_H}{\pi} \frac{2M_{VV}}{(M_{VV}^2 - M_H^2)^2 + (M_H \Gamma_H)^2}$$

- For ZZ, ~7.6% of the total cross-section is at high mass

	Tot[pb]	$M_{ZZ} > 2 M_Z$ [pb]	R[%]
$gg \rightarrow H \rightarrow \text{all}$	19.146	0.1525	0.8
$gg \rightarrow H \rightarrow ZZ$	0.5462	0.0416	7.6



Measuring the Width

- This can be used to set a constraint on the Higgs width as follows

$$\sigma_{gg \rightarrow H \rightarrow ZZ}^{\text{on-peak}} = \frac{\kappa_g^2 \kappa_Z^2}{r} (\sigma \cdot \text{BR})_{\text{SM}} \equiv \mu (\sigma \cdot \text{BR})_{\text{SM}}$$

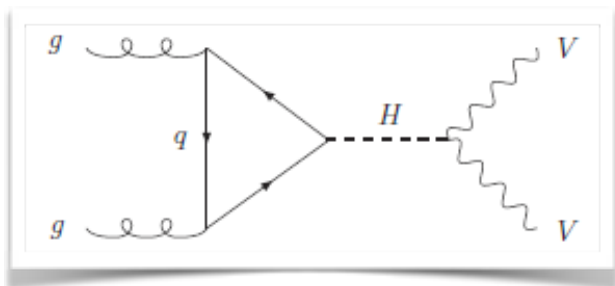
$$\frac{d\sigma_{gg \rightarrow H \rightarrow ZZ}^{\text{off-peak}}}{dm_{ZZ}} = \kappa_g^2 \kappa_Z^2 \frac{d\sigma_{gg \rightarrow H \rightarrow ZZ}^{\text{off-peak, SM}}}{dm_{ZZ}} = \mu r \frac{d\sigma_{gg \rightarrow H \rightarrow ZZ}^{\text{off-peak, SM}}}{dm_{ZZ}}$$

$$\kappa_g = g_{ggH} / g_{ggH}^{\text{SM}}$$

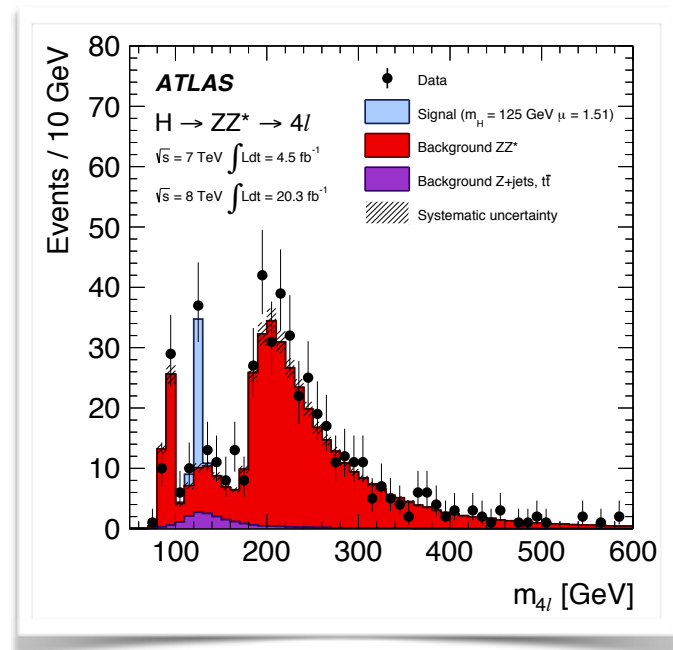
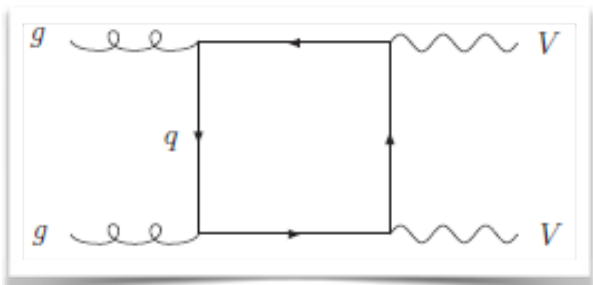
$$\kappa_Z = g_{HZZ} / g_{HZZ}^{\text{SM}}$$

$$r = \Gamma_H / \Gamma_H^{\text{SM}}$$

- Determine r by measuring ratio of off-peak to on-peak cross-section

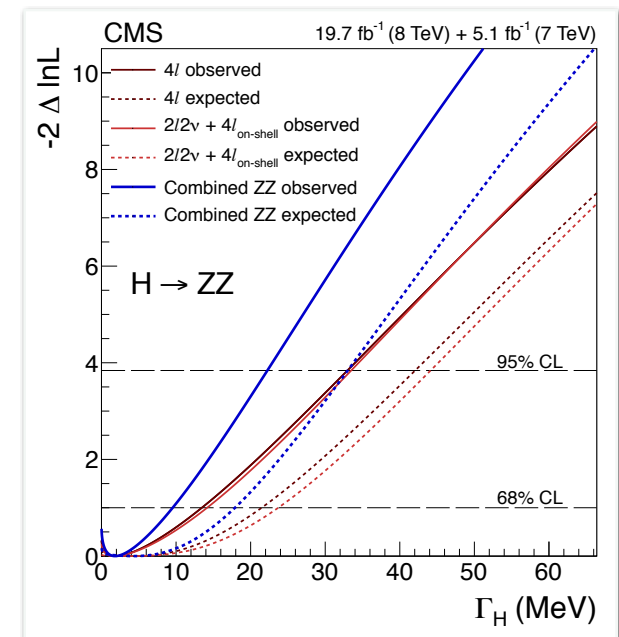
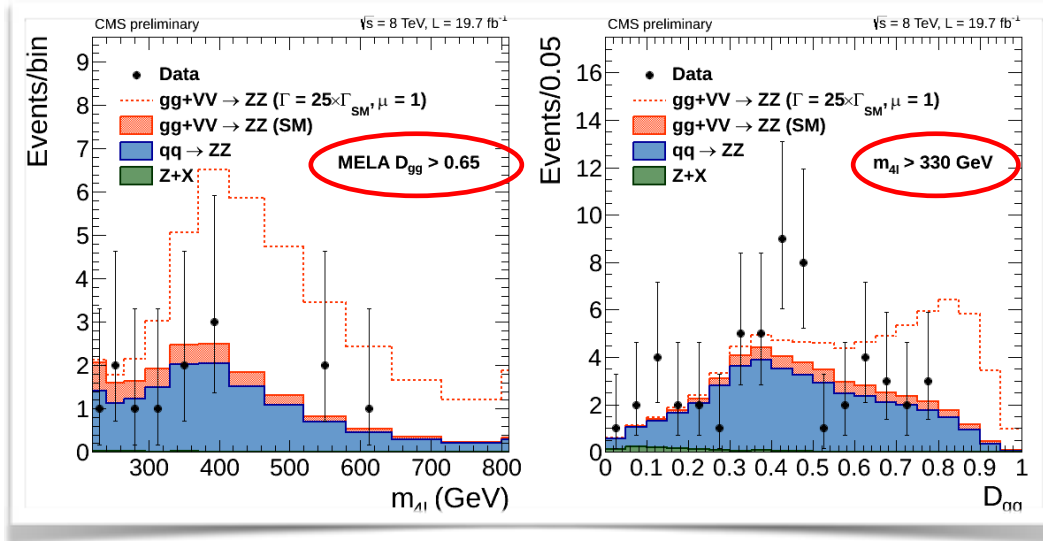
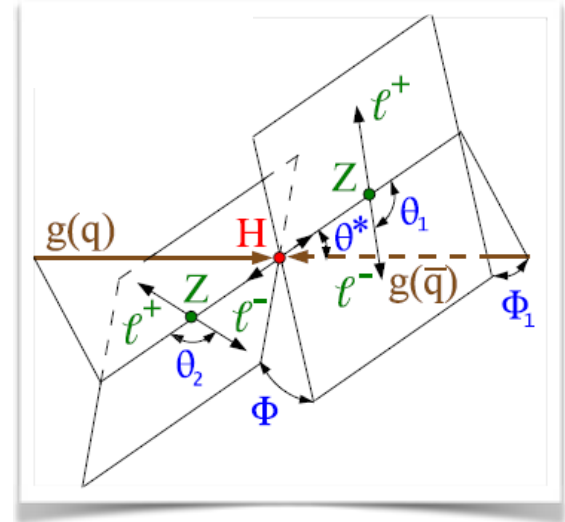


Significant interference with the SM
VV background at high mass



CMS measurement of the width

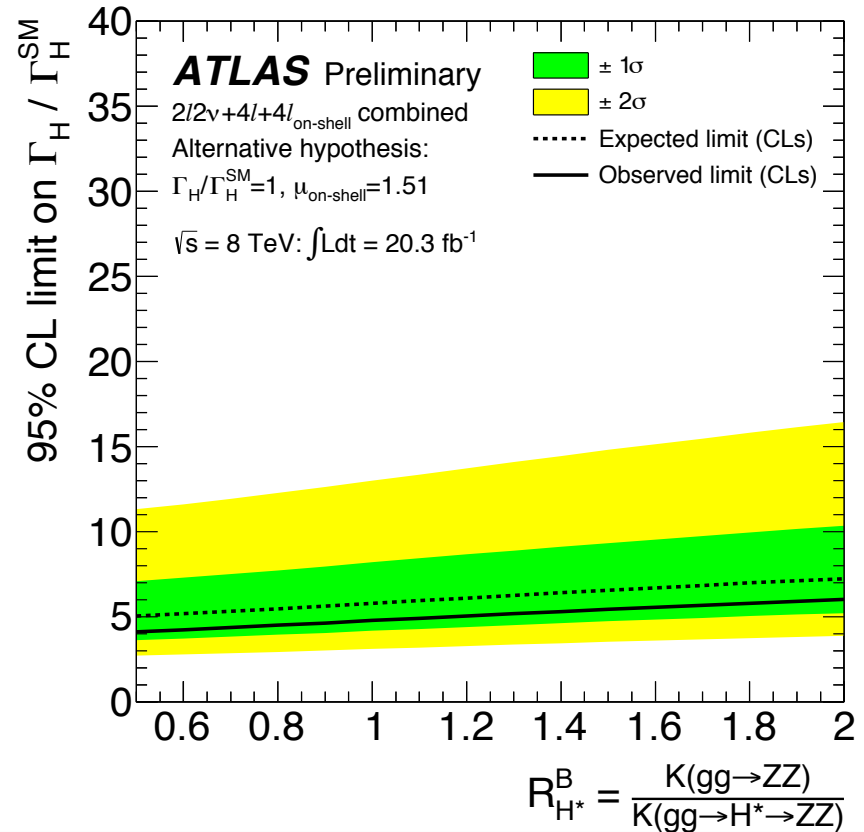
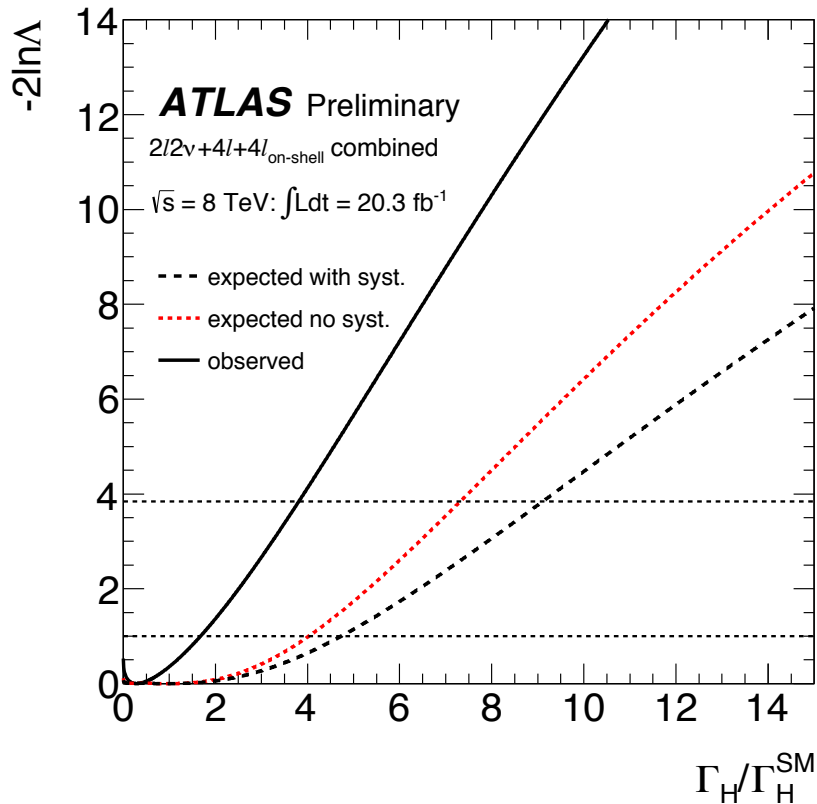
- First measured by CMS (Moriond 2014) using the 4l and 2l2v using a matrix element likelihood approach (MELA)
- Combined observed (expected) values
 - $r < 4.2$ (8.5) @ 96% CL
 - $\Gamma < 17.4$ (35.3) MeV
- Two orders of magnitude better than direct measurements



ATLAS width result

ATLAS-CONF-2014-042

- Similar result from ATLAS during 2014
- Additionally, showed the dependence on the k-factor for the ZZ background
 - No strong dependence observed





Couplings Combination

Combinations

- As we've only seen a small* amount of data from the LHC so far, we want to combine the results from ATLAS and CMS so that we get the most accurate measurements

$$\mu = 1.20 \pm 0.15$$

$$\mu = 0.97 \pm 0.14$$

Average: 1.08

But wait, what about the errors?

$$\bar{x} = \frac{\sum_{i=1}^n x_i \sigma_i^{-2}}{\sum_{i=1}^n \sigma_i^{-2}}$$

Weighted average: 1.07

Perhaps some of the errors should be correlated? e.g. the theoretical uncertainties

Full combination

2 years?

EUROPEAN ORGANISATION FOR NUCLEAR RESEARCH (CERN)



Submitted to: JHEP



CERN-EP-2016-100
8th June 2016

Measurements of the Higgs boson production and decay rates and constraints on its couplings from a combined ATLAS and CMS analysis of the LHC pp collision data at $\sqrt{s} = 7$ and 8 TeV

The ATLAS and CMS Collaborations

Combination Inputs

Channel	References for individual publications		Signal strength [μ] from results in this paper (Section 5.2)		Signal significance [σ]	
	ATLAS	CMS	ATLAS	CMS	ATLAS	CMS
$H \rightarrow \gamma\gamma$	[91]	[92]	1.14 ^{+0.27} _{-0.25} (+0.26) (-0.24)	1.11 ^{+0.25} _{-0.23} (+0.23) (-0.21)	5.0 (4.6)	5.6 (5.1)
$H \rightarrow ZZ$	[93]	[94]	1.52 ^{+0.40} _{-0.34} (+0.32) (-0.27)	1.04 ^{+0.32} _{-0.26} (+0.30) (-0.25)	7.6 (5.6)	7.0 (6.8)
$H \rightarrow WW$	[95, 96]	[97]	1.22 ^{+0.23} _{-0.21} (+0.21) (-0.20)	0.90 ^{+0.23} _{-0.21} (+0.23) (-0.20)	6.8 (5.8)	4.8 (5.6)
$H \rightarrow \tau\tau$	[98]	[99]	1.41 ^{+0.40} _{-0.36} (+0.37) (-0.33)	0.88 ^{+0.30} _{-0.28} (+0.31) (-0.29)	4.4 (3.3)	3.4 (3.7)
$H \rightarrow bb$	[100]	[101]	0.62 ^{+0.37} _{-0.37} (+0.39) (-0.37)	0.81 ^{+0.45} _{-0.43} (+0.45) (-0.43)	1.7 (2.7)	2.0 (2.5)
$H \rightarrow \mu\mu$	[102]	[103]	-0.6 ^{+3.6} _{-3.6} (+3.6) (-3.6)	0.9 ^{+3.6} _{-3.5} (+3.3) (-3.2)		
ttH production	[77, 104, 105]	[107]	1.9 ^{+0.8} _{-0.7} (+0.7) (-0.7)	2.9 ^{+1.0} _{-0.9} (+0.9) (-0.8)	2.7 (1.6)	3.6 (1.3)

No theory assumptions, independent

Production process	Decay channel				
	$H \rightarrow \gamma\gamma$	$H \rightarrow ZZ$	$H \rightarrow WW$	$H \rightarrow \tau\tau$	$H \rightarrow bb$
ggF	$(\sigma \cdot B)_{ggF}^{\gamma\gamma}$	$(\sigma \cdot B)_{ggF}^{ZZ}$	$(\sigma \cdot B)_{ggF}^{WW}$	$(\sigma \cdot B)_{ggF}^{\tau\tau}$	-
VBF	$(\sigma \cdot B)_{VBF}^{\gamma\gamma}$	$(\sigma \cdot B)_{VBF}^{ZZ}$	$(\sigma \cdot B)_{VBF}^{WW}$	$(\sigma \cdot B)_{VBF}^{\tau\tau}$	-
WH	$(\sigma \cdot B)_{WH}^{\gamma\gamma}$	$(\sigma \cdot B)_{WH}^{ZZ}$	$(\sigma \cdot B)_{WH}^{WW}$	$(\sigma \cdot B)_{WH}^{\tau\tau}$	$(\sigma \cdot B)_{WH}^{bb}$
ZH	$(\sigma \cdot B)_{ZH}^{\gamma\gamma}$	$(\sigma \cdot B)_{ZH}^{ZZ}$	$(\sigma \cdot B)_{ZH}^{WW}$	$(\sigma \cdot B)_{ZH}^{\tau\tau}$	$(\sigma \cdot B)_{ZH}^{bb}$
ttH	$(\sigma \cdot B)_{ttH}^{\gamma\gamma}$	$(\sigma \cdot B)_{ttH}^{ZZ}$	$(\sigma \cdot B)_{ttH}^{WW}$	$(\sigma \cdot B)_{ttH}^{\tau\tau}$	$(\sigma \cdot B)_{ttH}^{bb}$

No
measurement

Insufficient
precision

Take SM
values

23 parameter fit

Only quote results of 20 parameters: ZH/WH
and ttH have too low sensitivity

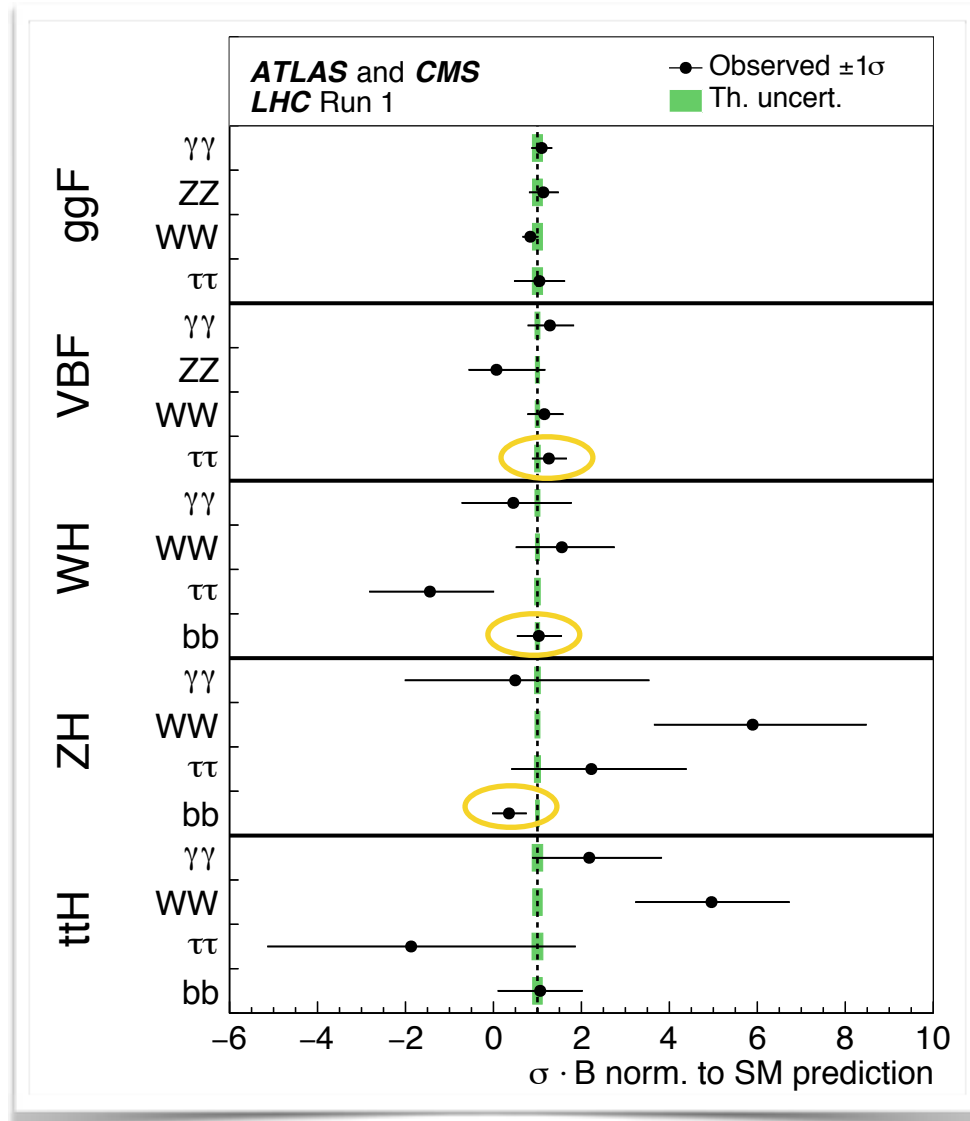
Independent cross-section and BR Results

Very generic: assume all cross-sections and BR to be independent

SM Higgs @ 1
No Higgs @ 0

Very good agreement

Larger errors, typically agreement at the 1 and 2 σ level



Statistical
+systematic
uncertainty as error
bars

Very large scale!

Correlations

Significant anti-correlations

