Recent Heavy Flavour results from ATLAS and CMS



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Selected recent results

Outline : ATLAS + CMS @ LHC

Di-charmonia resonances (ATLAS-CONF-2022-040, CMS PAS BPH-21-003) **Triple J/** ψ production (arXiv:2111.05370 accepted by Nature Physics) **B**_c⁺ production and decays (PRD 104(2021), JHEP 08(2022)87) *B*_(s)⁰ $\rightarrow \mu^+\mu^-$ decays (FCNC) (CMS PAS BPH-21-006) **b** quark fragmentation (JHEP 12 (2021) 131) **B**⁰_(s) $\rightarrow \psi(2S)K^0_S(X)$ decays (EPJC 82 (2022) 499) **Prospects & Summary**

ATLAS+CMS @ LHC

ATLAS : weight : ~ 7000 tons 44m



Run I : 2010 - 2012



Run II : 2015 - 2018





Dimuon triggers

Di-charmonia tetraquarks (cccc)

ATLAS-CONF-2022-040

$J/\psi J/\psi$ and $J/\psi + \psi(2S)$ in 4μ final state studied at ATLAS using 139 fb⁻¹ of pp at Vs = 13 TeV

Blind analysis

Di-charmonium Process	Generator	PDF	Parton shower	Tune
SPS DPS Non-prompt X(6900)	Рутніа 8.244 Рутніа 8.244 Рутніа 8.244 ЈНU	NNPDF23LO NNPDF23LO NNPDF23LO CTEQ6L1	Pythia 8.244+NNPDF23LO	A14 A14 A14 A14 A14

ΔR > 0.25



Event selection, reconstruction and definition of signal and control regions

Signal region	Signal region SPS/DPS control region non-prompt region			
E	Di-muon or tri-muon triggers,			
Opposite charged	I muons from the same J/ψ or $\psi(2S)$ v	ertex,		
Loose muon ID, $p_{\rm T}^{1,2,3,4} >$	• 4, 4, 3, 3 GeV and $ \eta_{1,2,3,4} < 2.5$ for t	the four muons		
$m_{J/\psi} \in \{2.94,$	3.25} GeV, or $m_{\psi(2S)} \in \{3.56, 3.80\}$ C	GeV,		
Loose vertex cuts $\chi^2_{4\mu}/N < 40$ and $\chi^2_{di-\mu}/N < 100$,				
Vertex $\chi^2_{4\mu}/N < 3$,				
$L_{xy}^{4\mu} < 0.2 \text{ mm}, L_{xy}^{\text{di-}\mu} < 0.3 \text{ mm},$ Vertex $\chi_{4\mu}^2/N > 6$,				
$m_{4\mu} < 7.5$ GeV,	$m_{4\mu} < 7.5 \text{ GeV},$ 7.5 GeV $< m_{4\mu} < 12.0 \text{ GeV} (\text{SPS})$ $ L_{xy}^{\text{di-}\mu} > 0.4 \text{ mm}$			
$R < 0.25$ between charmonia 14.0 GeV < $m_{4\mu} < 25.0$ GeV (DPS)				

ΔR < 0.25



Di-charmonia tetraquarks (cccc)

ATLAS-CONF-2022-040



0

7.5

8.5

9

m_{4µ}^{con} [GeV]

8



Di-J/ ψ tetraquarks ($c\bar{c}c\bar{c}$)

CMS PAS BPH-21-003

loose muon identification

 $J/\psi J/\psi$ ($\Rightarrow 4\mu$) spectrum studied at CMS using 135 fb⁻¹ of pp collisions at $\sqrt{s} = 13$ TeV (2016-2018)

Event selection and reconstruction:

- **3-µ trigger**: $\mu^+\mu^-$ from J/ ψ + third muon (on muons from J/ ψ : $p_+(\mu^+\mu^-) > 3.5$ GeV in 2017-2018)
- blinded signal region $m(J/\psi J/\psi)$ in [6.2, 7.8] GeV (from preliminary investigation on 2011-2012 data)
- $p_{\tau}(\mu) > 2.0 \text{ GeV};$ $|\eta(\mu)| < 2.4;$
- $m(\mu^+\mu)$ in [2.95, 3.25] GeV; $p_{\tau}(\mu^+\mu) > 3.5 \text{ GeV}$ $P_{vt}(\mu^+\mu) > 0.5\%$
- common vertex fit: $P_{_{vtx}}(4\mu) > 0.5\%$
- Arbitration of multiple candidates:
 - Select best combination of same 4μ (from MC: 0.2%) Ο
 - Keep all candidates arising from more than four muons (from MC: 0.2%) Ο

Background model:

- NRSPS: threshold function * pol2 * exponential
- NRDPS: threshold function * pol2 * exponential
- Relativistic Breit-Wigner near $J/\psi J/\psi$ threshold BWO:
 - inadequacy of NRSPS near threshold 0
 - feed-down of partially reconstructed higher mass states 0
 - possible coupled-channel interactions, pomeron-exchange processes, etc. 0



No interference assumed

$$\chi_m^2 = \left(\frac{m_1(\mu^+\mu^-) - M_{J/\psi}}{\sigma_{m_1}}\right)^2 + \left(\frac{m_2(\mu^+\mu^-) - M_{J/\psi}}{\sigma_{m_2}}\right)^2$$

$$= \left(\frac{m_1(\mu^+\mu^-) - M_{J/\psi}}{\sigma_{m_1}}\right)^- + \left(\frac{m_2(\mu^+\mu^-) - M_{J/\psi}}{\sigma_{m_2}}\right)^-$$

Di-J/ψ tetraquarks (*cc̄cc̄*)

CMS PAS BPH-21-003

25 MeV

Candidates /

Unc.

CMS Preliminary

CMS signal + background model

Three Relativistic Breit-Wigner (J^P = 0⁺) are considered

Mass (MeV)	Width (MeV)	Local stat. signif.
6552 ± 10 ± 12	124 ± 29 ± 34	> 5.7ơ
6927 ± 9 ± 5	122 ± 22 ± 19	> 9.4ơ
7287 ± 19 ± 5	95 ± 46 ± 20	> 4.1o
	Mass (MeV) 6552 ± 10 ± 12 6927 ± 9 ± 5 7287 ± 19 ± 5	Mass (MeV) Width (MeV) 6552 ± 10 ± 12 124 ± 29 ± 34 6927 ± 9 ± 5 122 ± 22 ± 19 7287 ± 19 ± 5 95 ± 46 ± 20

first error is statistic, second is systematic error

LHCb signal models + CMS background

- Model 1:
 - X(6900) parameters in agreement
 - but dip at 6.7 not well described
- Model 2:
 - Larger X(6700) amplitude
 - X(7200) region not well described



LHCb non-interference model I

X(6900) confirmed at CMS

Values consistent with LHCb

LHCb interference model II

135 fb⁻¹ (13 TeV)

Data - Fit

BW1 BW2[X(6900)]

BW3 - Background

Di-charmonia tetraquarks (*cccc*)

(GeV)	m_0	Γ_0	m_1	Γ_1
di_J/y_{ℓ}	$6.22 \pm 0.05^{+0.04}_{-0.05}$	$0.31 \pm 0.12^{+0.07}_{-0.08}$	$6.62 \pm 0.03 ^{+0.02}_{-0.01}$	$0.31 \pm 0.09^{+0.06}_{-0.11}$
ur o y φ	m_2 Γ_2		-	_
	$6.87 \pm 0.03^{+0.06}_{-0.01}$	$0.12 \pm 0.04 \substack{+0.03 \\ -0.01}$	-	_
(GeV)		<i>m</i> ₃	Γ_3	
J/u/u+u/(2S)	model A	$7.22 \pm 0.03 \substack{+0.02 \\ -0.03}$	$0.10\substack{+0.13+0.06\\-0.07-0.05}$	_
• / • · • (=3)	model B	$6.78 \pm 0.36^{+0.35}_{-0.54}$	$0.39 \pm 0.11 \substack{+0.11 \\ -0.07}$	—

CMS

	BW1	BW2	BW3
т	$6552\pm10\pm12$	$6927\pm9\pm5$	$7287 \pm 19 \pm 5$
Γ	$124\pm29\pm34$	$122\pm22\pm19$	$95\pm46\pm20$
N	474 ± 113	492 ± 75	156 ± 56

X(6900) well seen by all 3 experiments

Signatures for a bump at 7.2-7.3 GeV In all 3 experiments

Nature of the low-mass bump to be clarified

Role of reflections from decays to $J/\psi+\psi(2S)$, $J/\psi+\chi_c$ to be studied

ATLAS-CONF-2022-040





Sci.Bull. 65 (2020) 1983



Observation of triple J/ ψ meson production



CMS, arXiv:2111.05370 accepted by Nature Physics

$$N_{\rm sig}^{\rm 3J/\psi} = 5.0^{+2.6}_{-1.9}$$

5.5 σ (MC toys)

 $\sigma(pp \rightarrow J/\psi J/\psi J/\psi X) =$ 272⁺¹⁴¹₋₁₀₄(stat)±17(syst) fb

f_{SPS} = 6%

f_{DPS} = 74%

f_{TPS} = 20%

Observation of triple J/ ψ meson production

$$\begin{aligned} p_{DPS}^{pp\rightarrow q_{1}q_{2}+x_{n}} = \begin{pmatrix} m \\ 2 \end{pmatrix} \underbrace{p_{PS}^{pp\rightarrow q_{1}+x_{n}} e_{SPS}^{pp\rightarrow q_{2}+x_{n}} e_{SPS}^{pp\rightarrow q_{2}+x_{$$

CMS, arXiv:2111.05370

accepted by Nature Physics

Measurement of the relative B_c^{\pm}/B^{\pm} production cross section with the ATLAS detector at $\sqrt{s} = 8$ TeV PHYSICAL REVIEW D 104, 012010 (2021)



 $B^+ \rightarrow J/\psi K^+$



 B_c^+ and B^+ yields measured using di-muon trigger

Their ratios, corrected for acceptances and efficiencies, measured in two p_T bins (13-22 GeV, >22 GeV) and two |y| bins (<0.75, 0.75-2.3)

B_c^+ / B^+ x-section ratios at 8 TeV

 $\frac{\sigma(B_c^{\pm}) \cdot \mathcal{B}(B_c^{\pm} \to J/\psi\pi^{\pm})}{\sigma(B^{\pm}) \cdot \mathcal{B}(B^{\pm} \to J/\psiK^{\pm})} = (0.34 \pm 0.04_{\text{stat}} + 0.06_{-0.02}_{-0.02} \text{syst} \pm 0.01_{\text{lifetime}})\%$

Compatible with CMS/LHCb



The ratio decreases with p_T

Differences in production? hadronization?

No significant |y| dependence

B_c^+ / B^+ x-section ratios at LHC

$$\frac{\sigma(B_c^{\pm}) \cdot \mathcal{B}(B_c^{\pm} \to J/\psi\pi^{\pm})}{\sigma(B^{\pm}) \cdot \mathcal{B}(B^{\pm} \to J/\psi K^{\pm})} =$$

0.683 \pm 0.018 \pm 0.009**pT < 20 GeV**, 2.0 < |y| < 4.5</th>LHCb at 8 TeV**0.48** \pm 0.05 \pm 0.03 \pm 0.05**pT > 15 GeV**, |y| < 1.6</th>CMS at 7 TeV**0.44** \pm 0.07 $^{+0.09}_{-0.04}$ \pm 0.01**13 < pT < 22 GeV**, |y| < 2.3</th>ATLAS at 8 TeV**0.24** \pm 0.04 $^{+0.05}_{-0.01}$ \pm 0.01**pT > 22 GeV**, |y| < 2.3</th>ATLAS at 8 TeV

The ratio decreases with p_T

Differences in production? hadronization?



Measurement of $B_c^+ \rightarrow J/\psi D_s^{(*)+}$

ATLAS, JHEP 08(2022)87



Figure 1: Feynman diagrams for $B_c^+ \to J/\psi D_s^{(*)+}$ decays: (a) colour-favoured spectator, (b) colour-suppressed spectator, and (c) annihilation topology.

Normalization mode: $B_c^+ \rightarrow J/\psi \pi^+$

Run-I measurements: LHCb

Phys. Rev. D 87 (2013) 112012,

ATLAS Eur. Phys. J. C **76** (2016) 4.

2D unbinned fit of two subsets simultaneously



Normalization mode: $B_c^+ \rightarrow J/\psi \pi^+$

- $\chi^2(B_c^+)/N_{dof} < 1.8 \ (N_{dof} = 4 \text{ in the } B_c^+ \rightarrow J/\psi \pi^+ \text{ vertex fit}),$
- $L_{xy}(B_c^+) > 0.3 \text{ mm},$
- $|d_0^{\text{PV}}(B_c^+)/\sigma_{d_0^{\text{PV}}}(B_c^+)| < 3 \text{ and } |z_0^{\text{PV}}(B_c^+)/\sigma z_0^{\text{PV}}(B_c^+)| < 3,$
- $p_{\rm T}(B_c^+) / \sum p_{\rm T}({\rm trk}) > 0.10$,



Parameter	Value
$m_{B_c^+}$ [MeV] $\sigma_{B_c^+}$ [MeV] $N_{B_c^+ \rightarrow J/\psi \pi^+}$	$\begin{array}{r} 6274.5 \pm 1.5 \\ 47.5 \pm 2.5 \\ 8440^{+550}_{-470} \end{array}$

BCVEGPY 2.2 interfaced to Pythia 8.244

Mode	$\epsilon^{\mathrm{DS1}}_{B^+_c o J/\psi X}$ [%]	$\epsilon^{\mathrm{DS1\&2}}_{B^+_c \to J/\psi X} [\%]$
$B_c^+ \to J/\psi D_s^+$	0.971 ± 0.012	1.163 ± 0.013
$B_c^+ \rightarrow J/\psi D_s^{*+}, A_{00}$	0.916 ± 0.012	1.088 ± 0.012
$B_c^+ \rightarrow J/\psi D_s^{*+}, A_{\pm\pm}$	0.868 ± 0.010	1.049 ± 0.011
$B_c^+ o J/\psi \pi^+$	2.169 ± 0.018	_

Results:

$$R_{D_{s}^{(*)+}/\pi^{+}} = \frac{\mathcal{B}(B_{c}^{+} \to J/\psi D_{s}^{(*)+})}{\mathcal{B}(B_{c}^{+} \to J/\psi \pi^{+})} = \frac{N_{c}^{\text{DS1}}}{N_{B_{c}^{+} \to J/\psi \pi^{+}}} \times \frac{\epsilon_{B_{c}^{+} \to J/\psi \pi^{+}}}{\epsilon_{B_{c}^{+} \to J/\psi D_{s}^{(*)+}}} \times \frac{1}{\mathcal{B}(D_{s}^{+} \to \phi(K^{+}K^{-})\pi^{+})}$$

$$R_{D_{s}^{*+}/D_{s}^{*}} = \frac{\mathcal{B}(B_{c}^{+} \to J/\psi D_{s}^{*+})}{\mathcal{B}(B_{c}^{+} \to J/\psi D_{s}^{*})} = \frac{N_{B_{c}^{+} \to J/\psi D_{s}^{*+}}^{\mathrm{DS1\&2}}}{N_{B_{c}^{+} \to J/\psi D_{s}^{*}}^{\mathrm{DS1\&2}}} \times \frac{\epsilon_{B_{c}^{+} \to J/\psi D_{s}^{+}}^{\mathrm{DS1\&2}}}{\epsilon_{B_{c}^{+} \to J/\psi D_{s}^{*+}}^{\mathrm{DS1\&2}}} = r_{D_{s}^{*+}/D_{s}^{+}} \times \frac{\epsilon_{B_{c}^{+} \to J/\psi D_{s}^{+}}^{\mathrm{DS1\&2}}}{\epsilon_{B_{c}^{+} \to J/\psi D_{s}^{*+}}^{\mathrm{DS1\&2}}}$$

$$\Gamma_{\pm\pm}/\Gamma = f_{\pm\pm} \times \frac{\epsilon_{B_c^+ \to J/\psi D_s^{*+}}^{\text{DS1\&2}}}{\epsilon_{B_c^+ \to J/\psi D_s^{*+}, A_{\pm\pm}}^{\text{DS1\&2}}}$$

 $\begin{aligned} R_{D_s^*/\pi^+} &= 2.76 \pm 0.33 \pm 0.30 \pm 0.16 \\ R_{D_s^{*+}/\pi^+} &= 5.33 \pm 0.61 \pm 0.67 \pm 0.32 \\ R_{D_s^{*+}/D_s^+} &= 1.93 \pm 0.24 \pm 0.10 \\ \Gamma_{\pm\pm}/\Gamma &= 0.70 \pm 0.10 \pm 0.04 \end{aligned}$

Comparison with previous results and models



QCD PM (Colangelo, De Fazio) – good agreement, no uncertainties, no prediction for $\Gamma_{\pm\pm}/\Gamma$ CCQM (M.Ivanov et al) – agrees only for $R_{D_s^{*+}/\pi^+}$

Predictions for $\Gamma_{\pm\pm}/\Gamma$ are below or around 0.5 while data agree with naïve spin counting 2/3

Comparison with other B-meson decays



Figure 1: Feynman diagrams for $B_c^+ \to J/\psi D_s^{(*)+}$ decays: (a) colour-favoured spectator, (b) colour-suppressed spectator, and (c) annihilation topology.

$$R_{D_{s}^{*}/\pi^{+}} \approx \frac{\Gamma(B \to D^{*}D_{s}^{*})}{\Gamma(B \to \bar{D}^{*}\pi^{+})},$$

$$R_{D_{s}^{*+}/\pi^{+}} \approx \frac{\Gamma(B \to \bar{D}^{*}D_{s}^{*+})}{\Gamma(B \to \bar{D}^{*}\pi^{+})},$$

$$R_{D_{s}^{*+}/D_{s}^{*}} \approx \frac{\Gamma(B \to \bar{D}^{*}D_{s}^{*+})}{\Gamma(B \to \bar{D}^{*}D_{s}^{*})},$$

$$R_{D_{s}^{*+}/D_{s}^{*}} \sim \frac{\Gamma(B \to J/\psi K^{*})}{\Gamma(B \to J/\psi K)},$$

 $\Gamma_{\pm\pm}/\Gamma$

Assuming that the colour-favoured spectator diagram dominates the decay amplitudes

Assuming that the colour-supressed spectator diagram dominates the decay amplitudes

Can be also compared to $B_s^0 \rightarrow J/\psi \phi$

Comparison with ligher B-meson decays



Supports the assumtion that the colour-favoured spectator diagram dominates the decay amplitudes



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$B_{(s)}^{0} \rightarrow \mu^{+} \mu^{-}$ (FCNC, CKM and helicity suppressed)



$$\begin{split} & \mathcal{B}(\mathsf{B}^0_{\mathrm{s}} \to \mu^+ \mu^-) = \mathcal{B}(\mathsf{B}^+ \to \mathsf{J}/\psi\mathsf{K}^+) \frac{N_{\mathsf{B}^0_{\mathrm{s}} \to \mu^+ \mu^-}}{N_{\mathsf{B}^+ \to \mathsf{J}/\psi\mathsf{K}^+}} \frac{\varepsilon_{\mathsf{B}^+ \to \mathsf{J}/\psi\mathsf{K}^+}}{\varepsilon_{\mathsf{B}^0_{\mathrm{s}} \to \mu^+ \mu^-}} \frac{f_{\mathrm{u}}}{f_{\mathrm{s}}}, \\ & \mathcal{B}(\mathsf{B}^0_{\mathrm{s}} \to \mu^+ \mu^-) = \mathcal{B}(\mathsf{B}^0_{\mathrm{s}} \to \mathsf{J}/\psi\phi) \frac{N_{\mathsf{B}^0_{\mathrm{s}} \to \mu^+ \mu^-}}{N_{\mathsf{B}^0_{\mathrm{s}} \to \mathsf{J}/\psi\phi}} \frac{\varepsilon_{\mathsf{B}^0_{\mathrm{s}} \to \mathsf{J}/\psi\phi}}{\varepsilon_{\mathsf{B}^0_{\mathrm{s}} \to \mu^+ \mu^-}}, \\ & \mathcal{B}(\mathsf{B}^0 \to \mu^+ \mu^-) = \mathcal{B}(\mathsf{B}^+ \to \mathsf{J}/\psi\mathsf{K}^+) \frac{N_{\mathsf{B}^0 \to \mu^+ \mu^-}}{N_{\mathsf{B}^+ \to \mathsf{J}/\psi\mathsf{K}^+}} \frac{\varepsilon_{\mathsf{B}^+ \to \mathsf{J}/\psi\mathsf{K}^+}}{\varepsilon_{\mathsf{B}^0 \to \mu^+ \mu^-}} \frac{f_{\mathrm{u}}}{f_{\mathrm{d}}}, \end{split}$$

CMS PAS BPH-21-006

$$\begin{split} \mathcal{B}(\mathrm{B}^0_{\mathrm{s}} \to \mu^+\mu^-) &= \left[3.83^{+0.38}_{-0.36} \ (\mathrm{stat}) \, {}^{+0.19}_{-0.16} \, (\mathrm{syst}) \, {}^{+0.14}_{-0.13} \, (f_{\mathrm{s}}/f_{\mathrm{u}}) \right] \times 10^{-9}, \\ \mathcal{B}(\mathrm{B}^0 \to \mu^+\mu^-) &= \left[0.37^{+0.75}_{-0.67} \ (\mathrm{stat}) \, {}^{+0.08}_{-0.09} \, (\mathrm{syst}) \right] \times 10^{-10}. \end{split}$$



$${\cal B}({
m B}^0 o\mu^+\mu^-)<1.5 imes10^{-10}$$
 at 90% CL,
 ${\cal B}({
m B}^0 o\mu^+\mu^-)<1.9 imes10^{-10}$ at 95% CL,

Released for

$B_{(s)}^{0} \rightarrow \mu^{+} \mu^{-}$ (FCNC, CKM and helicity suppressed)

ATLAS, JHEP 04(2019)098





agreement with SM: 2.4 σ



agreement with sm: 2.1 σ



CMS PAS BPH-21-006 Released for ICHEP 2022

Perfect agreement!

Not an anomaly anymore

$B_{(s)}^{0} \rightarrow \mu^{+} \mu^{-}$ (FCNC, CKM and helicity suppressed)



Good agreement

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Measurement of *b*-quark fragmentation properties in jets using the decay $B^{\pm} \rightarrow J/\psi K^{\pm}$ in *p p* collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector

ATLAS, JHEP 12 (2021) 131



- All Pythia fragmentation models give a decent description.
- \bigcirc Herwig7 with dipole parton shower overestimates the low z tail at low p_{T}
 - Iarger fraction of jets arising from gluon splittings
- Sherpa (mainly cluster hadronisation model) differs for very high z



All Pythia fragmentation models give a decent description.

Herwig7 with dipole PS overestimates for p_T^{rel} in [1.5, 4.0] GeV at low p_T
 Sherpa (mainly cluster HM) discrepant for low values of p_T^{rel}, gets worse for higher jet p_T.

Measurement of *b*-quark fragmentation properties
in jets using the decay $B^{\pm} \rightarrow J/\psi K^{\pm}$ in ppATLAS, JHEP 12 (2021) 131collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector



- Pythia (A14*) predicts slightly larger <z> and slightly lower <p_T^{rel}>
 Both Herwig7 discrepant at 15-20% level in <p_T^{rel}> profile
- Sherpa (cluster) disagreeing at 10% to 25% for $< p_T^{rel} >$

Observation of $B^0 \rightarrow \psi(2S) K^0_S \pi^+ \pi^-$ and $B^0_S \rightarrow \psi(2S) K^0_S$ decays EPJC 82 (2022) 499



$$R_{\pi^{+}\pi^{-}} = \frac{\mathcal{B}(B^{0} \to \psi(2S) K_{S}^{0} \pi^{+} \pi^{-})}{\mathcal{B}(B^{0} \to \psi(2S) K_{S}^{0})} =$$

0.480 ± 0.013 (stat) ± 0.032 (syst)
$$\mathcal{B}(B^{0} \to \psi(2S) K_{S}^{0} \pi^{+} \pi^{-}) =$$

(13.9 ± 0.4 (stat) ± 0.9 (syst) ± 1.2 (\mathcal{B})) × 10^{-5},

$$R_{\rm s} = \frac{\mathcal{B}({\rm B}_{\rm s}^{0} \to \psi(2{\rm S}){\rm K}_{\rm S}^{0})}{\mathcal{B}({\rm B}^{0} \to \psi(2{\rm S}){\rm K}_{\rm S}^{0})} =$$

$$(3.33 \pm 0.69\,({\rm stat}) \pm 0.11\,({\rm syst}) \pm 0.34\,(f_{\rm s}/f_{\rm d})) \times 10^{-2}$$

 $\mathcal{B}(B^0_s \to \psi(2S)K^0_S) =$

 $(0.97 \pm 0.20 \,(\text{stat}) \pm 0.03 \,(\text{syst}) \pm 0.22 \,(f_{\rm s}/f_{\rm d}) \pm 0.08 \,(\mathcal{B})) \times 10^{-5}$

Observation of $B^0 \rightarrow \psi(2S) K^0_S \pi^+ \pi^-$ and $B^0_S \rightarrow \psi(2S) K^0_S$ decays

EPJC 82 (2022) 499



Iterative 1D reweightings done over 2- and 3-body invariant mass distributions

No significant exotic narrow structures

More data are needed?

Observation of $B^0 \rightarrow \psi(2S) K^0_S \pi^+ \pi^-$ and $B^0_S \rightarrow \psi(2S) K^0_S$ decays EPJC 82 (2022) 499



No significant exotic narrow structures

More data are needed?

Prospects





Shutdown/Technical stop Protons physics Ions Commissioning with beam Hardware commissioning/magnet training

Run 3: ~300 fb-1 HL-LHC: ~3000 fb-1



 $L_{max} = 25.3 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$



HL-LHC projection



Summary : main messages



Exciting results on di-charmonia resonances from ATLAS. CMS and LHCb 2-4 cccc tetraquark candidates actual states and their spin-parities to be clarified

Precision measurements of B_c^+ production and decays the B_c^+ / B^+ ratio decreases with p_T challenging for theory results on $B_c^+ \rightarrow J/\psi D_s^{(*)+}$

supporting small $\sigma_{eff,DPS}$ value for di-onia production

First results on triple J/ψ meson production

 $B_{(s)}^{0} \rightarrow \mu^{+}\mu^{-}$ measurement with full Run-II statistics good agreement with SM - not an anomaly anymore

b quark fragmentation properties with **B**⁺ in jets PYTHIA fragmentation gives a decent description

Observation of $B^0 \rightarrow \psi(2S)K_S^0\pi^+\pi^-$ and $B_s^0 \rightarrow \psi(2S)K_S^0$ decays no significant exotic narrow structures – more data are needed?

New results with full Run-II and first Run-III data for spring/summer conferences

Backup

 $J/\psi J/\psi$ ($\Rightarrow 4\mu$) spectrum studied at LHCb using 9 fb⁻¹ of pp collisions at $\sqrt{s} = 7, 8, 13$ TeV

Sci.Bull. 65 (2020), 23

Two structures are reported:

- A narrow resonance, X(6900), renamed T_{tub}(6900)
- A broad structure near the di-J/ ψ mass threshold

Background contribution for J/ψ pair production:

- Non-Resonant Single Parton Scattering (NRSPS)
- Non-Resonant Double Parton Scattering (DPS)

Two signal + background fit models are considered:

- Model 1 (top) poor description of the "dip" at 6.7 GeV
 - background
 - Breit-Wigner for X(6900)
 - two auxiliary Breit-Wigner (near threshold)
- Model 2 (bottom)
 - a "virtual" X(6700) to interfere with NRSPS is added



Preselection:

- $\chi^2(B_c^+)/N_{dof} < 2$, where $\chi^2(B_c^+)$ is the quality of the B_c^+ cascade vertex fit and $N_{dof} = 8$.
- $L_{xy}(B_c^+) > 0.3$ mm, where $L_{xy}(B_c^+)$ is the transverse distance between the primary vertex (PV) and the B_c^+ candidate vertex projected onto the direction of the B_c^+ transverse momentum. The PV is chosen as the one giving the smallest three-dimensional impact parameter of the B_c^+ candidate. To avoid biasing L_{xy} , the PV position is recalculated after removing any tracks used in the reconstruction of the B_c^+ meson candidate. To remove poorly reconstructed candidates, $L_{xy}(B_c^+)$ is also required to not exceed 10 mm.
- $L_{xy}(D_s^+) > 0$ mm, where $L_{xy}(D_s^+)$ is the transverse distance between the B_c^+ vertex and the D_s^+ vertex projected onto the direction of the D_s^+ transverse momentum. To remove poorly reconstructed candidates, $L_{xy}(D_s^+)$ is also required to not exceed 10 mm.
- $|d_0^{PV}(B_c^+)/\sigma_{d_0^{PV}}(B_c^+)| < 5$ and $|z_0^{PV}(B_c^+)/\sigma_{z_0^{PV}}(B_c^+)| < 5$, where d_0^{PV} and z_0^{PV} are the transverse and longitudinal impact parameters with respect to the PV, and $\sigma_{d_0^{PV}}(B_c^+)$ and $\sigma_{z_0^{PV}}(B_c^+)$ are their respective uncertainties. The uncertainties are calculated from the covariance matrix of the PV and the covariance matrix of the B_c^+ pseudo-track extracted from the cascade vertex fit. These two requirements are used to ensure that the B_c^+ candidate points back to the PV.
- *p*_T(*B*⁺_c)/∑ *p*_T(trk) > 0.10, where the sum is taken over all tracks originating from the selected PV, including the tracks of the *B*⁺_c candidate. Due to the characteristic hard fragmentation of *b*-quarks, this requirement removes a sizeable fraction of the combinatorial background while having almost no effect on the signal.

 $p_T(B_c) > 15 \text{ GeV}$ $|\eta(B_c)| < 2$ $p_T(B_c) > 15 \text{ GeV}$ $|\eta(B_c)| < 2$ $p_T($

m(K⁺K⁻) within ±7 MeV of m(ϕ)_{PDG}

in case of $D_s^* \rightarrow D_s \pi^0 / \gamma : \pi^0 / \gamma$ escapes detection 3 helicity amplitudes for PS \rightarrow V + V decay: A₀₀, A₊₊, A₋₋



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BDT selection:

$p_{\rm T}$ of the D_s^+ meson candidate, the $L_{xy}(D_s^+)$ variable, and four angular variables:

- $\cos \theta^*(\pi^+)$, where $\theta^*(\pi^+)$ is the angle between the pion momentum in the $K^+K^-\pi^+$ rest frame and the $K^+K^-\pi^+$ combined momentum in the laboratory frame. The signal distribution of $\cos \theta^*(\pi^+)$ is flat before kinematic selection because the pseudoscalar D_s^+ meson decays isotropically, but it increases as $\cos \theta^*(\pi^+)$ approaches +1 for the background events.
- $|\cos^3 \theta'(K^+)|$, where $\theta'(K^+)$ is the angle between one of the kaons and the pion in the K^+K^- rest frame. The decay of the pseudoscalar D_s^+ meson to the ϕ (vector) plus π^+ (pseudoscalar) final state results in the spin of the ϕ meson being aligned perpendicularly to the direction of motion of the ϕ relative to the D_s^+ . Consequently, the distribution of $\cos \theta'(K)$ follows a $\cos^2 \theta'(K)$ shape, implying a uniform distribution for $\cos^3 \theta'(K)$. In contrast, the $\cos \theta'(K)$ distribution of the combinatorial background is approximately uniform and its $\cos^3 \theta'(K)$ distribution peaks at zero.
- $\cos \theta^*(D_s^+)$, where $\theta^*(D_s^+)$ is the angle between the D_s^+ momentum in the B_c^+ rest frame and the B_c^+ flight direction in the laboratory frame. The distribution of $\cos \theta^*(D_s^+)$ is uniform for the decays of pseudoscalar B_c^+ mesons before kinematic selection, while it tends to increase towards negative values for the background.
- $\cos \theta'(\pi^+)$, where $\theta'(\pi^+)$ is the angle between the J/ψ momentum and the pion momentum in the $K^+K^-\pi^+$ rest frame. Its distribution is nearly uniform for the signal processes but peaks towards -1 and +1 for the background.

 $S/\sqrt{S+B}$

efficiency of 81%





Parameter	Value	
$m_{B_c^+}$ [MeV]	6274.8 ± 1.4	agree with PDG
$\sigma_{B_c^+}$ [MeV]	11.5 ± 1.5	agree with MC
$r_{D_{s}^{*+}/D_{s}^{+}}$	1.76 ± 0.22	
$f_{\pm\pm}$	0.70 ± 0.10	
$N^{\mathrm{DS1}}_{B^+_c o J/\psi D^+_s}$	193 ± 20	
$N^{\mathrm{DS2}}_{B^+_c \to J/\psi D^+_s}$	49 ± 10	
$N_{B_c^+ \to J/\psi D_s^{*+}}^{\mathrm{DS1}}$	338 ± 32	
$N_{B_c^+ \to J/\psi D_s^+}^{ m DS1\&2}$	241 ± 28	
$N^{\mathrm{DS1\&2}}_{B^+_c \to J/\psi D^{*+}_s}$	424 ± 46	

2D unbinned fit of two subsets simultaneously : results

Systematic uncertainties

Source	Uncertainty [%]			-	
	$R_{D_s^+/\pi^+}$	$R_{D_s^{*+}/\pi^+}$	$R_{D_s^{*+}/D_s^+}$	$\Gamma_{\pm\pm}/\Gamma$	_
Simulated $p_{\rm T}(B_c^+)$ spectrum	1.5	1.9	0.4	0.1	
Simulated $ \eta(B_c^+) $ spectrum	0.7	0.7	0.1	0.2	
B_c^+ lifetime	0.1	< 0.1	_	_	
D_s^+ lifetime	0.4	0.4	_	_	
Tracking efficiency	1.0	1.0	< 0.1	< 0.1	
Pile-up effects	1.0	1.0	_	_	
$\chi^2/N_{\rm dof}$ cut efficiency	3.2	3.2	_	_	
Impact parameter cuts efficiency	0.2	0.2	_	_	•
BDT cut efficiency	1.3	1.3	-	-	
Trigger efficiency	0.6	0.6	_	-	
Other D_s^+ decay modes	1.6	1.6	_	_	
$B_c^+ \rightarrow J/\psi D_s^{(*)+}$ signal fit:					- ·
D_s^+ signal mass modelling	1.8	0.5	1.3	0.8	
D_s^{*+} signal mass modelling	0.6	1.2	1.7	2.7	
Signal angular modelling	0.4	< 0.1	0.4	0.6	
Background mass modelling	6.0	9.0	3.2	1.0	
Background angular modelling	0.9	1.3	2.1	2.4	
$B_s^0 \to \mu^+ \mu^- \phi$ triggers	0.8	0.5	1.3	4.0	
D_s^{*+} branching fractions	< 0.1	< 0.1	< 0.1	0.7	
$B_c^+ \to J/\psi \pi^+$ signal fit:					
Signal modelling	4.2	4.2	_	_	<u>/</u>
PRD/comb. background modelling	5.8	5.8	_	_	N
CKM-suppr. background modelling	1.0	1.0	_	_	_
MC statistics	1.5	1.5	1.7	1.5	<
Total	10.8	12.6	5.0	5.9	
$\mathcal{B}(D_s^+ \to \phi(K^+K^-)\pi^+)$	5.9	5.9	_	_	

omparison	with previ	ous results,	models and	lighter B-meson de
$R_{D_s^+/\pi^+}$	$R_{D_s^{*+}/\pi^+}$	$R_{D_s^{*+}/D_s^+}$	$\Gamma_{\pm\pm}/\Gamma$	Reference
2.76 ± 0.47	5.33 ± 0.96	1.93 ± 0.26	0.70 ± 0.11	ATLAS Run 2
2.90 ± 0.62 3.8 ± 1.2	- 10.4 ± 3.5	$\begin{array}{c} 2.37 \pm 0.57 \\ 2.8^{+1.2}_{-0.9} \end{array}$	0.52 ± 0.20 0.38 ± 0.24	LHCb Run 1 ATLAS Run 1
2.6	4.5	1.7	_	QCD potential model
1.3	5.2	3.9	_	QCD sum rules
1.29 ± 0.26	5.09 ± 1.02	3.96 ± 0.80	0.46 ± 0.09	CCQM
2.2	_	_	_	BSW
2.06 ± 0.86	_	3.01 ± 1.23	_	LFQM
$3.45_{-0.17}^{+0.49}$	_	$2.54^{+0.07}_{-0.21}$	0.48 ± 0.04	pQCD
3.7832	_		0.410	RIQM
3.257 ± 0.293	_	_	_	FNCM
1.67 ± 0.36	3.49 ± 0.52	2.09 ± 0.52	_	$B^+ \to \bar{D}^{*0} D_s^{(*)+} / \bar{D}^{*0} \pi^+$

 0.48 ± 0.05

 0.94 ± 0.18

 0.396 ± 0.023

 0.429 ± 0.007

 0.4774 ± 0.0034

 2.21 ± 0.35

 1.402 ± 0.083

 1.425 ± 0.065

 6.46 ± 0.60

 7.2 ± 2.1

 2.92 ± 0.42

 $B^0 \to D^{*-} D_s^{(*)+} / D^{*-} \pi^+$

 $B_s^0 \to D_s^{*-} D_s^+ / D_s^{*-} \pi^+$

 $B^+ \rightarrow J/\psi K^{(*)+}$ $B^0 \rightarrow J/\psi K^{(*)0}$

 $B_s^0 \to J/\psi \phi$

Modified Gaussian

$$G_{\text{mod}} \propto \exp\left(-0.5 \times t^{1} + 1/(1+t/2)\right)$$

where $t = |m(J/\psi D_s^+) - m_{B_c^+}|/\sigma_{B_c^+}$

 $B_c^+ \rightarrow J/\psi D_s^+$ signal mass shape modelling effects are tested with alternative models for the $B_c^+ \rightarrow J/\psi D_s^+$ signal $m(J/\psi D_s^+)$ distribution: a double-Gaussian function and a double-sided Crystal Ball function [34–36], fixing the tail parameters to the values extracted from simulation.

Is it really $B_c^+ \rightarrow J/\psi D_s^{(*)+}$?

EPJ C 75 (2016) 1

(not $B_c^{+} \rightarrow \mu^+ \mu^- \varphi \pi^+, \ \mu^+ \mu^- D_s^+, \ J/\psi \varphi \pi^+$)



Cascade fit w/o mass constraints

events with 5.9 < $m(J/\psi D_{s^{+}})$ < 6.4 GeV

 $N(J/\psi) = 568 \pm 28$

a lot of J/ψ in background

 $N(D_{s}^{+}) = 175 \pm 36$

non-significantly above $N(B_c^+ \rightarrow J/\psi D_s^+) + N(B_c^+ \rightarrow J/\psi D_s^{*+})$

Observation of excited $B_c (\rightarrow B_c \pi^+ \pi^-)$



ATLAS, PRL 113 (2014) 212004

Large B_c family is expected although only ground state has been known until today

 $B_c^{\pm}(2S)$ 6835–6917 MeV $2S/1S \simeq 0.6$ $2^1S_0 \rightarrow 1^1S_0 + \pi\pi$



 $p_{\tau}(\pi) > 4 \text{ GeV}, m(J/\psi)$ constrained to PDG



Observation of excited $B_c (\rightarrow B_c \pi^+ \pi^-)$ ATLAS, PRL 113 (2014) 212004

 $p_{\tau}(\pi^{\pm}) > 400 \text{ MeV}, m(J/\psi)$ constrained to PDG



Significance from $\Delta \ln L$ of pseudo-experiments: 5.4 σ (local)

5.2 σ ("look elsewhere")

Q = 288.3 ± 3.5 ± 4.1 MeV M = 6842 ± 4 ± 5 MeV

Both mass value and decay mode agree with expectations for $B_c^{\pm}(2S)$



Charmonium production

Non-prompt (from B decays) – probes open b quark production, g fragmentation and B-decay kinematics FONLL, matched NLO+NLL ("massive" NLO + resummation) GM-VFNS ("massless" NLO + mass-dependent terms)

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Prompt (not from B decays) – probes specific mechanisms of $Q\bar{Q}$ system production and transformation to a meson



NRQCD: Color Singlet (CS) and Color Octet (CO) terms. Long-distance matrix elements (LDME) determined from experimental data. Color Singlet Model (CSM) – only CS diagrams. Color Evaporation Model (CEM) – only one LDME.

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Charmonium production at 13 TeV with 139 fb⁻¹

ATLAS-CONF-2019-047

Uses a single-muon trigger, with threshold at 50 GeV, un-prescaled on the full integrated luminosity of Run II, 139 fb⁻¹

 p_T range covered: 60-360 GeV for J/ ψ in 11 bins (60-140 GeV for ψ (2S))

Rapidity range |y| < 2 covered in three bins

Yields for J/ ψ and ψ (2S), prompt and non-prompt (from B decays), determined using 2D fit (mass and "pseudo-proper" lifetime)

 $\boldsymbol{\tau} = \frac{m L_{xy}}{c P_T}$



15

Charmonium non-prompt fractions



Charmonium non-prompt x-sections



0.5⊨

60

70 80 90 100

New fragmentation tuning? Fixing of technical FONLL problems at high p_T?

300

 $p_{_{T}}(\mu\mu)$ [GeV]

200

Charmonium prompt x-sections



Waiting NRQCD predictions for high- p_T charmonium production

Measurement of J/ψ production in association with a W^{\pm} boson with pp data at 8 TeV A

ATLAS, JHEP 01 (2020) 95





 $\Lambda_b \to J/\psi \; p \; K^{\text{-}}$ signal is seen on the top of

- large combinatorial background
- very large $B \rightarrow J/\psi K^+ \pi^-$ contribution
- large $B_s \rightarrow J/\psi K^+ K^-$ contribution
- tails from small $B \rightarrow J/\psi \pi^+ \pi^-$ and $B_s \rightarrow J/\psi \pi^+ \pi^-$ contributions

P_c⁺ at 7 - 8 TeV



1010±140 direct $\Lambda_b \rightarrow J/\psi, p, K$

 $\Lambda_b \rightarrow J/\psi, p, K$ decays analysis: 2 pentaquark hypothesis



 $\chi^2/N_{dof} = 49.0/43 \text{ (p-value} = 0.25)$

P_c signal parameters and yields from fit:

	ĸ	μ ^μ
	[J/W H
Ab	Pc	p

Parameter	Value	LHCb value
$N(P_{c1})$	$400^{+130}_{-140}(\text{stat})^{+110}_{-100}(\text{syst})$	-
$N(P_{c2})$	$150^{+170}_{-100}(\text{stat})^{+50}_{-90}(\text{syst})$	-
$N(P_{c1} + P_{c2})$	$540^{+80}_{-70}(\text{stat})^{+70}_{-80}(\text{syst})$	-
$\Delta \phi$	$2.8^{+1.0}_{-1.6}(\text{stat})^{+0.2}_{-0.1}(\text{syst})$ rad	<u></u>
$m(P_{c1})$	4282 ⁺³³ ₋₂₆ (stat) ⁺²⁸ ₋₇ (syst) MeV	$4380 \pm 8 \pm 29$ MeV
$\Gamma(P_{c1})$	140^{+77}_{-50} (stat) $^{+41}_{-33}$ (syst) MeV	$205 \pm 18 \pm 86$ MeV
$m(P_{c2})$	4449 ⁺²⁰ ₋₂₉ (stat) ⁺¹⁸ ₋₁₀ (syst) MeV	$4449.8 \pm 1.7 \pm 2.5$ MeV
$\Gamma(P_{c2})$	51^{+59}_{-48} (stat) $^{+14}_{-46}$ (syst) MeV	$39 \pm 5 \pm 19$ MeV

 $\Lambda_h \rightarrow J/\psi, p, K$ decays analysis: 4 pentaquark hypothesis



Similar fits (no interference, Breit-Wigner amplitudes) has been performed on our data with masses, widths and relative yields of narrow states fixed to LHCb values. Parameters of P_c (4380) kept free.

ATLAS data is consistent with LHCb Run II results.

No pentaquark fits: extended *A** decay model



Projection of 2D M(J/ ψ ,p) vs M(J/ ψ ,K) + 1D M(p,K) fit w/o pentaquarks using extended Λ^* decay model (left)

Result of 1D χ 2 M(J/ ψ ,p) fit with the same model (right): χ ²/NDF = 42.0/23 **p-val = 9.1 x 10**⁻³

This model shows a 'border-line agreement' with data.

J/ψ

 Λ_h

Measurements of CP violation with $B_s \rightarrow J/\psi \phi$

The time evolution of B_s meson mixing is characterized by

- \checkmark the mass difference Δm_s of the heavy (B_H) and light (B_L) mass eigenstates
- ✓ the CP-violating mixing phase φ_s
- ✓ the width difference of $\Delta\Gamma_s = \Delta\Gamma_L \Delta\Gamma_H$

Interference between the B_s decays amplitudes to the CP eigenstates $J/\psi \varphi$ or via mixing gives rise to a measurable CP violating phase φ_s



New Physics could modify φ_s and $\Delta \Gamma_s / \Delta m_s$ If new particles contributes to box diagrams

$$PS \rightarrow VV$$
 decay gives orbital angular $L = 0$ or 2 are CP -even
momentum $L = 0, 1$ or 2 $L = 1$ is CP -odd 55

Measurement of the *CP*-violating phase ϕ_s in $B_s^0 \rightarrow J/\psi \phi$ decays in ATLAS at 13 TeV

Eur. Phys. J. C 81 (2021) 342 DOI: 10.1140/epjc/s10052-021-09011-0



angular analysis

Measurement of the *CP*-violating phase ϕ_s in $B_s^0 \rightarrow J/\psi \phi$ decays in ATLAS at 13 TeV



Eur. Phys. J. C 81 (2021) 342 DOI: 10.1140/epjc/s10052-021-09011-0

ϕ_s	=	-0.087	± 0.036	$(stat.) \pm 0$	0.021	(syst.) 1	rad
Γ_s	=	0.0657	± 0.0043	$(stat.) \pm 0$	0.0037	(syst.) j	ps ⁻¹
Γ_s	=	0.6703	± 0.0014	$(stat.) \pm 0$	0.0018	(syst.) j	ps ⁻¹

13 TeV: Data 2015 - 2017

Data 2018 to be included in next publication

Измерение СР-нарушающей фазы ϕ_s и $\Delta\Gamma_s$ на 100 fb⁻¹

В согласии с LHCb, CMS и CM по ϕ_s рассогласование по $\Delta\Gamma_s$ и Γ_s

New combination after full Run-2 results publication

Measurement of the *CP*-violating phase ϕ_s in $B_s^0 \rightarrow J/\psi \phi$ decays in ATLAS at 13 TeV

Eur. Phys. J. C 81 (2021) 342 DOI: 10.1140/epjc/s10052-021-09011-0



Studies of methodical differences New combination after full Run-2 results publication

Problems with lifetime measurements? ATLAS is working on high precision measurement of $B^o \to J/\psi K^{*0}$ lifetime

$B^0 \rightarrow K^{*0} \mu^+ \mu^-$ (FCNC) ATLAS, JHEP 10 (2018) 047



polarised K*'s S_i – angular coefficients

Full angular analysis Performed in bins of $q^2 = m^2(\mu^+\mu^-)$







Angular analysis on B \rightarrow K^{*}µµ at 8 TeV



Results are compatible with theoretical calculations & fits

Deviations of about 2.5 σ (2.7 σ) from DHMV in P'₄(P'₅) in [4,6] GeV²

$b \rightarrow s\ell^+\ell^-$ transitions



 $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ (FCNC) JHEP 10 (2018) 047



Angular observables in $B \rightarrow K^* \mu \mu$ show about 3.4 σ discrepancy



More data are needed!

Angular analysis of $B^+ \rightarrow K^*(892)^+ \mu^+ \mu^-$ at CMS

JHEP 04 (2021) 124



- The first angular analysis of the $B^+ \rightarrow K^*(892)^+ \mu^+ \mu^-$ decay is performed;
- In three bins of the dimuon invariant mass squared (q²), a 3D fits are performed;
- The muon forward/backward asymmetry (A_{FB}) and $K^*(892)^+$ longitudinal fraction (F_L) are consistent with SM predictions;

q^2 (GeV ²)	Y_S	$A_{ m FB}$	$F_{\rm L}$
1 - 8.68	22.1 ± 8.1	$-0.14^{+0.32}_{-0.35}\pm0.17$	$0.60^{+0.31}_{-0.25}\pm0.13$
10.09 - 12.86	25.9 ± 6.3	$0.09^{+0.16}_{-0.11}\pm0.04$	$0.88^{+0.10}_{-0.13}\pm0.05$
14.18 – 19	45.1 ± 8.0	$0.33^{+0.11}_{-0.07}\pm0.05$	$0.55^{+0.13}_{-0.10}\pm0.06$

• ATLAS earlier performed similar analysis of $B^0 \rightarrow K^* \mu^+ \mu^-$ decays

JHEP 10 (2018) 047

Observation of new bottom-strange baryon $\Xi_b^-(6100)$ in $\Xi_b^-\pi^+\pi^-$ channel at CMS



Topologies with different Ξ_b^- decay channels



• A narrow resonance $\Xi_b^-(6100)$ is observed at a $\Xi_b^-\pi^+\pi^-$ invariant mass of:

 $M = 6100.3 \pm 0.2(stat) \pm 0.1(syst) \pm 0.6(\bar{z}_{b})$

- Results are consistent between fully reconstructed channels and partially reconstructed channel with $\Sigma^0 \rightarrow \Lambda \gamma$
- Upper limit is set 1.9MeV on a natural width of the new state (95% CL)
- New state is consistent with orbitally excited Ξ_b baryon with spin/parity of 3/2⁻.

X(3872) production in different collision systems

First evidence using 1.7 nb⁻¹ of PbPb collisions data (2018) at CMS at $\sqrt{s_{NN}} = 5.02$ TeV per nucleon pair <u>PRL 128 (2022) 032001</u>

UML fit to extract signal yields for $\psi(2S)$ and X(3872) Final state: $J/\psi(\rightarrow \mu^+\mu^-) \pi^+\pi^-$ Significance for inclusive X(3872): 4.2 σ

Prompt fraction estimated with MC studies Yields corrected by acceptance and overall efficiency

Ratio of corrected yields for prompt production in PbPb collisions Q^{pp}:

- compatible with 1 (within 1 σ)
- compatible with $Q^{pp} \approx 0.1$ (within 2σ)

Much larger data sample expected in Run-3 at LHC in order to improve the measurement and understand the internal structure of X(3872) and the differences of its production mechanism w.r.t. $\psi(2S)$ kinematical range: 15 < p_T < 50 GeV/c, lyl < 1.6





Lepton Flavour Universality (Violation?)

$$R_{K^{(*)}} = \frac{\Gamma(B \to K^{(*)}\mu^+\mu^-)}{\Gamma(B \to K^{(*)}e^+e^-)} \qquad \mathcal{R}_{K^{*0}} = \frac{\mathcal{B}(B^0 \to K^{*0}\mu^+\mu^-)}{\mathcal{B}(B^0 \to K^{*0}J/\psi(\to \mu^+\mu^-))} \bigg/ \frac{\mathcal{B}(B^0 \to K^{*0}e^+e^-)}{\mathcal{B}(B^0 \to K^{*0}J/\psi(\to e^+e^-))}$$

ATLAS potentially can do $R(K), R(\phi), R(pK) = BR(\Lambda_b \rightarrow$ $pK\mu\mu$)/BR($\Lambda_b \rightarrow pKee$)

In 2018, a di-electron high-level trigger implemented and being analysed now





Phys. Rev. Lett. 122 (2019) 191801 2.5σ from the SM LHCb



Aiming at R(K*) measurement

Run-3 data will add statistic

Lepton Flavour Universality (Violation?)

$$R_{D^*} = \frac{\Gamma(\overline{B}^0 \to D^{*+} \tau^- \overline{\nu_{\tau}})}{\Gamma(\overline{B}^0 \to D^{*+} \mu^- \overline{\nu_{\mu}})}$$

 \sim

ATLAS added dedicated trigger branches for Run 3:



Prospects



Shutdown/Technical stop Protons physics Ions Commissioning with bear

Ions Commissioning with beam Hardware commissioning/magnet training

Run 3: ~300 fb-1 HL-LHC: ~3000 fb-1



Last updated: January 2022