

Hadronic production of double heavy baryon Ξ_{cc} with excited diquark

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$QQ'q$ – system:

$$\Xi_{QQ'} = QQ'q, \quad q = (u, d)$$

$$\Omega_{QQ'} = QQ'q, \quad q = s$$



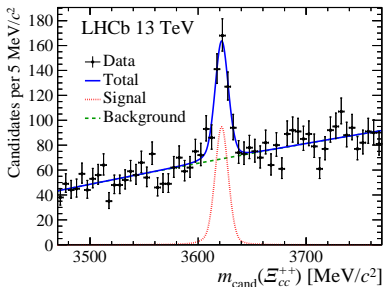
Key features:

- Hierarchy of interaction scales: $m_{Q,Q'} \gg m_{Q,Q'} \cdot v \gg \Lambda_{QCD}$
- In the limit $m_Q \rightarrow \infty$ diquark acts as a local heavy source of gluon field
- Sequential interaction:
 - $[QQ']_{\bar{3}_c}$
 - $[QQ']_{\bar{3}_c} + q_{3_c}$

SEP 2017.

Observation of Ξ_{cc}^{++} in decay mode

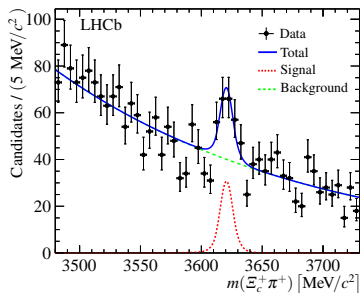
$$\Xi_{cc}^{++} \rightarrow \Lambda_c K^- \pi^+ \pi^+$$



OCT 2018.

Observation of the decay

$$\Xi_{cc}^{++} \rightarrow \Xi_c^+ \pi^+$$



[R. Aaij et al. (LHCb), Phys. Rev. Lett. 119, 112001 (2017)

R. Aaij et al. 1807.01919 (2018)]

In case $Q = Q'$ wave function of a diquark is AS:

- S,D-wave $S = 1$
- P-wave $S = 0$

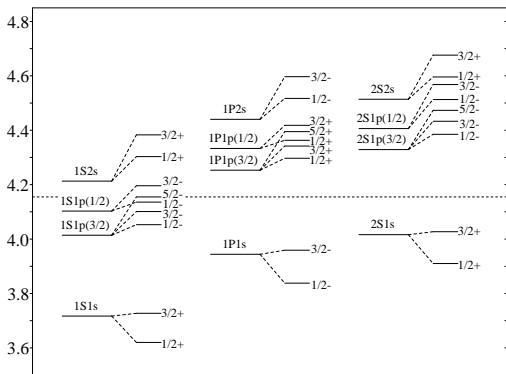
Metastable states:
 $\Xi_{cc}(1P)$, $\Xi_{cc}(2P)$

Transitions

$$1P1s \rightarrow 1S1s$$

$$2P1s \rightarrow 1S1s$$

require a simultaneous change of the diquark's spin and angular momentum



Mass spectrum of Ξ_{cc} baryon in GeV. Dashed line shows the $\Lambda_c D$ threshold.

Transitions are suppressed as $\Lambda_{\text{QCD}}^2/m_c^2$. Both neutral and charged modes are expected. Soft photon is being lost.

$$\Xi_{cc}^{P*} (J^P = 3/2^-) : M + \Delta M^* + \Delta M$$

$$\Xi_{cc}^P (J^P = 1/2^-) : M + \Delta \tilde{M}$$

$$\Xi_{cc}^* (J^P = 3/2^+) : M + \Delta M^*$$

$$\Xi_{cc} (J^P = 1/2^+) : M$$

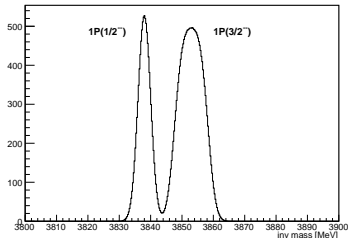
$$\Xi_{cc}^{P*} \rightarrow \Xi_{cc}^* \pi \rightarrow \Xi_{cc} \pi \gamma$$

$$\Xi_{cc}^P \rightarrow \Xi_{cc} \pi$$

Peak corresponding to $1P(3/2^-)$:

- is shifted to $\Delta M^* \approx 100$ MeV
- has an extra widening ≈ 10 MeV

$$M_{\text{inv}}^{\text{max}} - M_{\text{inv}}^{\text{min}} \approx 2\Delta M^* \sqrt{(\Delta M/M)^2 - (m_\pi/M)^2}$$



Possible mass distribution

Production of double heavy baryons

Several stages:

- Hard production of two heavy quark pairs
- Soft formation of the diquark in color antitriplet
- Hadronization into baryon

$$Q\bar{Q}Q'\bar{Q}' \rightarrow [QQ']_{\bar{3}_c} \rightarrow QQ'q$$

- DPS mechanism is not involved in production of QQ'
- In very rough approximation

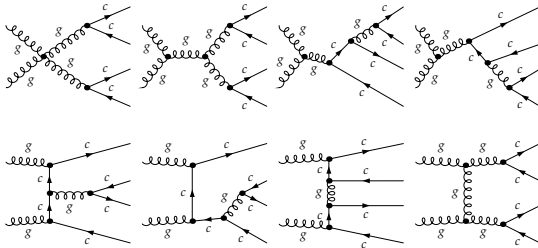
$$P\left([QQ']_{\bar{3}_c} \rightarrow QQ'q\right) = 1$$

the probability of diquark dissociation is not accounted,
the whole momentum of the diquark is transmitted to the baryon

Hadronic production of $c\bar{c}c\bar{c}$

- gg contribution dominates towards $q\bar{q}$
- gq diagrams may be neglected given that $p_T \sim m_c$
- 36 diagrams of the 4th order by α_s in perturbative QCD

$$g \ g \rightarrow c \ c \ c \ c$$



- FeynArts
Generation and
visualization of
feynman
diagrams

FeynArts: analytical expressions

[T. Hahn, Comput. Phys. Commun. 140, 418-431 (2001)]

- Spin selection for S-wave state:

$$N(1, 1) = |\uparrow\uparrow\rangle$$

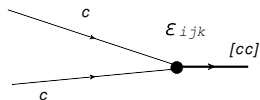
$$N(1, 0) = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)$$

$$N(1, -1) = |\downarrow\downarrow\rangle$$

- Spin selection for P-wave state:

$$N(0, 0) = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

- Colour factor



- Convolution with the wave function of the diquark:

$$A^{S J j_z} = \int T_{c\bar{c}c\bar{c}}^{S s_z}(p_i, k(\vec{q})) \cdot (\Psi_{cc}^{L l_z}(\vec{q}))^* \cdot C_{s_z l_z}^{J j_z} \frac{d\vec{q}}{(2\pi)^3}$$

$T(p_i, k(\vec{q}))$ — amplitude of the hard production of two $c\bar{c}$ pairs

$\Psi_{c\bar{c}}(\vec{q})$ — wave function of the diquark

p_i — four momenta of $[c\bar{c}]_{\bar{3}_c}$, c_1, c_2

\vec{q} — three momentum of c -quark in the diquark rest frame

$(0, \vec{q}) \rightarrow k(\vec{q})$ — boost from the diquark rest frame to the frame of $c\bar{c}c\bar{c}$ production

$C_{s_z l_z}^{J j_z}$ — Clebsh-Gordon coefficients

Under assumption of small dependence of $T_{c\bar{c}c\bar{c}}^{Ssz}$ on $k(\vec{q})$:

$$A \sim \int d^3q \Psi(\vec{q}) \left\{ T(p_i, \vec{q}) \Big|_{\vec{q}=0} + \vec{q} \frac{\partial}{\partial \vec{q}} T(p_i, \vec{q}) \Big|_{\vec{q}=0} + \dots \right\}$$

- For S -wave state:

$$A^{sz} \sim R_s(0) \cdot T_{c\bar{c}c\bar{c}}^{sz}(p_i) \Big|_{\vec{q}=0}$$

- For scalar P -wave state:

$$A^{lz} \sim R'_p(0) \cdot \mathcal{L}^{lz} T_{c\bar{c}c\bar{c}}(p_i) \Big|_{\vec{q}=0}$$

Differential operator:

$$\mathcal{L}^{lz} = \begin{cases} \mathcal{L}^{-1} = \frac{1}{\sqrt{2}} \left(\frac{\partial}{\partial q_x} + i \frac{\partial}{\partial q_y} \right) \\ \mathcal{L}^0 = \frac{\partial}{\partial q_z} \\ \mathcal{L}^{+1} = -\frac{1}{\sqrt{2}} \left(\frac{\partial}{\partial q_x} - i \frac{\partial}{\partial q_y} \right) \end{cases}$$

[A.V. Berezhnoy, V.V. Kiselev, A.K. Likhoded, Phys. Lett. B 381 (1996) 341–347]

Input parameters:

$$m, M, R_s(0), R'_p(0), \sqrt{s_{pp}}$$

Fortran code:

- Calculation of matrix elements with spin selection
- Differentiation:

$$\left. \frac{\partial T(p_i, k(\vec{q}))}{\partial q_j} \right|_{\vec{q}=0} \approx \frac{T(p_i, k(\varepsilon^j)) - T(p_i, 0)}{\varepsilon}$$

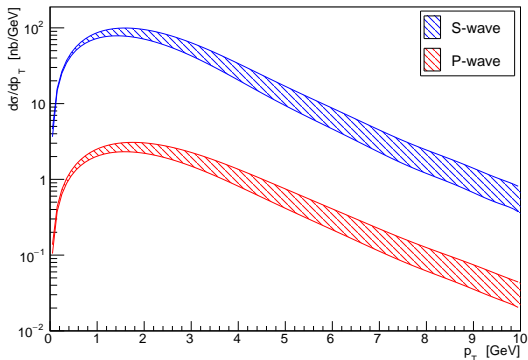
- Integration over phase space within RAMBO generator
- Convolution with proton distribution functions CT14 at varied scale:

$$\sigma_{pp} = \int \hat{\sigma}_{gg}(\hat{s}_{gg}) f_{g1}(x_1) f_{g2}(x_2) dx_1 dx_2$$

[S. Dulat, T.J. Hou, J. Gao et al., EPJ Web Conf. 120, 07003 (2016)]

Cross sections for Ξ_{cc}
excitations at
 $\sqrt{s} = 13\text{TeV}$

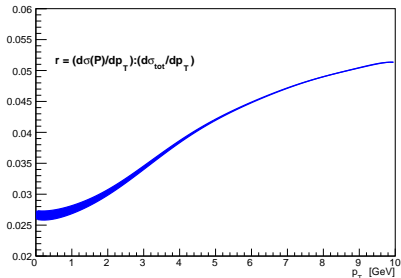
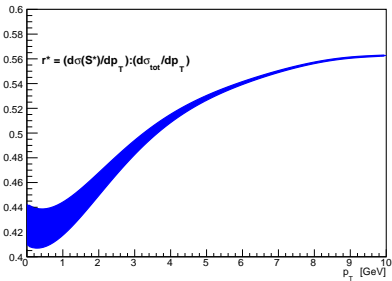
state	σ , [nb]
$1S$	$120 \div 170$
$2S$	$60 \div 90$
$3S$	$40 \div 70$
$1P$	$4 \div 6$
$2P$	$4 \div 5$



Relative yields

$r(\Xi_{cc}^*) = \sigma(\Xi_{cc}^*) / \sigma(\Xi_{cc})$ for
 Ξ_{cc} excitations

state	r, %
1S	49 ÷ 52
2S	26 ÷ 27
3S	18 ÷ 20
1P	2
2P	1 ÷ 2



- Relative yields for production of S-wave and P-wave excitations of Ξ_{cc} have been estimated in LHCb' kinematics:

$$\sigma(\Xi_{cc}^P) / \sigma(\Xi_{cc}) \approx 3\%$$

$$\sigma(\Xi_{cc}^{S*}) / \sigma(\Xi_{cc}) \approx 45\%$$

- Nowadays LHCb collaboration observes hundreds of events for Ξ_{cc} ground state. 313 ± 33 particles are identified in decay:

$$\Xi_{cc}^{++} \rightarrow \Lambda_c^+ K^- \pi^+ \pi^+$$

- It is worth expecting that the excited states of Ξ_{cc} will be extracted at Run 3 in single pion transitions:

$$\Xi_{cc}^P \rightarrow \Xi_{cc} \pi$$

$$\Xi_{cc}^{S*} \rightarrow \Xi_{cc} \pi$$

Thank you for attention!

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