



# Model analysis of transverse momentum fluctuations in NICA and SPS energy range

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# Motivation



Figure 1: QCD diagram of phase transition.



- transition from negative to positive correlations: good description with help of multipomeron model [N. Armesto et al., Phys. of Atom. Nucl., 71. 2087-2095 (2008)]
- limit theoretical model (PYTHIA, Herwig++ and etc.)

#### Models

- EPOS (Energy-conserving quantum mechanical multiple scattering approach, based on Partons (parton ladders), Off-shell remnants, and Splitting of parton ladders). This model [K. Werneret al., Phys.Rev. C74, 044902 (2006)] takes into account the multiple scattering approach based on partons and pomerons (parton ladders). The basis is a string model.
- SMASH (Simulating Many Accelerated Strongly-interacting Hadrons) is a relativistic hadronic transport approach [J. Weil et al., arXiv:1606.06642 [nucl-th] (2017)]. Include all well-established hadrons up to a mass of ~ 2 GeV as degrees of freedom.
- PHSD (*Parton-Hadron-String Dynamics*)[E. Bratkovskaya et al., arXiv:1908.00451 [nucl-th] (2019)] is a microscopic off-shell transport approach for the description of strongly interacting hadronic and partonic matter in and out-of equilibrium.
- UrQMD (Ultra-relativistic Quantum Molecular Dynamics)[M. Bleicher et al., arXiv:hep-ph/9009407 [hep-ph] (1999).] is a microscopic model used to simulate (ultra)relativistic heavy ion collisons in the energy range from Bevalac and SIS up to AGS, SPS and RHIC.

These models have been applied to nucleus-nucleus collisions from low Super-Proton-Synchrotron (SPS) to Large-Hadron-Collider (LHC) energies in order to explore the space-time regions of 'partonic matter'.

$$\Sigma[p_t, N] = \frac{1}{\langle N \rangle \omega [p_t]} [\langle N \rangle \omega [P_t] + \langle P_t \rangle \omega [N]]$$
(1)

$$\Delta[p_t, N] = \frac{1}{\langle N \rangle \omega[p_t]} \left[ \langle N \rangle \omega[P_t] + \langle P_t \rangle \omega[N] - 2 \left( \langle P_T N \rangle - \langle P_T \rangle \langle N \rangle \right) \right]$$
(2)

M. Gorenstein, M. Gazdzicki, Phys. Rev. C 84, 014904 (2011)

where  $P_T = \sum_{i=1}^{N} p_{T_i}$  and  $\omega[p_T]$  is the scaled variance of the inclusive  $p_T$  spectrum.  $\Delta[p_T, N] = \sum[p_T, N] = 1$  - for independent particle production model,  $\Delta[p_T, N] = \sum[p_T, N] = 0$  in the

absence of fluctuations.

Also, another strongly intensive quantity is introduced for analysis [M. Cody, S. Gavin, B.Koch et al., arXiv:2110.04884 [nucl-th]]:

$$\langle N \rangle D[p_t, N] = \frac{1}{\langle N \rangle} \left[ \left( \langle P_T N \rangle - \langle P_T \rangle \langle N \rangle \right) - \langle P_t \rangle \omega [N] \right]$$
(3)

#### Definitions and observables: cumulants

The *n*-particle  $p_T$  correlator in one event is defined as

$$C_{n} = \frac{\sum_{i_{1} \neq \dots \neq i_{n}} \omega_{i_{1}} \dots \omega_{i_{n}} (p_{\tau, i_{1}} - \langle \langle p_{\tau} \rangle \rangle) \dots (p_{\tau, i_{n}} - \langle \langle p_{\tau} \rangle \rangle)}{\sum_{i_{1} \neq \dots \neq i_{n}} \omega_{i_{1}} \dots \omega_{i_{n}}}$$
(4)

$$C_2 = \frac{\overline{p}_{11}^2 - \overline{p}_{22}}{1 - \tau_1} \tag{5}$$

$$C_{3} = \frac{\overline{p}_{11}^{3} - 3\overline{p}_{22}\overline{p}_{11} + 2\overline{p}_{33}}{1 - 3\tau_{1} + 2\tau_{2}} \tag{6}$$

$$K_2 = \frac{\langle C_2 \rangle}{\langle \langle \rho_T \rangle \rangle^2},\tag{7}$$

$$K_{3} = \frac{\langle C_{3} \rangle}{\langle \langle p_{T} \rangle \rangle^{3}},\tag{8}$$

Bhatta S. et. al. Phys. Rev. C 105, 024904

## Results for strongly intensive observables



Figure 3: Dependence of  $\Sigma[p_T, N]$  on beam energy for proton-proton collisions.

Figure 4: Dependence of  $\Delta[p_T, N]$  on beam energy for proton-proton collisions.

- The resonances-to-strings transition in models produce a "wave", which is observed on the figures of the Σ[p<sub>T</sub>, N] and Δ[p<sub>T</sub>, N] on energy collisions
- For various models, this transition from resonances to strings occurs with various energies

### Results for strongly intensive observables



Figure 5: Energy dependence of  $\langle N \rangle D[p_t, N]$ .

In our calculations, it was revealed  $D \neq 0$  that was predicted by the model PYTHIA/Angantyr simulations of proton-proton (p + p) and nucleus-nucleus (A + A) collisions [M. Cody et al., arXiv:2110.04884 [nucl-th]].

#### Results for second and third order cumulants



Figure 6: Dependence of the second-order cumulant for the transverse momentum in proton-proton collisions.

Figure 7: Dependence of the third-order cumulant for the transverse momentum in proton-proton collisions.

- In HIJING models  $\kappa_2$  and  $\kappa_3$  have positive means at high energies [S. Bhatta et al., Phys. Rev. C 105, 024904].
- At low energies  $\kappa_2$  has only negative values,  $\kappa_3$  has both positive (SMASH, PHSD) and negative(UrQMD, EPOS) values

The key findings are:

• significant discrepancies between predictions of EPOS, SMASH, PHSD and UrQMD models are observed indicating that data on p + p collisions from the NICA experiment, which will be obtained in the future, will limit the prediction of this models, as well as clarify the results obtained.

• in the figures of the dependence of  $\Sigma[p_T, N]$ ,  $\Delta[p_T, N]$ ,  $\langle N \rangle D[p_T, N]$  on energy appear a "wave" that is caused by the resonance-to-string transition.

• non-trivial  $p_T$  cumulants collision energy dependence predicted by the models for a 'baseline' p + p reaction emphasizes difficulties in interpreting future results for A + A collisions and requires further investigations.

• At low energies  $\kappa_2$  has only negative values for all models,  $\kappa_3$  has both positive (SMASH, PHSD) and negative(UrQMD, EPOS) values. This means that future NICA experiments will make an important contribution to the study of these quantities.

# Thank you for your attention!



Figure 8: Dependence on energy of the  $p_t - N$  coefficient correlation in proton-proton collisions for EPOS model.

Figure 9: Dependence on energy of the  $p_t - N$  coefficient correlation in proton-proton collisions for SMASH model.