SOLITON CONFIGURATIONS AND GROUND STATES IN MAXIMAL GAUGED SUPERGRAVITY

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1 Introduction



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- 2 The model
 - Explicit solutions
 - Thermodynamics
- BPS solutions



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3 Conclusions





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Solitons: special role in classical physics as well as in quantum and string theory, determining a richer structure of the full non-perturbative regime:

- originally used as "bounce solutions" to discuss the possible instability of the pure Kaluza-Klein vacuum ground state;
- generalizations of these soliton solutions have been also considered in the analysis of the semiclassical stability of non-susy AdS gravity;
- soliton configurations can turn out to be the lowest energy solution with chosen boundary conditions, leading to a new kind of positive energy conjecture;





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 BPS configurations preserving some of the supercharges can be obtained analysing the explicit form of the Killing spinors equations.





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- lacktriangledown The scalar fields of the theory can be characterized by means of the geometry of the chosen non-linear σ -model.
- We consider an explicit solutions in the T^3 model, the latter resulting in a single dilaton truncations of the maximal SO(8) gauged supergravity in D=4.





We restrict to purely magnetic solutions. The action has the explicit form:

$$\mathscr{S} = \frac{1}{8\pi G} \int d^4x \; \sqrt{-g} \left(\frac{R}{2} - \frac{1}{2} \left(\partial \varphi \right)^2 + \frac{3}{L^2} \; \text{cosh} \left(\sqrt{\frac{2}{3}} \, \varphi \right) - \frac{1}{4} \, e^{3\sqrt{\frac{2}{3}} \, \varphi} \left(F^1 \right)^2 - \frac{1}{4} e^{-\sqrt{\frac{2}{3}} \, \varphi} \left(F^2 \right)^2 \right).$$



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- in the model we consider there are two Wilson lines,

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and there is a one-parameter family of values of the Wilson lines which give supersymmetric solitons;





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$$\begin{split} & \varphi \; = \; \pm \ell^{-1} \, \text{In}(x) \,, \qquad F^{\Lambda}_{\mu\nu}(x,\Gamma^{\Lambda}) \,, \\ & ds^2 \; = \; \Upsilon(x) \left(L^2 \, dt^2 - \frac{\eta^2}{f(x)} \, dx^2 - f(x) \, d\psi^2 - L^2 \, dz^2 \right) \,; \end{split}$$

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- for special boundary conditions, can be found both susy and non-susy solutions
 - \implies new kind of degeneracy of supersymmetric solutions;
 - surprisingly, there is a family of non-susy solutions of lower energy and free energy than the susy ones.

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$$\Delta^{-1} = \left| \frac{1}{4\pi\eta} \frac{\mathrm{df}}{\mathrm{dx}} \right|_{x=x_0} ;$$





Solutions with non-zero charges have net magnetic fluxes at infinit

$$\begin{split} &\Phi_\text{M}^1 = \int F^1 = \oint A^1 = Q_1 \, \Delta \left(1 - \kappa_0^{-2}\right) \equiv 2\pi L \, \psi_1 \text{,} \\ &\Phi_\text{M}^2 = \int F^2 = \oint A^2 = Q_2 \, \Delta \left(1 - \kappa_0^2\right) \equiv 2\pi L \, \psi_2 \, . \end{split} \label{eq:phi_M}$$



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The scalar field induces a vev of an operator in the dual theory,

$$\left\langle \mathfrak{O}\right\rangle =\varphi_{0}=\pm\frac{\sqrt{6}}{2}\;\frac{\pi\,x_{0}\left|\psi_{1}^{2}\left(1+2\,x_{0}^{2}\right)-\psi_{2}^{2}\right|}{\Delta}\;.\label{eq:eq:energy_energy}$$



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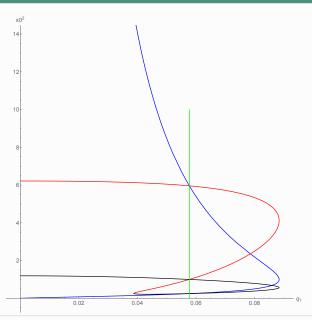
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- for the same fixed charge boundary conditions, surprisingly a family of non-susy solutions of lower energy and free energy than the supersymmetric ones can be found.

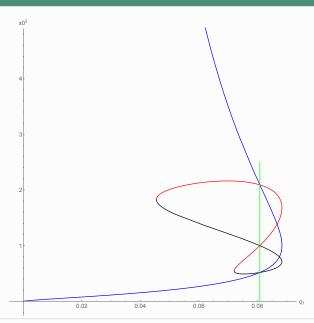






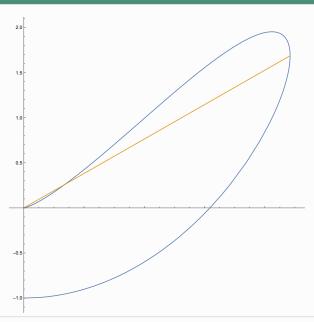
















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- The new solutions require a more in-depth study of the degeneracy of the susy configurations in the presence of generic boundary conditions.
- One branch of susy solutions has higher energy than a non-susy one with the same boundary conditions.



Thank you for listening!