

W-Ai measurements

with low $\langle mu \rangle$ data



Daniil Ponomarenko

10-March-2023

Plan of the talk

- Motivation
- Brief analysis overview
- Updates on MJ estimation



ATLAS Note

ANA-STDM-2020-07-INT1

3rd February 2023



Draft version 0.4

1

2

3

4

**Measurement of angular coefficients in $W \rightarrow \ell \nu$
events in low- μ pp collisions at $\sqrt{s} = 13$ TeV with the
ATLAS detector**

Link to CDS: <https://cds.cern.ch/record/2824052> (last modified 2022-09-30).
Latest version of supporting note is in attachments to the agenda.

Motivation

W mass measurement

$$M_W = 80369.5 \pm 6.8 \text{ (stat)} \pm 10.6 \text{ (exp.syst.)} \pm 13.6 \text{ (model.syst.) MeV}$$

Limited by data

Experimental: modelling of lepton momentum measurement and hadronic recoil

Theoretical: understanding of vector boson production and decay

Combined categories	Value [MeV]	Stat. Unc.	Muon Unc.	Elec. Unc.	Recoil Unc.	Bckg. Unc.	QCD Unc.	EW Unc.	PDF Unc.	Total Unc.	χ^2/dof of Comb.
$m_T-p_T^\ell, W^\pm, e-\mu$	80369.5	6.8	6.6	6.4	2.9	4.5	8.3	5.5	9.2	18.5	29/27

- Larger source of modelling uncertainty from PDFs
- Total uncertainty on ATLAS W mass measurement $\sim 19\text{MeV}$ still larger than 8MeV from electroweak fit

Why measure A_i in W events?

physics modelling

Data sample	7 TeV, $\mu \sim 9$	
Luminosity	4.5 fb ⁻¹	Improved
Nb. of candidates	$\sim 15 \times 10^6$	
Most sensitive dist.	p_T^ℓ	
Physics Modelling Unc.		
EW	5 → latest MC gen.	2
QCD (p_T^W)	6 → WpT meas.	< 3
QCD (A_i)	6 → data input	< 3
PDFs	9 → PDF profiling	6

- **W- A_i analysis (+ x-section):**
 - A_i : Stringent test of QCD & physics modelling!
 - X-section measurement as function of the Y .
 - All previous measurements at the LHC are done with lepton η .
 - Charge asymmetry measurement: Input for PDF fits
- **Full set of A_i have never been measured before for the W boson!**
 - Only A_2 and A_3 were measured by CDF ([Phys. Rev. D 73, 052002](#))

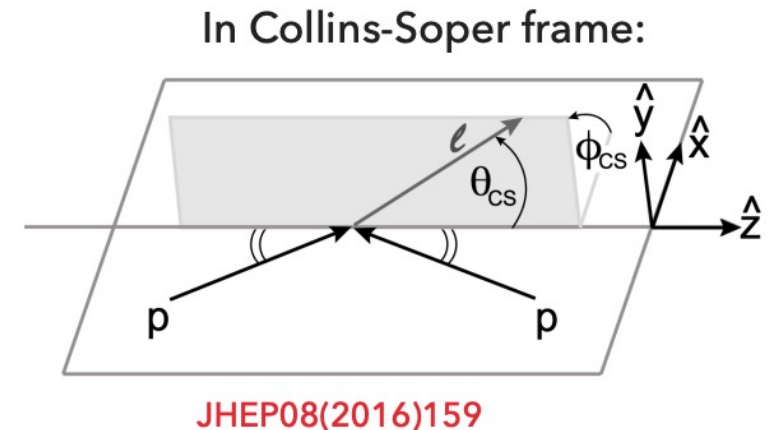
Vector boson production and decay

- Drell Yan $pp \rightarrow V \rightarrow l\bar{l}$ cross section can be factorised into an **unpolarised** cross-section and **8 angular coefficients** (A_i).

$$\frac{d\sigma}{dp_1 dp_2} = \left[\frac{d\sigma(m)}{dm} \right] \left[\frac{d\sigma(y)}{dy} \right] \left[\frac{d\sigma(p_T, y)}{dp_T dy} \left(\frac{d\sigma(y)}{dy} \right)^{-1} \right] \left[(1 + \cos^2 \theta) + \sum_{i=0}^7 A_i(p_T, y) P_i(\cos \theta, \phi) \right]$$

Breit Wigner
Parton shower
Fixed order pQCD

- **The angular coefficients:**
 - A_0, A_1, A_2 sensitive to vector boson **polarisation**
 - $A_0 - A_2 = 0$ but violated due to higher order QCD effects (e.g. multi-gluon emission..)
 - A_3, A_4 sensitive to **vector and axial couplings** of the boson
 - A_4 directly related to A_{FB} , probes electroweak mixing
 - A_5, A_6, A_7 non-zero only at $\mathcal{O}(\alpha_s^2)$
- QCD production dynamics fully contained in A_i coefficients, while decay kinematics fully contained in coupled angular polynomials



Analysis overview

W Ai analysis overview

- Supporting note: [ANA-STDM-2020-07-INT1](#)
- Dataset:
 - ATLAS 13 TeV low pileup datasets 335.18 pb^{-1}
 - Statistically limited by small datasets but provides hadronic recoil resolution needed
- Able to measure 1D in p_T^W and separately in y^W
- We focus on the larger coefficients A_0, A_2, A_3, A_4
 - Other coefficients are also get measured, but the analysis is far from having sensitivity to A_5, A_6, A_7
- Channels:
 - $W^- \rightarrow e^- \nu, W^+ \rightarrow e^+ \nu, W^- \rightarrow \mu^- \nu, W^+ \rightarrow \mu^+ \nu$
- Looking forward for EB request

Event selection

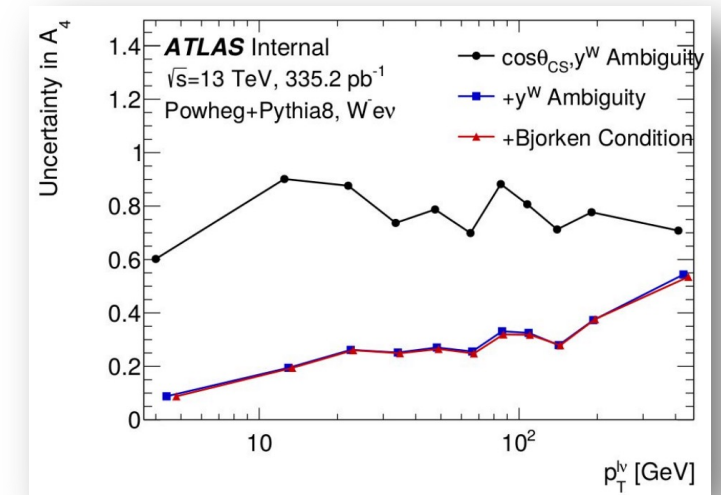
Electron	Muon
Tight ID	Medium ID
$p_T > 25 \text{ GeV}$	
$ \eta < 2.4$	
$1.37 < \eta < 1.52$	
$ptvarcone20/p_T < 0.1$	
$topoetcone20/p_T < 0.05$	
$ d_0 \text{ sig} < 5$	$ d_0 \text{ sig} < 3$
$\Delta z_0 \sin\theta < 0.5$	

- We use signal region selection based on pTW analysis with optimizations for better A_i sensitivity
 - Use same physics objects calibrations
 - + estimate and apply corrections specific for Wai analysis
- Relaxed kinematic cuts:
 - No $E_T^{miss} > 25 \text{ GeV}$ cut
 - No $m_T^W > 50 \text{ GeV}$ cut
- Use *TightLH* working point for electrons
- Tighter isolation selection:
 - $ptvarcone20/pt < 0.1$ && $topoetcone20/pt < 0.05$

Analysis

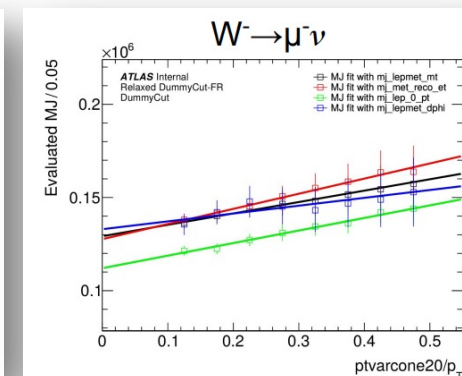
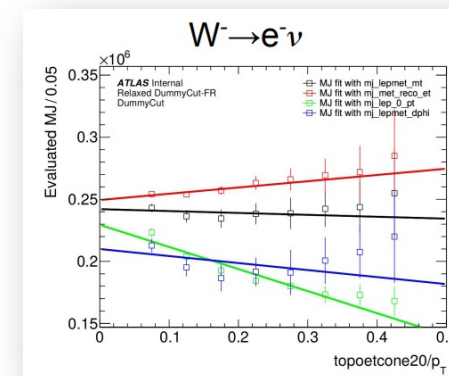
Angular definitions require fully constructed neutrino (will cover in next slides):

- Using Hadronic recoil solve for Neutrino p_z with mass constrain.
- Once the mass constraint is chosen φ_{CS} is solved while the 2 solutions corresponds to a sign ambiguity in $\cos\theta_{CS}$.
- Discovered that adding sign ambiguity only in y^W , not in $\cos\theta_{CS}$ drives our A_4 sensitivity.

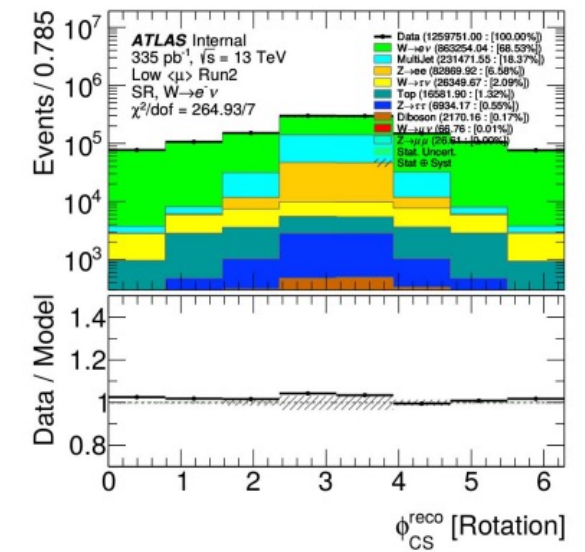
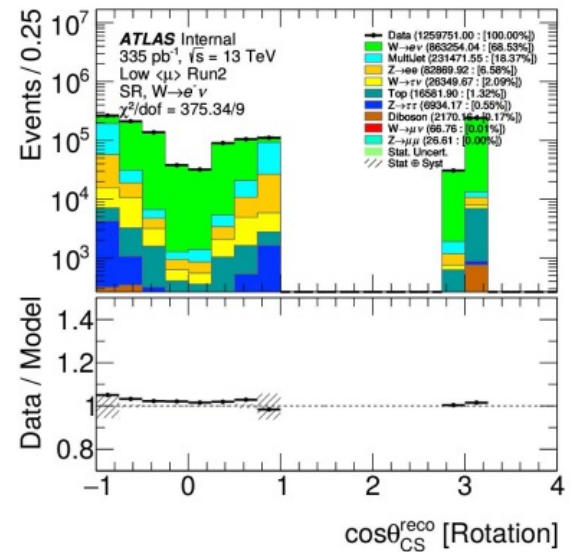
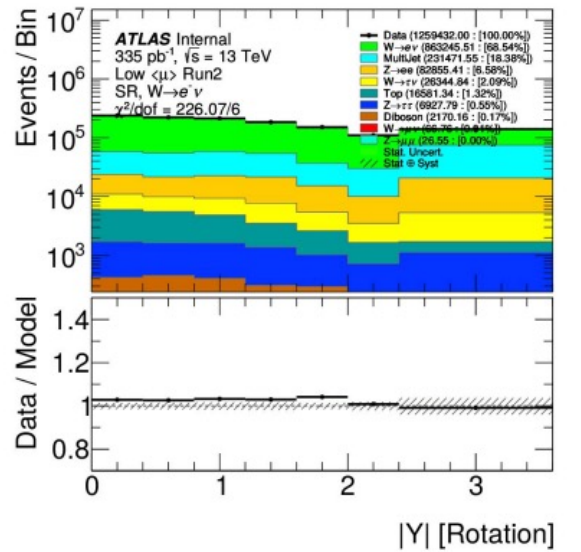
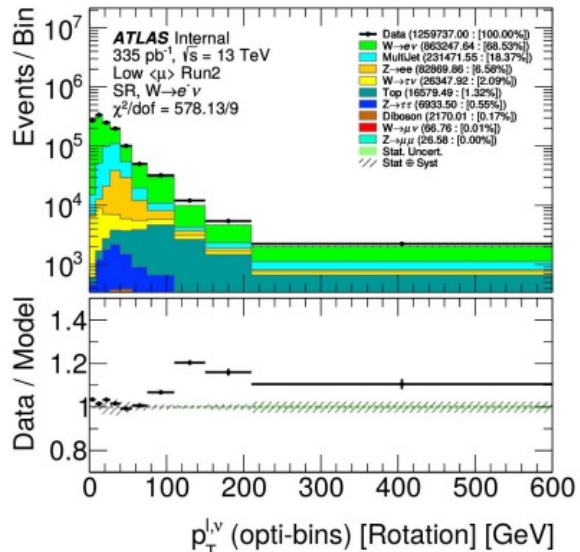
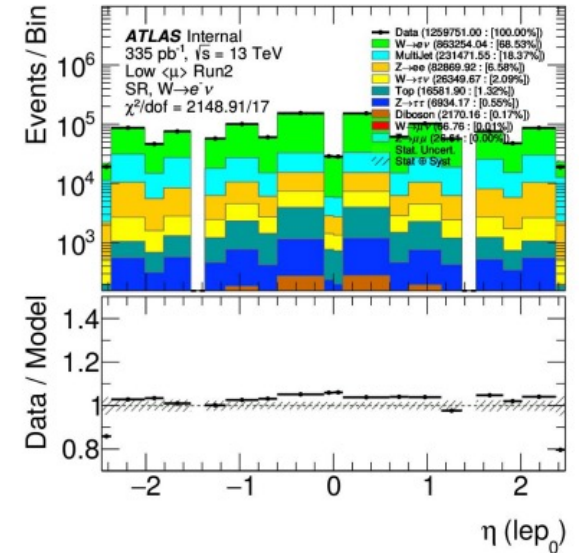
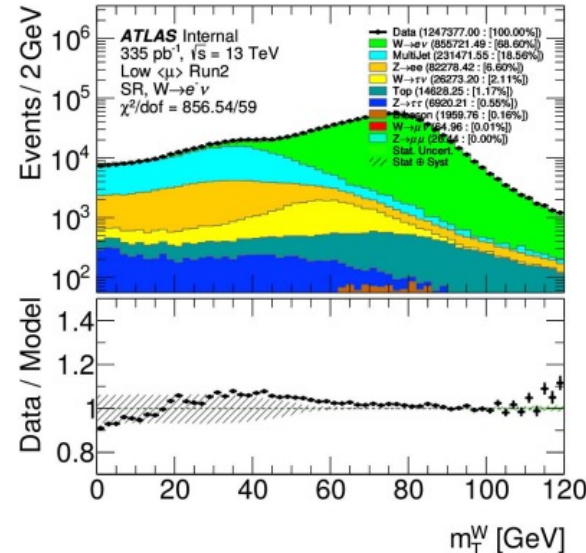
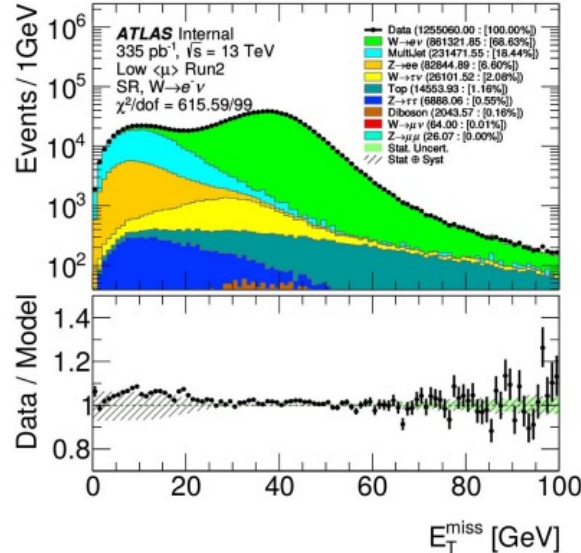
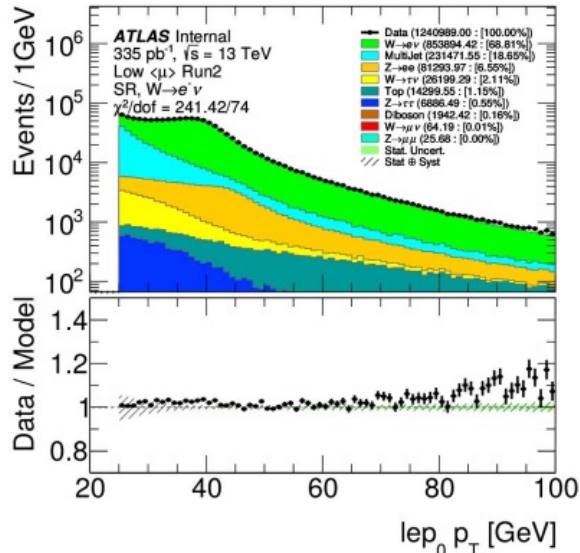


MJ shape estimation follows procedure used by pTW analysis ([slides in backup](#)):

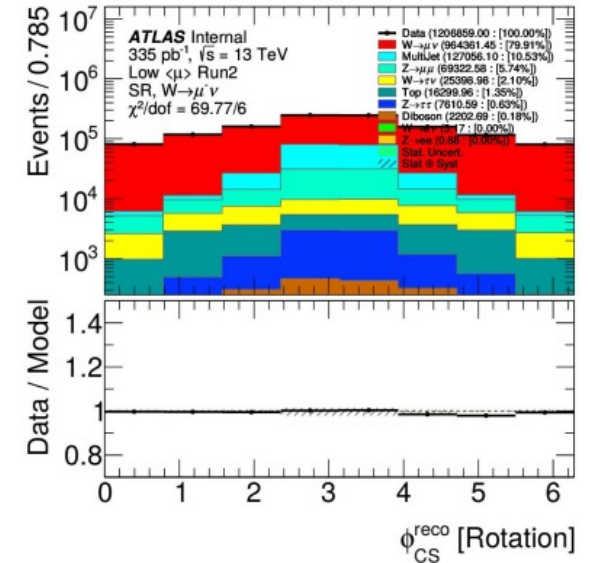
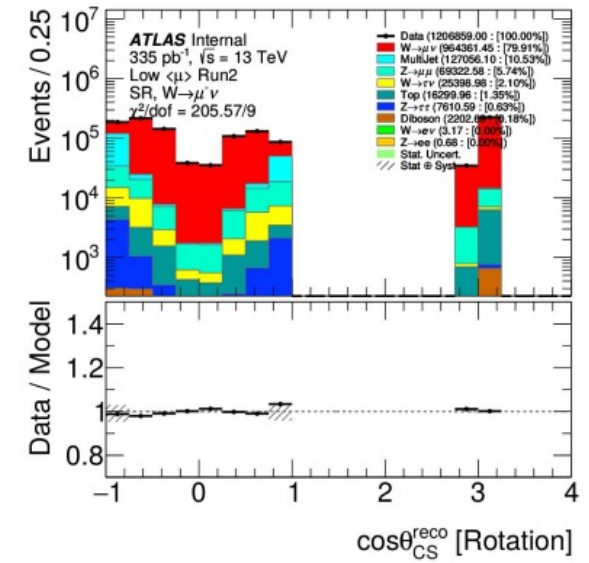
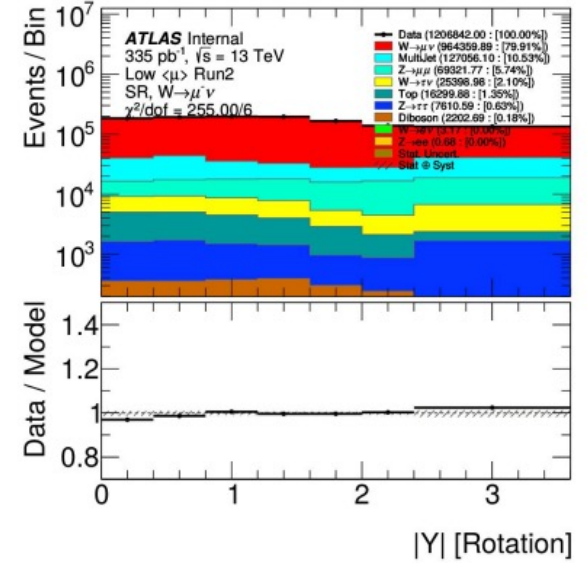
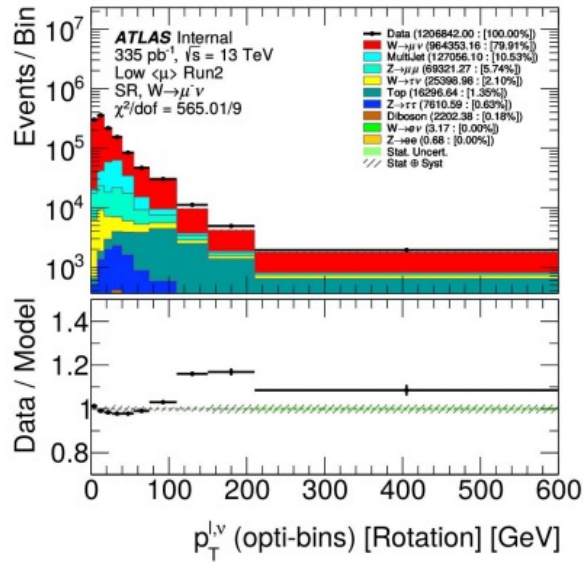
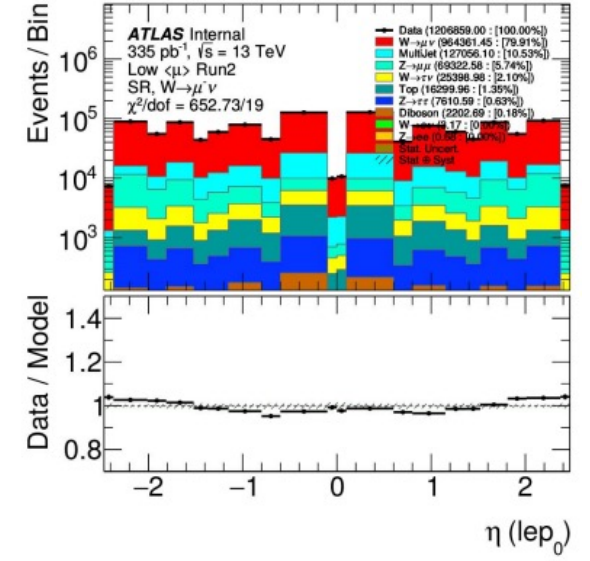
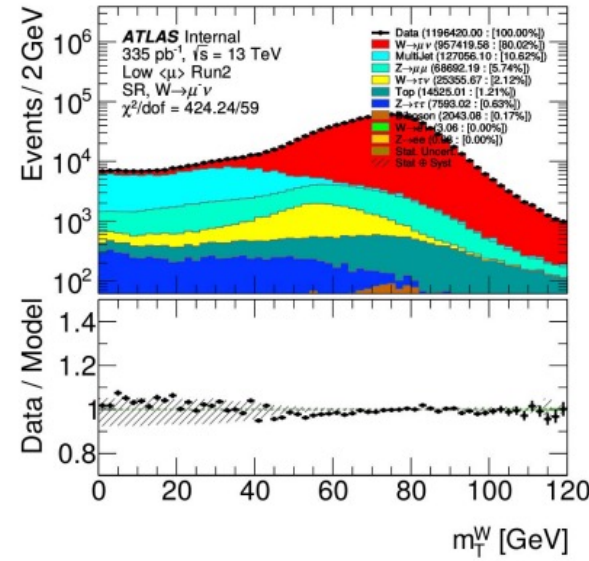
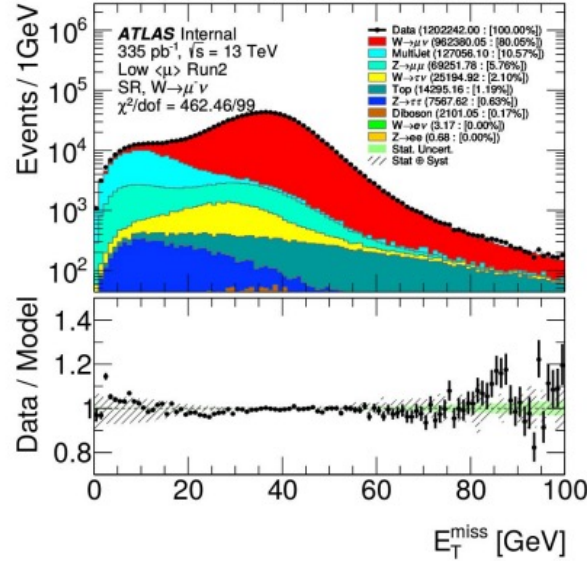
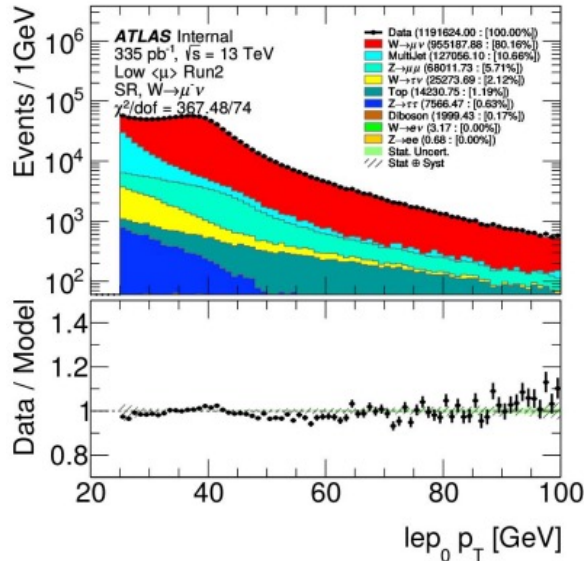
- Calculate MJ normalization by linear extrapolation fit for different anti-isolation slices
- Calculate MJ template shape by applying bin-by-bin linear shape extrapolation
- To get MJ shape as a function of $|Y|$ and $p_T^{l,\nu}$ MC samples used: $b\bar{b} + c\bar{c}$ for $W \rightarrow \mu\nu$ channels, JF17 for $W \rightarrow e\nu$ channels
- Calorimetric isolation used for $e\pm$ and track isolation used for $\mu\pm$ channels:
 - Slicing $ptvarcone20/pt$ for $W \rightarrow \mu\nu$: 8 slicing bins from 0.1 to 0.5
 - Slicing in $topoetcone20/pt$ for $W \rightarrow e\nu$: 8 slicing bins from 0.05 to 0.45



Control plots in the Signal Region for $W^- \rightarrow e^- \nu$



Control plots in the Signal Region for $W^- \rightarrow \mu^- \nu$



Fitting method

Reference coefficients created by taking moments of the polynomials using truth MC

$$\langle P_i(\theta, \varphi) \rangle = \frac{\int d\sigma(p_T, m, y, \theta, \varphi) P_i(\theta, \varphi) d\cos\theta d\varphi}{\int d\sigma(p_T, m, y, \theta, \varphi) d\cos\theta d\varphi} \longrightarrow \langle \cos\theta \rangle = \frac{1}{4} A_4$$

Truth phase space is folded into the reco phase space through polynomial templates.

$$T_{ij} = \sum_{event \in \Delta_j} \frac{P_i(\cos\theta_{CS}^{Reco}, \varphi_{CS}^{Reco}) w_{event}(r, t)}{\sigma_j \left\{ P_8(\cos\theta_{CS}^{Truth}, \varphi_{CS}^{Truth}) + \sum_{i=0}^7 A_{ij}^{ref} P_i(\cos\theta_{CS}^{Truth}, \varphi_{CS}^{Truth}) \right\}}$$

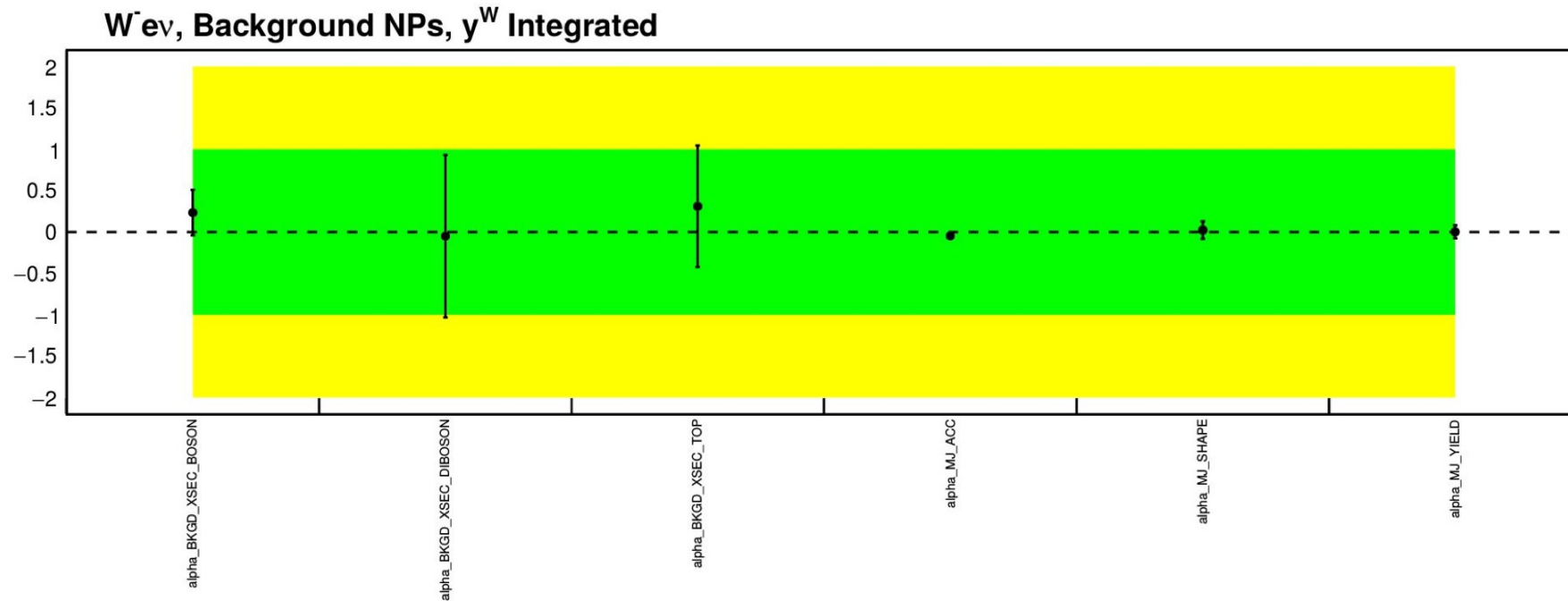
Extraction of coefficients through likelihood fit of MC templates in $p_T^{l,v}$ and $y^{l,v}$ bins.

Polynomial templates differ to ZAi analysis by extending $\cos\theta_{CS}$ range to save information for events with no real solution.

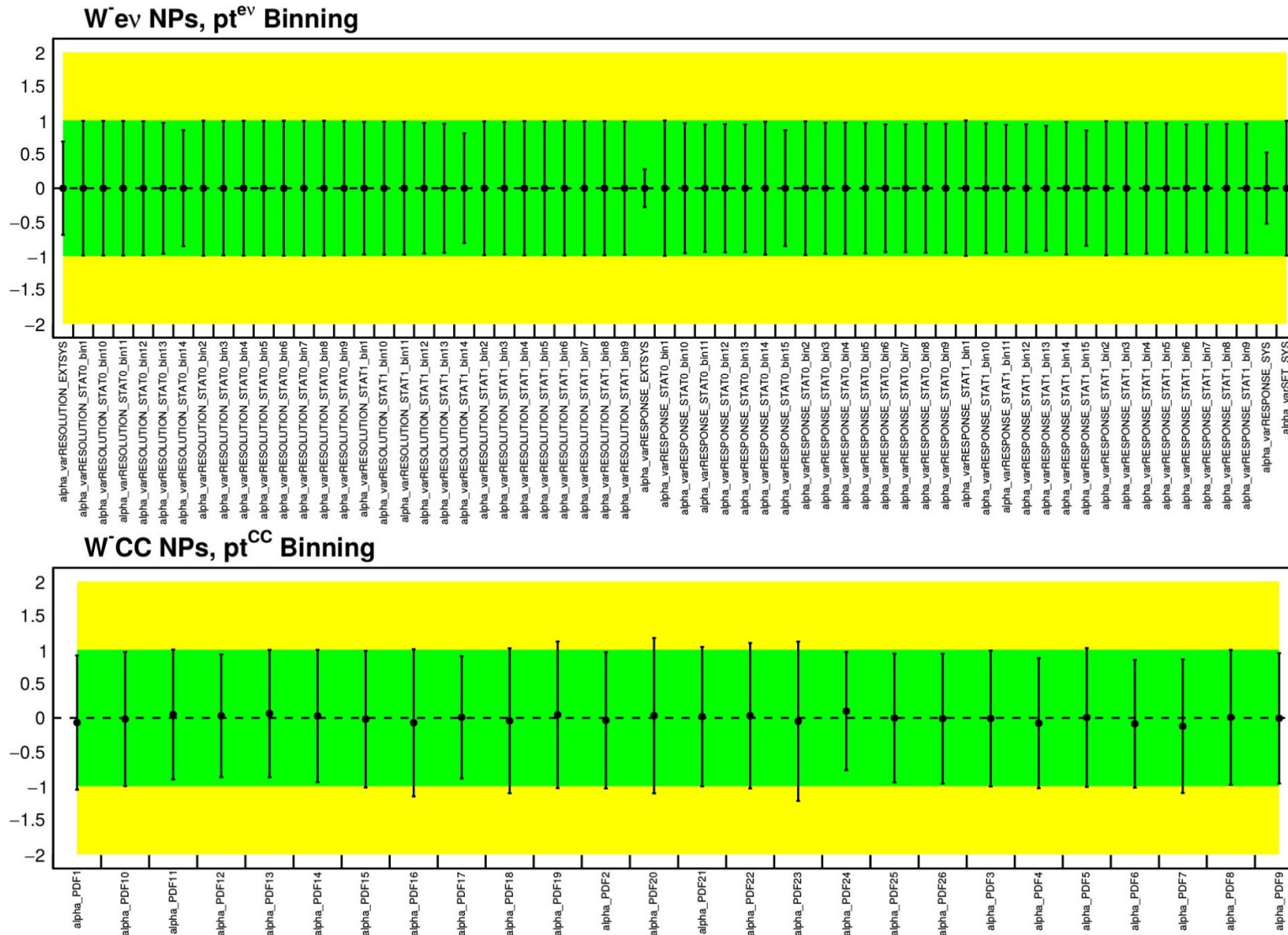
$$\mathcal{L}(A_{ij}, \mu_j | N) = \prod_{events} \left\{ \sum_{j=1}^{N_{p_T}^{bins}} \mu_j \left[T_{8,j} + \sum_{i=0}^7 A_{ij} \times T_{ij} \right] + \sum_B^{bkgds} T_B + T_{Fakes} \right\}$$

Systematics

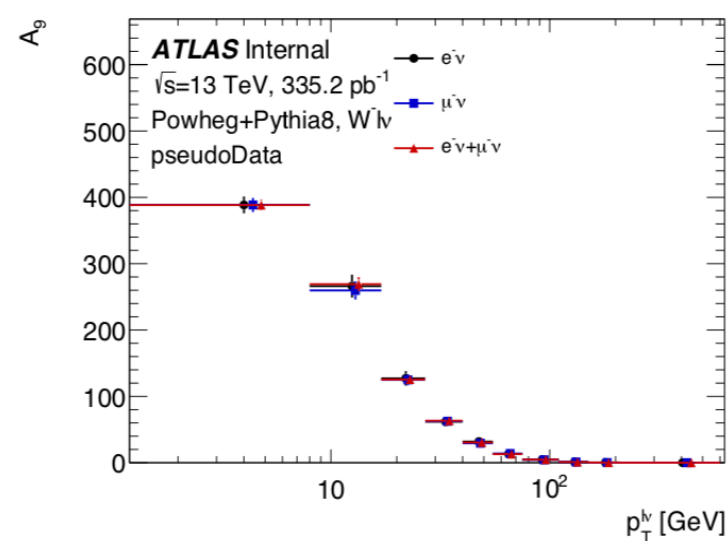
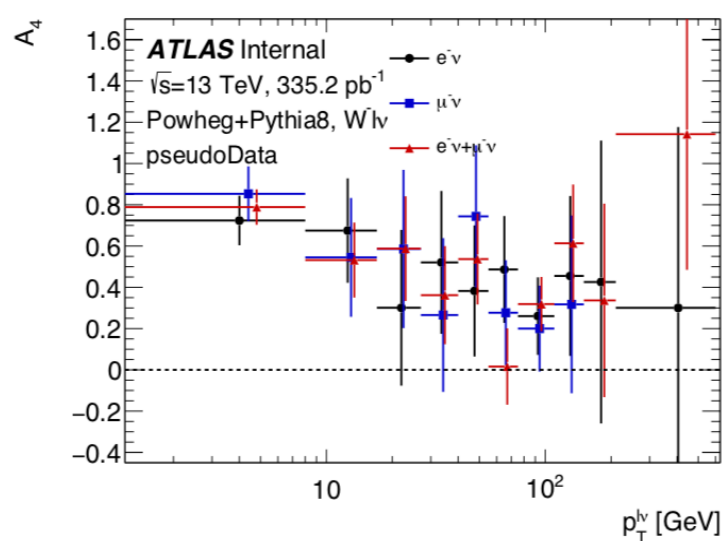
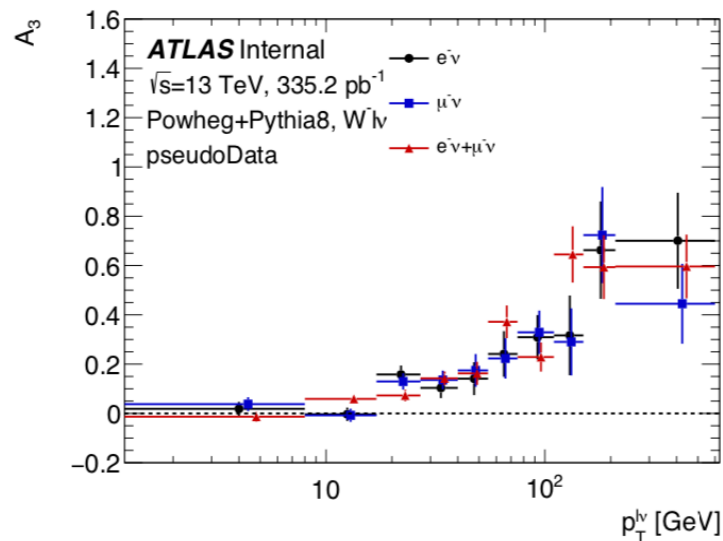
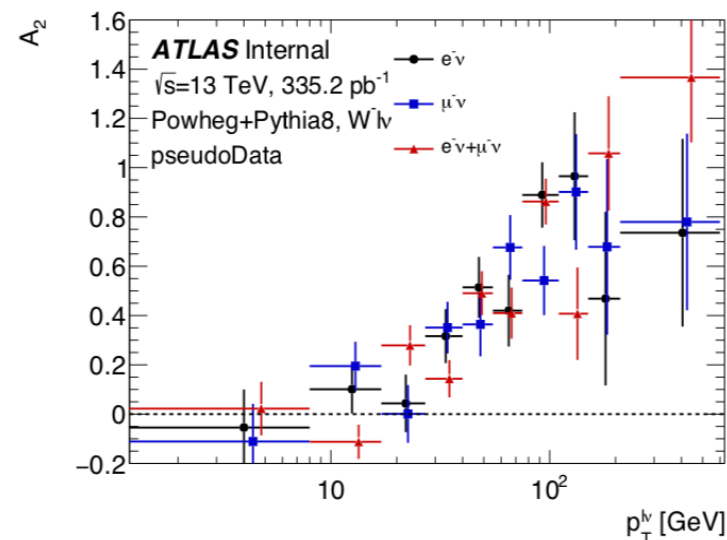
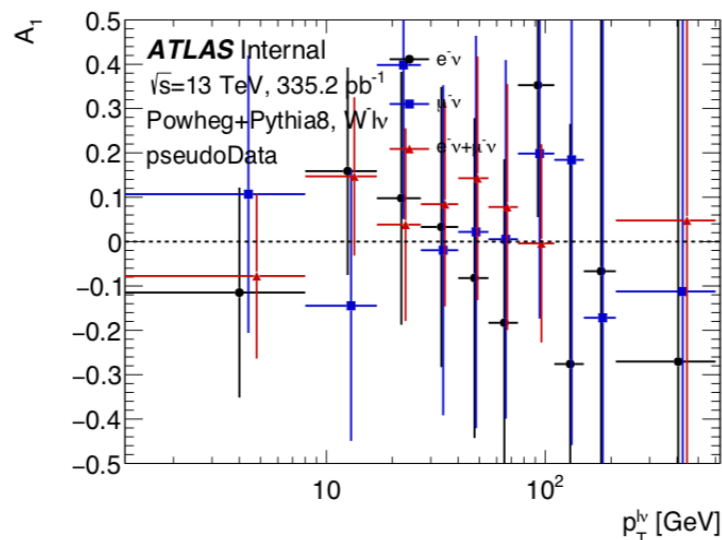
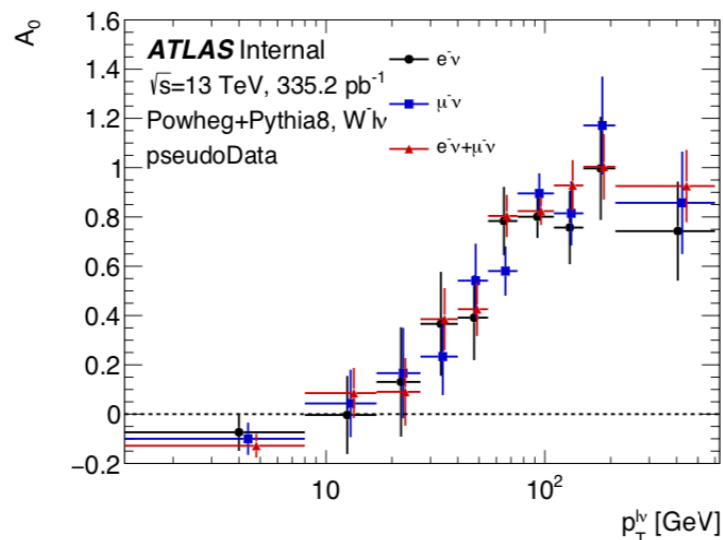
- ❖ Have largest expected systematics included (MJ, cross-section variation, hadronic recoil, and MC stat) and understood.
- ❖ Adding remaining systematics (PDF^[1] and lepton)
- ❖ Large constraint on MJ systematics from very conservative estimation^[2].



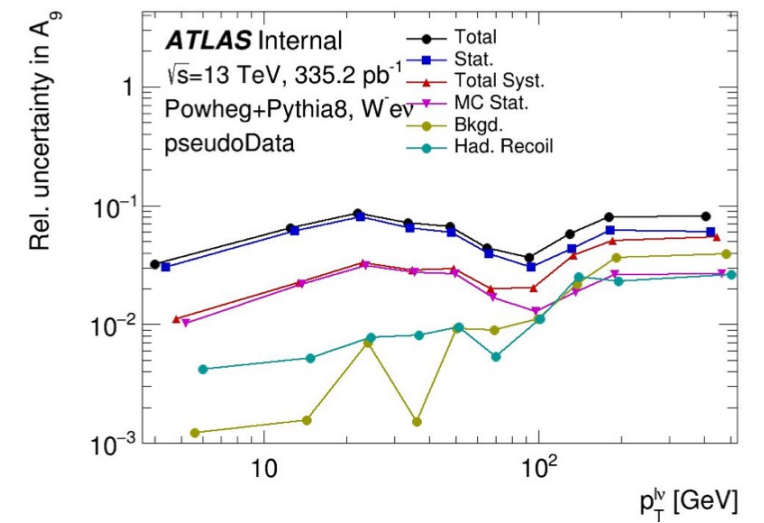
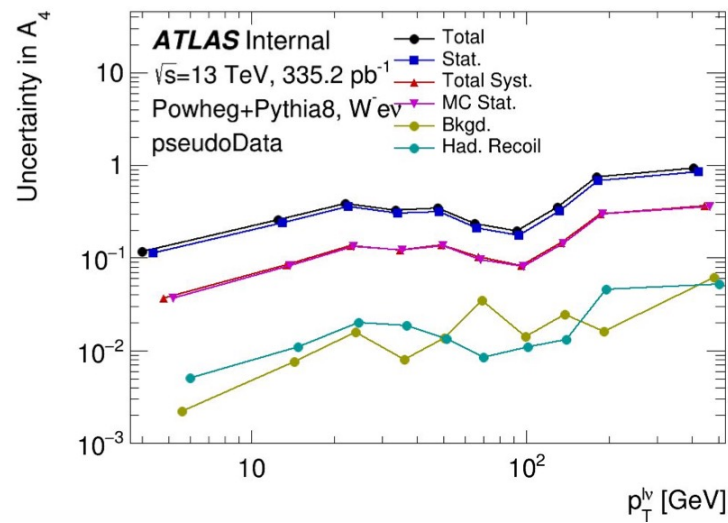
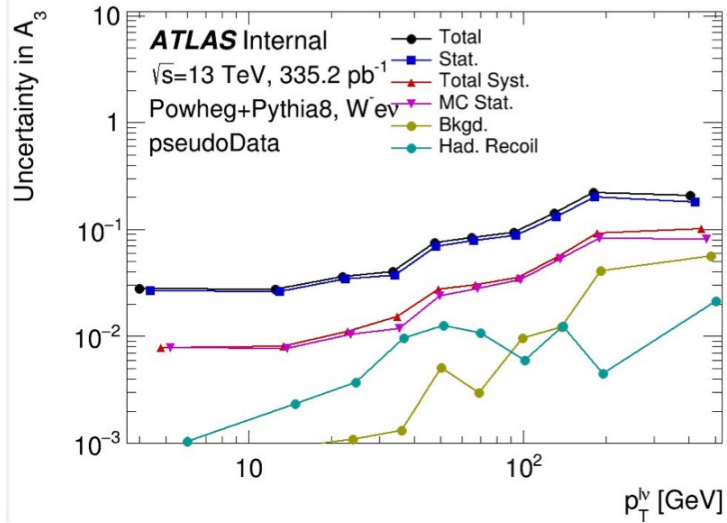
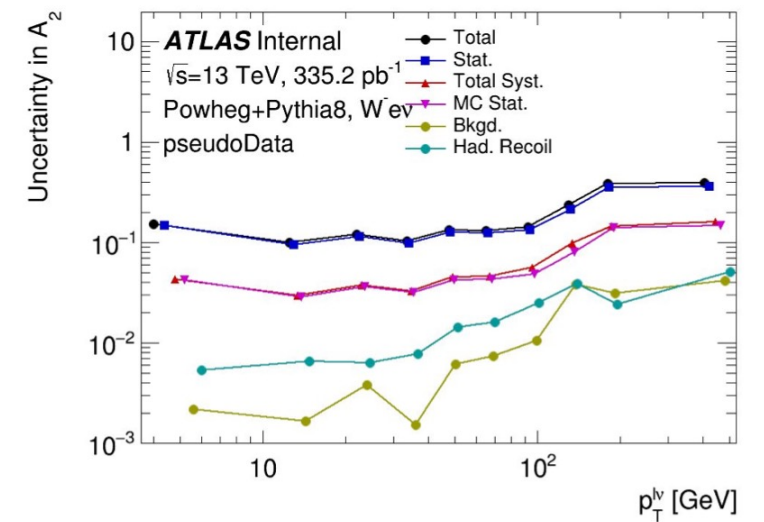
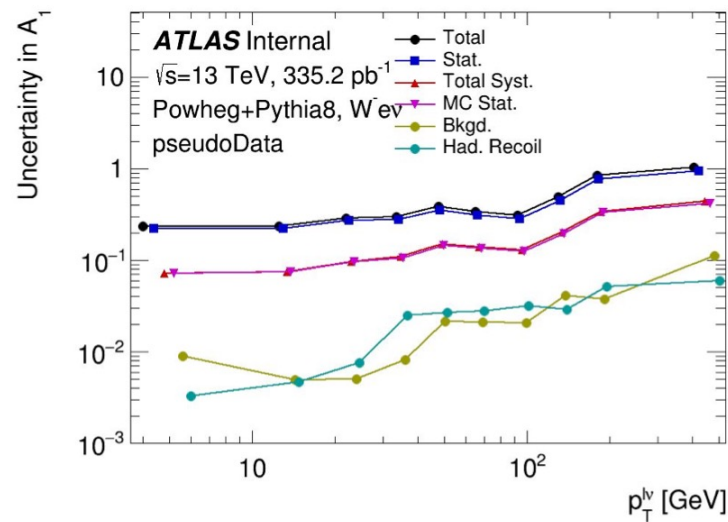
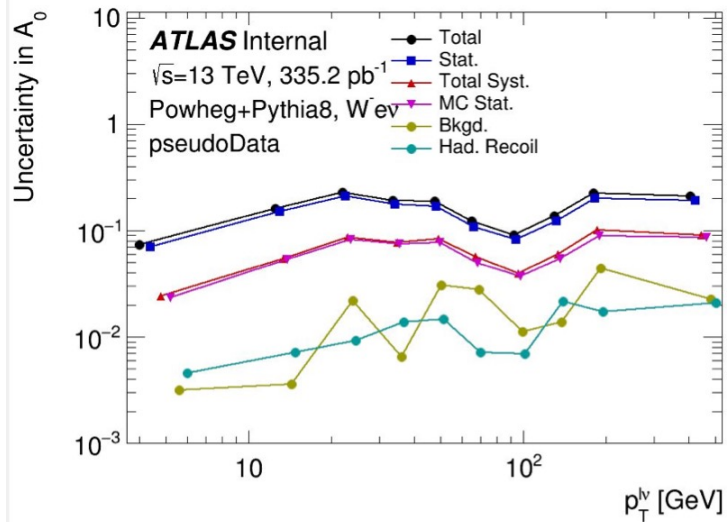
Systematics



Pseudo data ($p_T^{l,\nu}$)

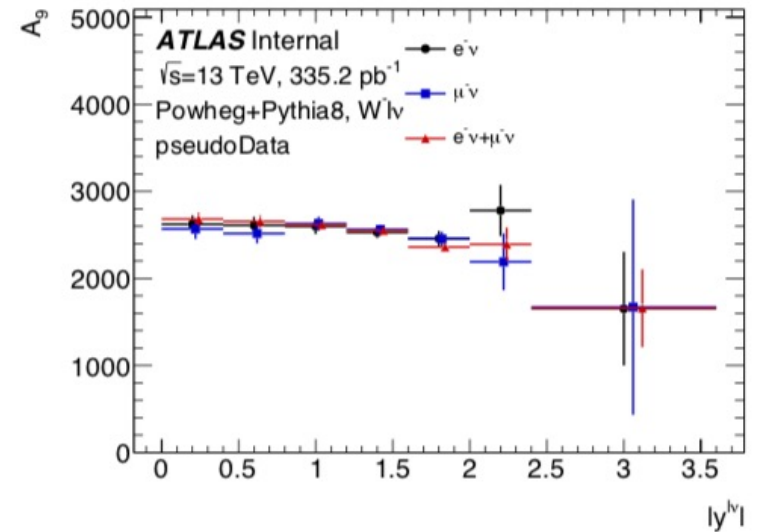
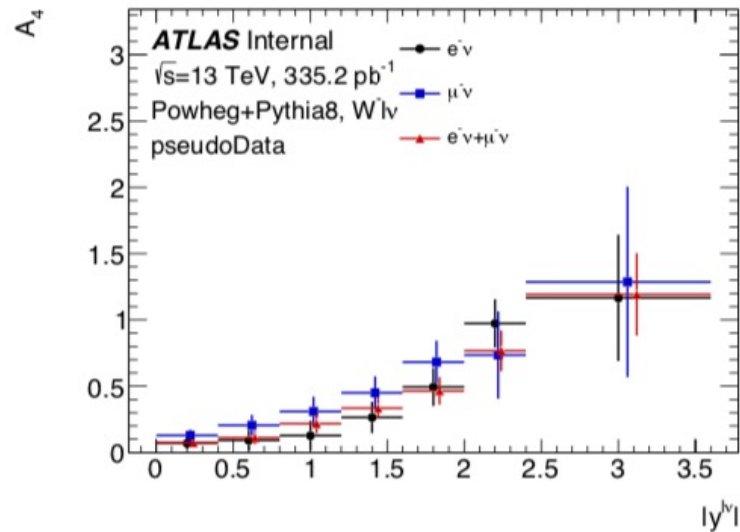
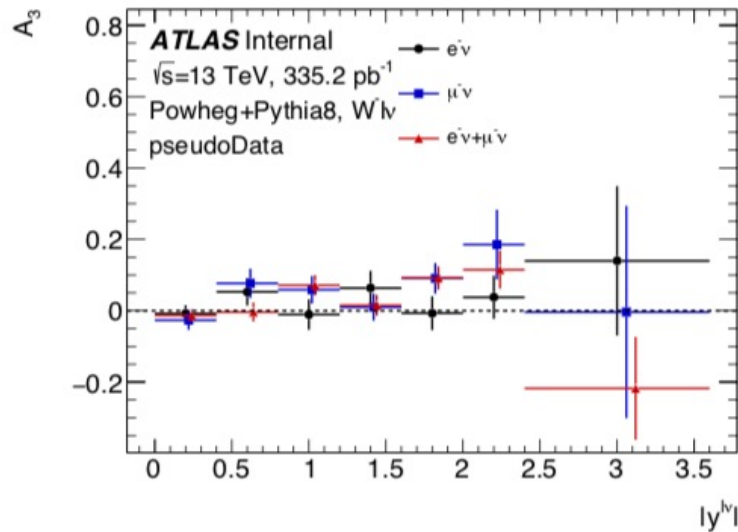
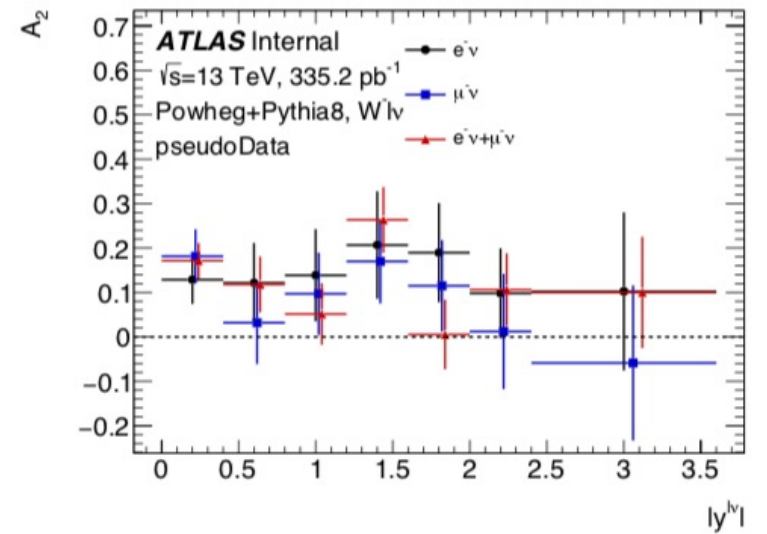
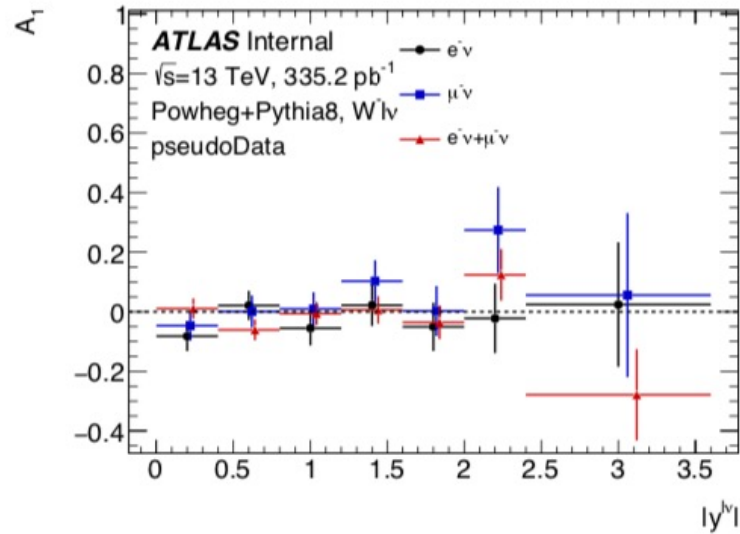
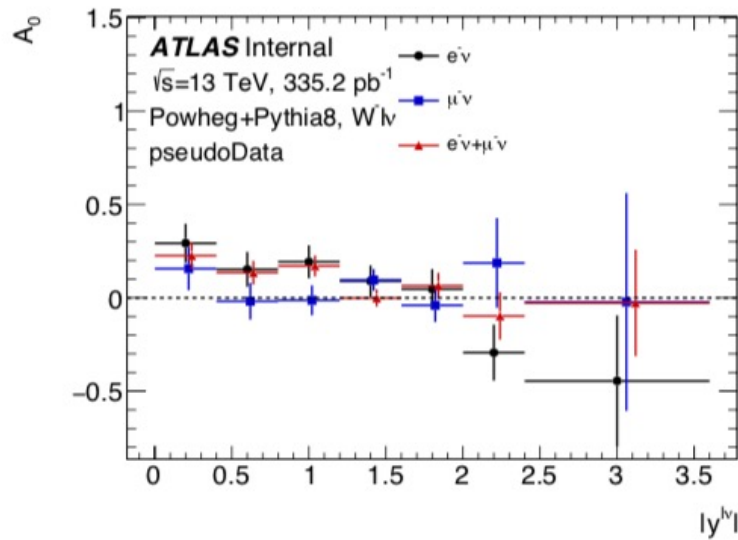


Systematic uncertainties ($p_T^{l,\nu}$)

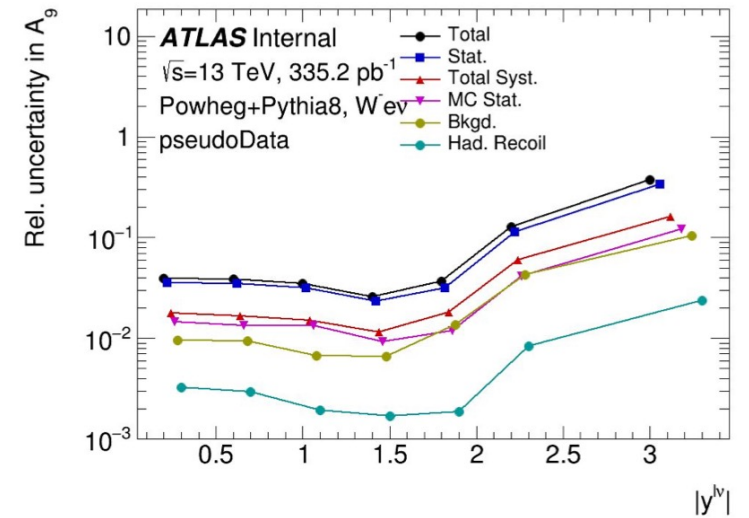
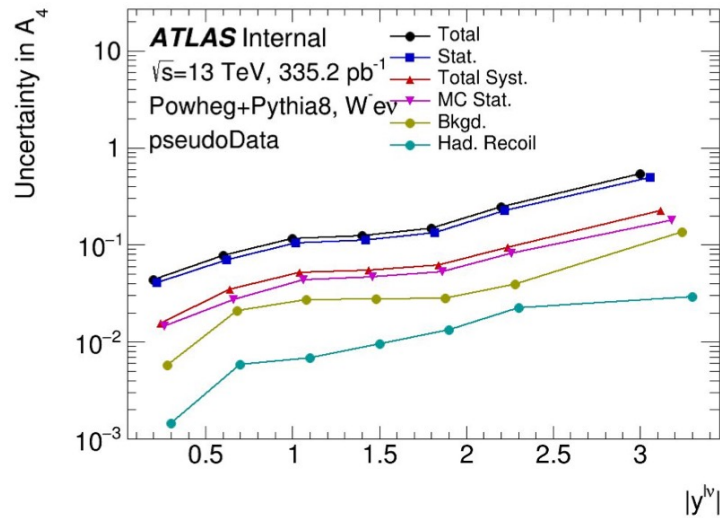
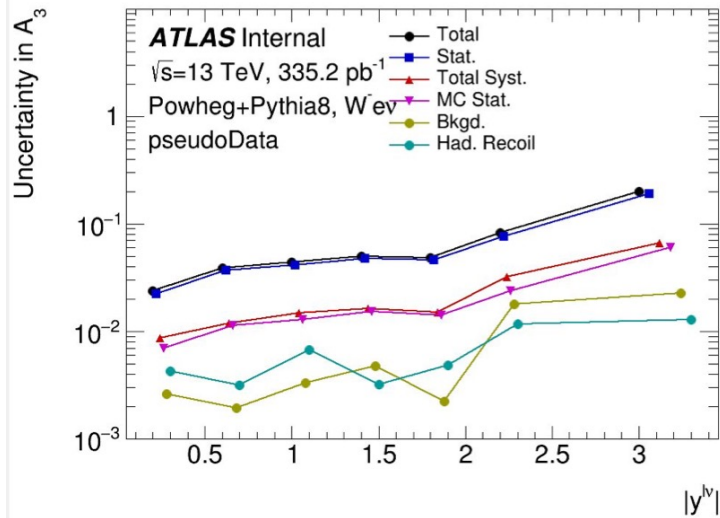
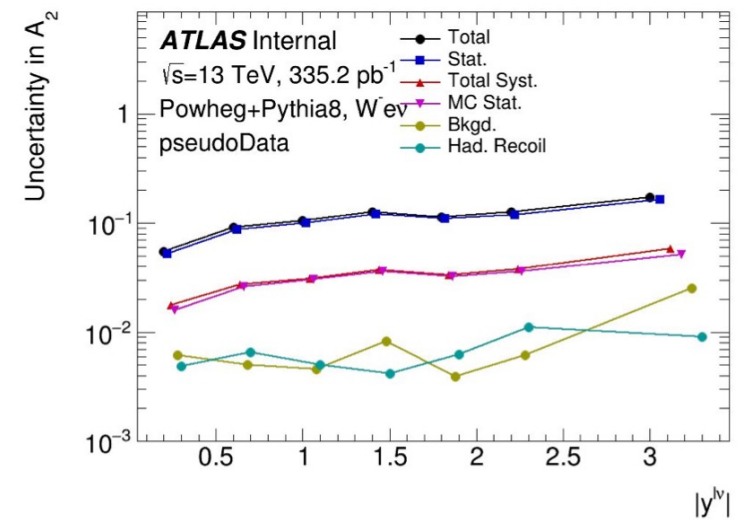
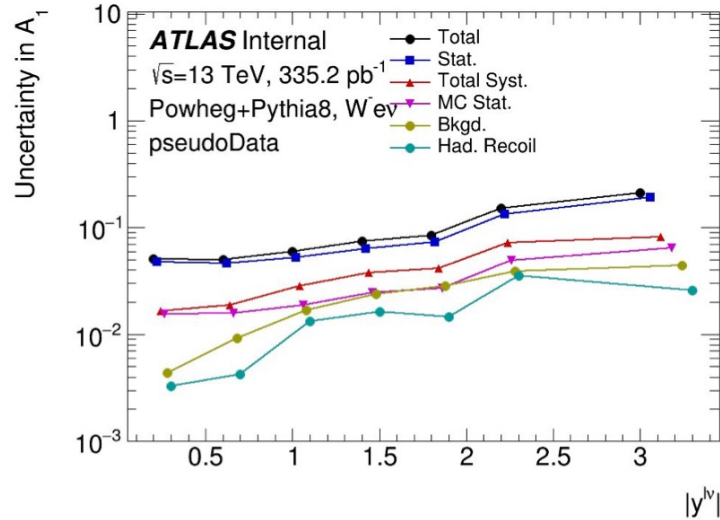
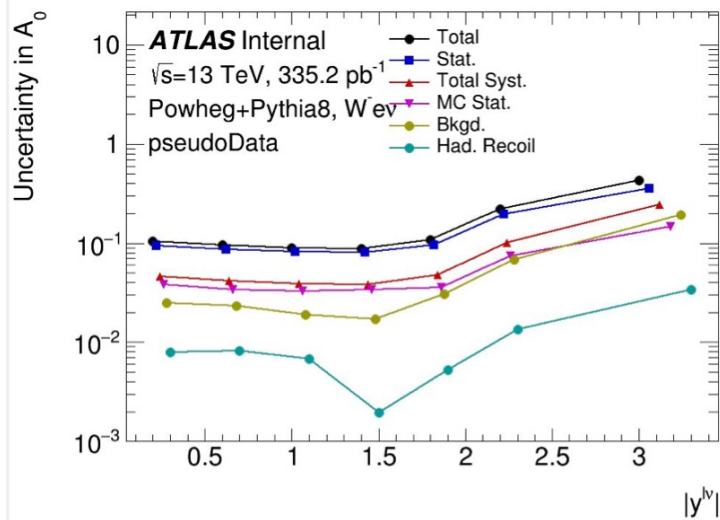


◦ Expect uncertainty to be stat dominated for all bins for both p_T^W and y^W differential measurements

Pseudo data ($y^{l,\nu}$)



Systematic uncertainties ($y^{l,\nu}$)



◦ Expect uncertainty to be stat dominated for all bins for both p^W and y^W differential measurements

Current status

- **Goal of the measurement**

- measure A_i in W events
- important input to W mass measurement and PDF fits
- main targets: A_4 (forward-backward asymm.), $A_0 - A_2$ (Lam-Tung), charge asymmetry

- **Measurement status:**

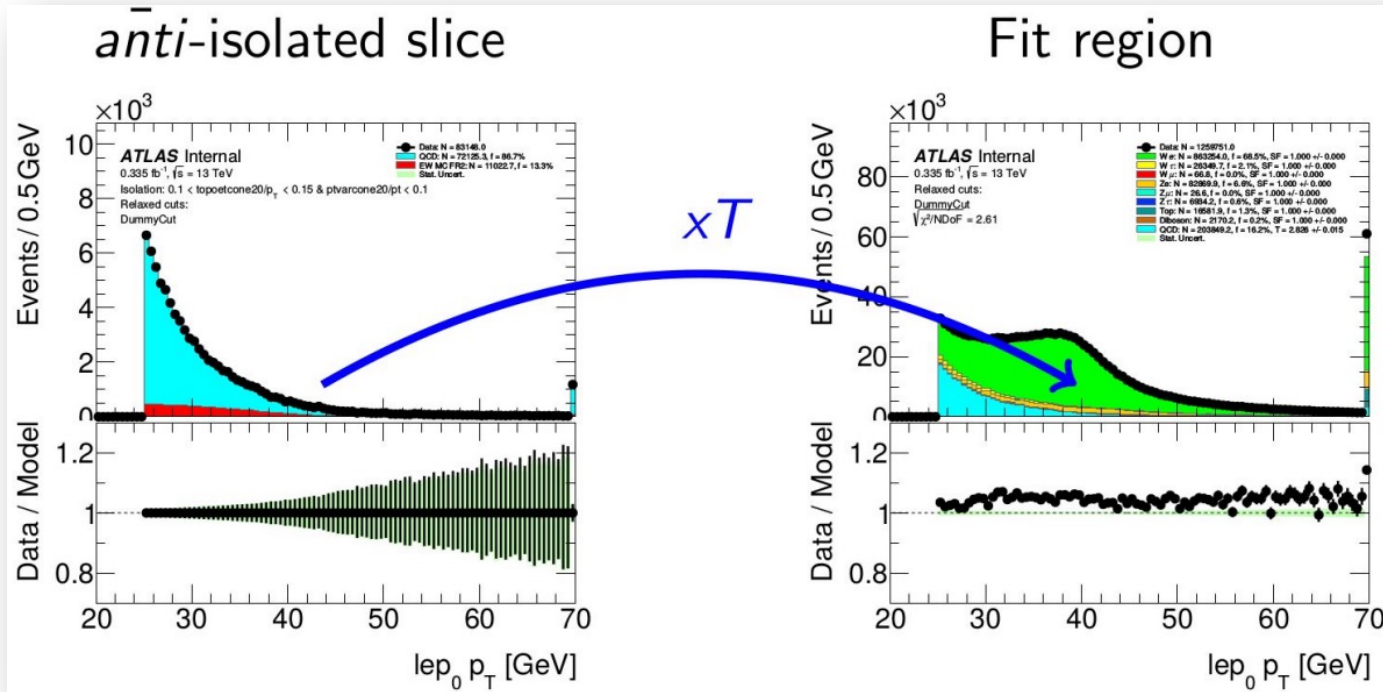
- Use W events in low- μ data (less statistics, but better hadronic recoil resolution)
- Solve for neutrino p_z via W mass constraint
- Study of sensitivity of this analysis to A_4
- MJ background estimate is in place. Other backgrounds added and understood.
- Inclusion of hadronic recoil systematics and background systematics.
- Electron and muon channel combined fit also working.

- **ToDo:**

- Refine MJ systematics correlation uncertainty model (MJ NPs heavily constrained, still being understood)
- Fit doesn't work with real data
- Add remaining systematics
- Prepare for EB request

MJ estimation updates

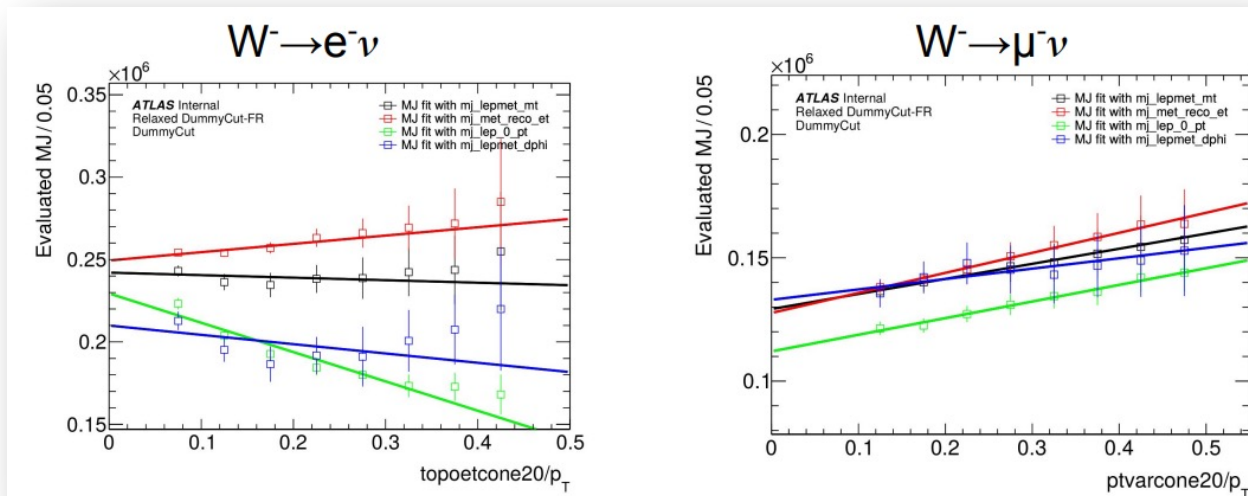
MJ background in WAi analysis (1): inclusive SR



- MJ bkgd estimate has 2 “steps”:
- Calculate MJ normalization
 - Repeat MJ estimation for different anti-isolation slices
 - Fit linear function
 - Extrapolate back to SR
- Calculate MJ template shape:
 - MJ distributions in anti-iso slices don’t match SR shape
 - Apply bin-by-bin linear shape extrapolation
 - Assign 100% uncertainty
- Use 4 discriminative variables:
 - $p_T^l, m_T, E_T^{miss}, |\Delta\varphi(l, E_T^{miss})|$
- In the fit use fixed EWT background normalization.
- To get MJ shape as a function of $|Y|$ and $p_T^{l,\nu}$ following samples used:
 - $b\bar{b} + c\bar{c}$ for $W \rightarrow \mu\nu$ channels
 - JF17 for $W \rightarrow e\nu$ channels

MJ background in WAi analysis (2): inclusive SR

MJ normalization



- The error bars are multiplied by $\sqrt{\chi^2/NDoF}$
- Take final MJ yield as mean at $ptvarcone20/pt=0.025$
- Less MJ background contribution for muon channel (as expected).
- Dominant MJ yield uncertainty comes from intersection point.

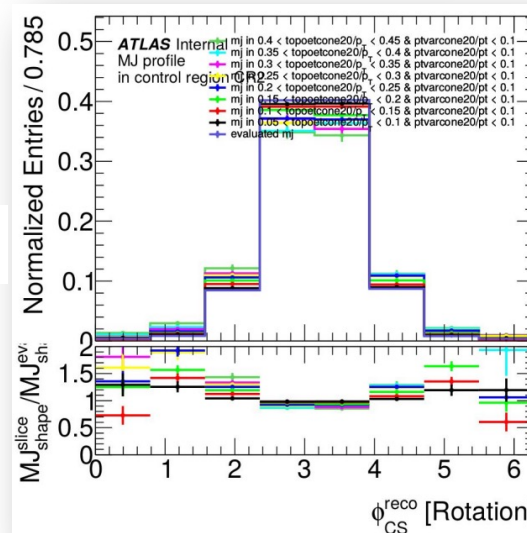
MJ template shape

In assumption, extrapolation is linear:

$$H_{MJ}^{[0.A,0.B]}[X] = H_{Data}^{[0.A,0.B]}[X] - H_{MC}^{[0.A,0.B]}[X]$$

$$\Delta H[X] = \frac{1}{4} \left\{ \frac{H_{MJ}^{0.1,0.15} - H_{MJ}^{0.3,0.35}}{4} + \frac{H_{MJ}^{0.15,0.2} - H_{MJ}^{0.35,0.4}}{4} + \frac{H_{MJ}^{0.2,0.25} - H_{MJ}^{0.4,0.45}}{4} + \frac{H_{MJ}^{0.25,0.3} - H_{MJ}^{0.45,0.5}}{4} \right\}$$

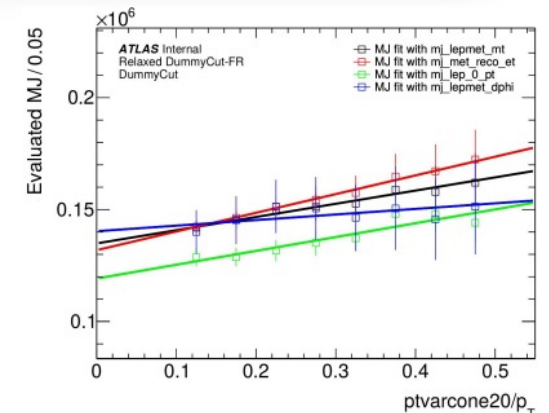
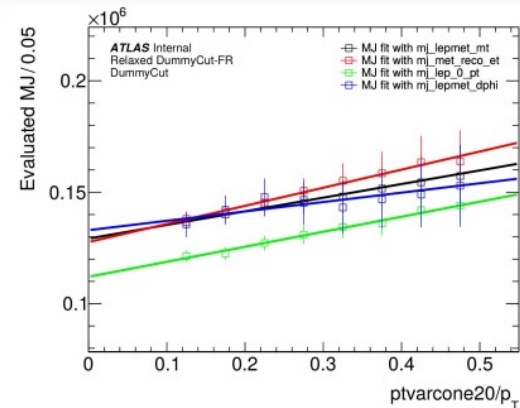
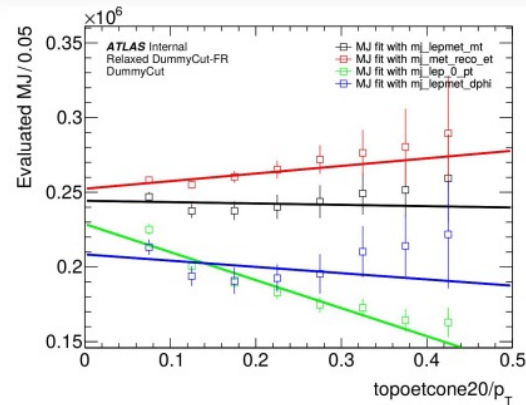
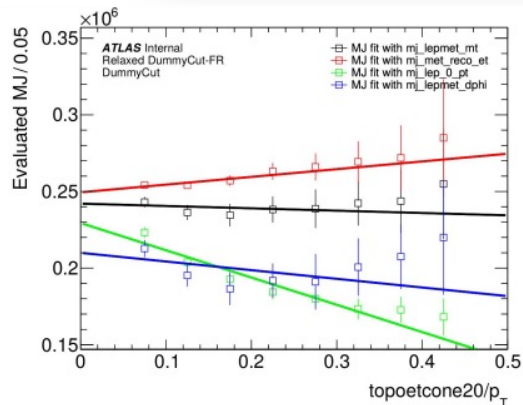
$$H_{MJ}^{sig}[X] = H_{MJ}^{0.1,0.15}[X] + 2 \cdot \Delta H[X]$$



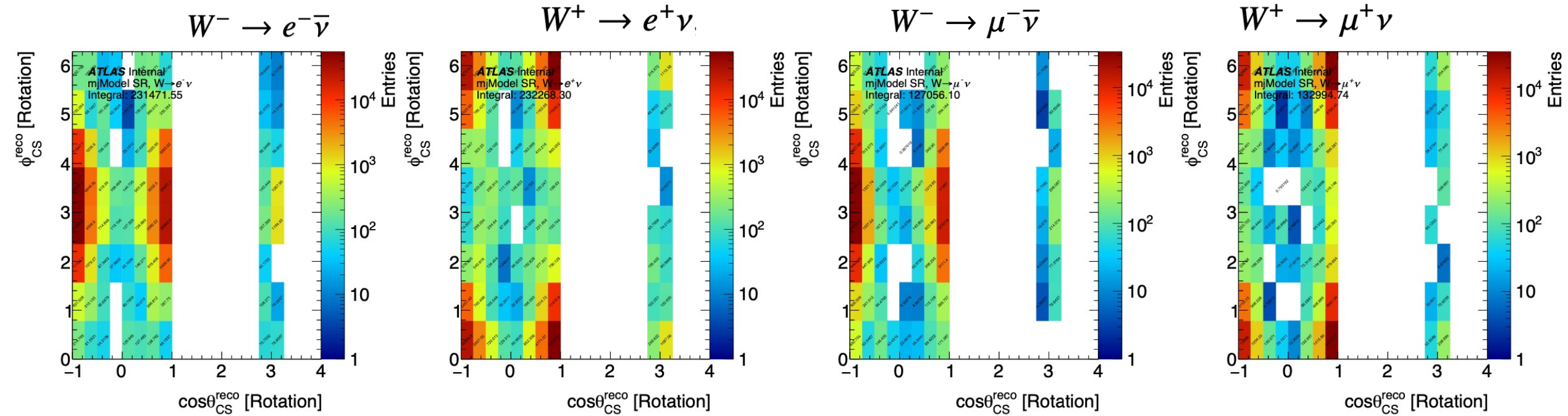
- Calculate shape correction using isolation slices for final MJ templates.
- Given the large statistical uncertainty and the linear approximation used, the shift $\Delta H[X]$ applied is assigned a 100% relative uncertainty.
- Small wrt intersection point.

MJ Systematics Summary (inclusive signal region)

Channel	Using <i>topoetcone20/p_T</i>		Using <i>ptvarcone20/p_T</i>	
	$W \rightarrow e^- \nu$	$W \rightarrow e^+ \nu$	$W \rightarrow \mu^- \nu$	$W \rightarrow \mu^+ \nu$
Total Number of MJ bkg	231472	232268	127056	132995
Luminosity and cross section	± 7199 (3.11%)	± 9098 (3.92%)	± 547 (0.43%)	± 698 (0.52%)
Intersection point	± 22928 (9.91%)	± 24910 (10.72%)	± 13313 (10.48%)	± 12276 (9.23%)
Extrapolation target	± 1246 (0.54%)	± 1172 (0.50%)	± 1575 (1.24%)	± 1433 (1.08%)
Choice of hists	± 7643 (3.30%)	± 8303 (3.57%)	± 4438 (3.49%)	± 4092 (3.08%)
Isolation correction	N/A	N/A	N/A	N/A
Correlated Uncertainty	± 24032 (10.38%)	± 26519 (11.42%)	± 13324 (10.49%)	± 12296 (9.25%)
Data Stat.	± 490 (0.21%)	± 463 (0.20%)	± 795 (0.63%)	± 809 (0.61%)
MC Stat.	± 262 (0.11%)	± 229 (0.10%)	± 990 (0.78%)	± 880 (0.66%)
Shape Correction	± 2961 (1.28%)	± 2769 (1.19%)	± 927 (0.73%)	± 915 (0.69%)
Uncorrelated Uncertainty	± 3012 (1.30%)	± 2816 (1.21%)	± 1572 (1.24%)	± 1505 (1.13%)



MJ templates for $\cos\theta_{CS}$ vs. ϕ_{CS} : inclusive SR



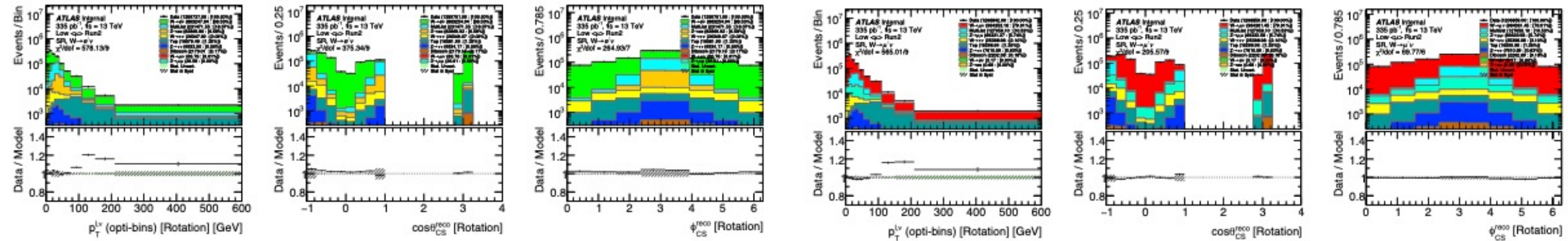
- The 2D templates are derived in the same way as described for 1D histograms

MJ shape as function of $|Y|$ and $p_T^{\ell,\nu}$: task definition

1D control plots in SR (only stat unc. + MJ systematics are shown.):

$W^- \rightarrow e^- \nu$

$W^- \rightarrow \mu^- \nu$



W-Ai analysis uses more complicated approach:

- Building 2D histograms $\cos \theta_{CS}$ vs. ϕ_{CS} as function of η or $p_T^{\ell,\nu}$

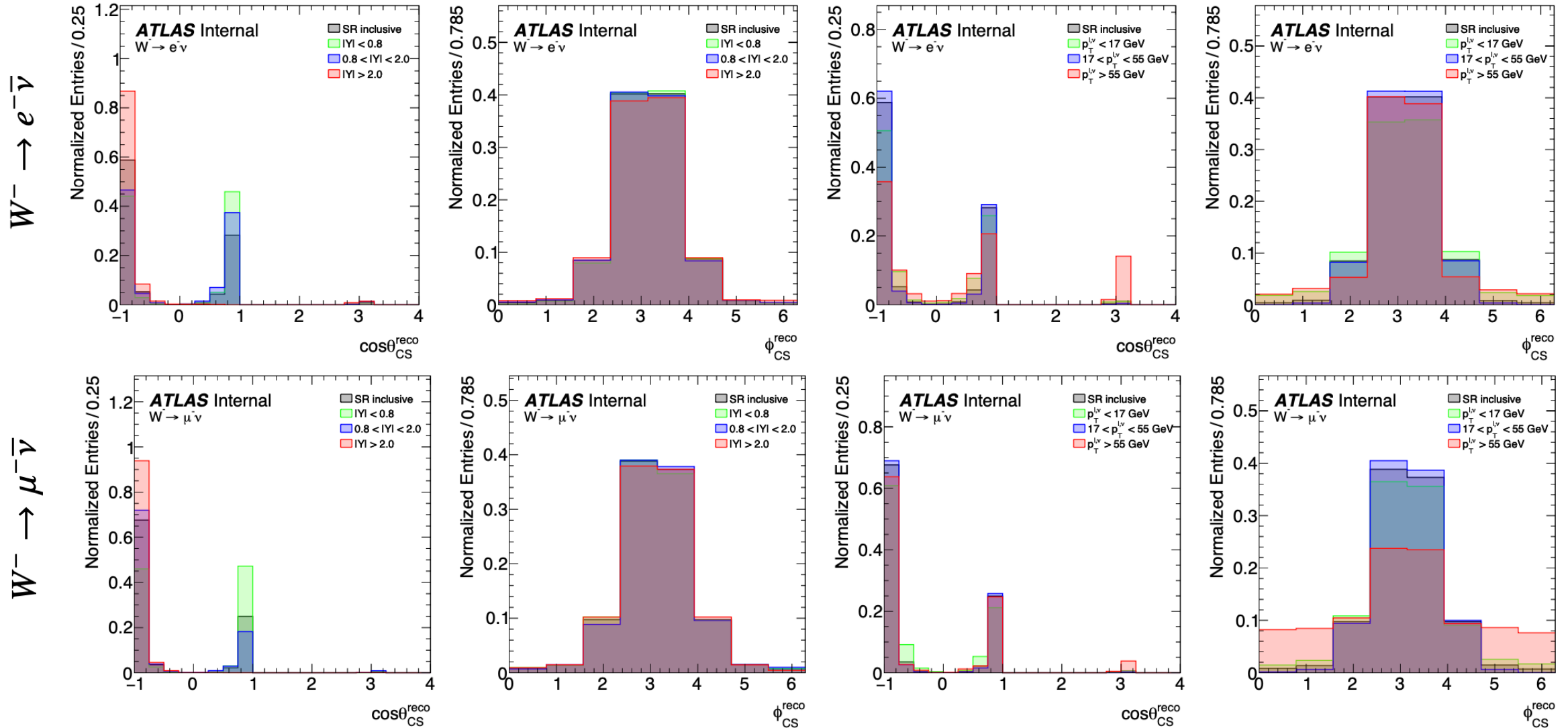
$$p_T^{\ell,\nu} : [0, 8, 17, 27, 40, 55, 75, 110, 150, 210, 600];$$

$$|Y| : [0, 0.4, 0.8, 1.2, 1.6, 2.0, 2.4, 3.6]$$

- Above 100 GeV the data/MC agreement in the $p_T^{\ell,\nu}$ degrades due to missing truth p_T^W reweighting correction above this range.

Closure test by DD method: coarse bins

$p_T^{\ell, \nu} : [0, 17, 55, 600];$
 $|Y| : [0, 0.8, 2.0, 3.6]$

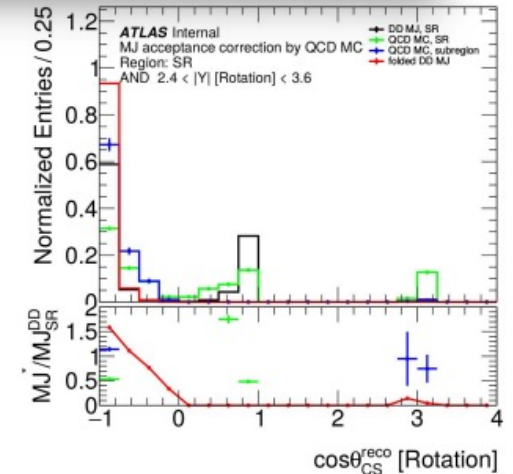
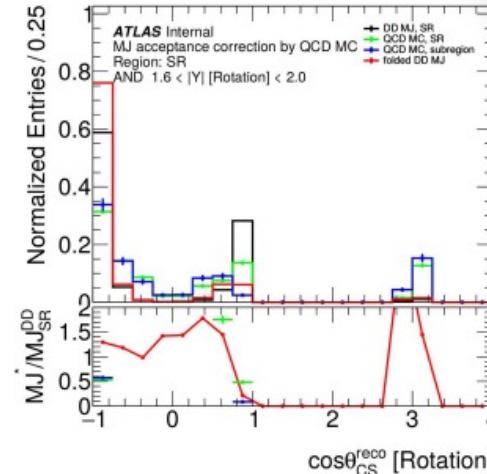
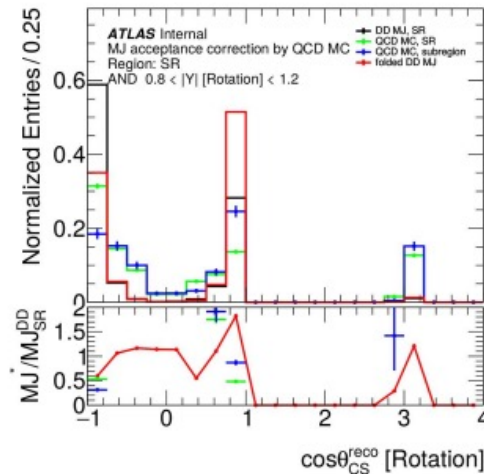
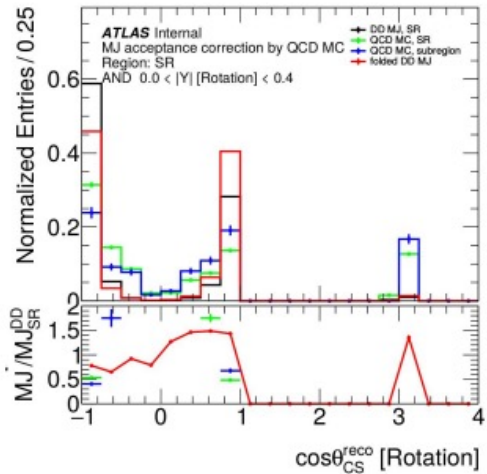
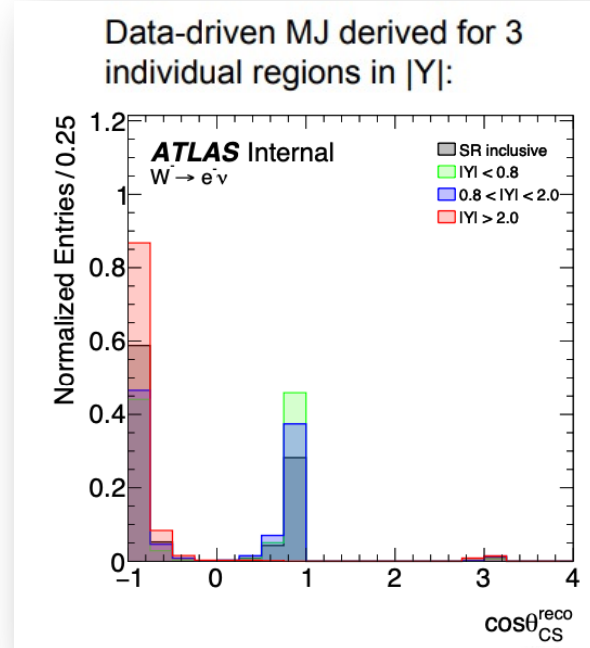


MJ shape as function of $|Y|$ and $p_T^{l,\nu}$: studies

$p_T^{l,\nu} : [0, 17, 55, 600];$
 $|Y| : [0, 0.8, 2.0, 3.6]$

- There are strong dependence of the $\cos\theta_{CS}$ MJ template shape as function of $|Y^{l,\nu}|$ and φ_{CS} as a function of $p_T^{l,\nu}$
- Acceptance correction functions were built to correct MJ data-driven template derived in the SR to the given $|Y^{l,\nu}|$ or $p_T^{l,\nu}$ slice using the samples:
 - $b\bar{b} + c\bar{c}$ for $W \rightarrow \mu\nu$ channels
 - JF17 for $W \rightarrow e\nu$ channels
- Acceptance correction functions for number of $|Y^{l,\nu}|$ bins, fully uncorrelated uncertainty on these acceptance correction functions is applied on the MJ estimate:

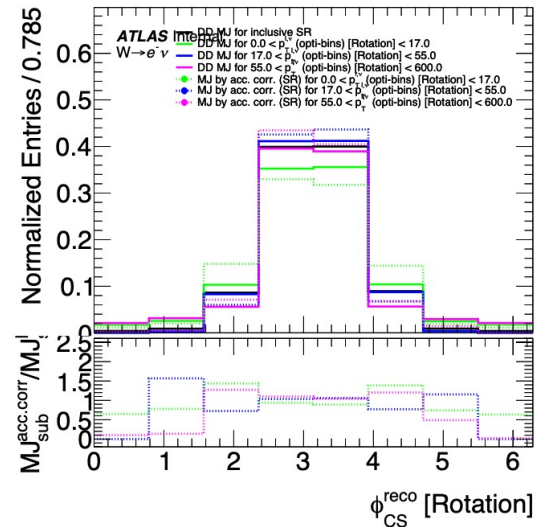
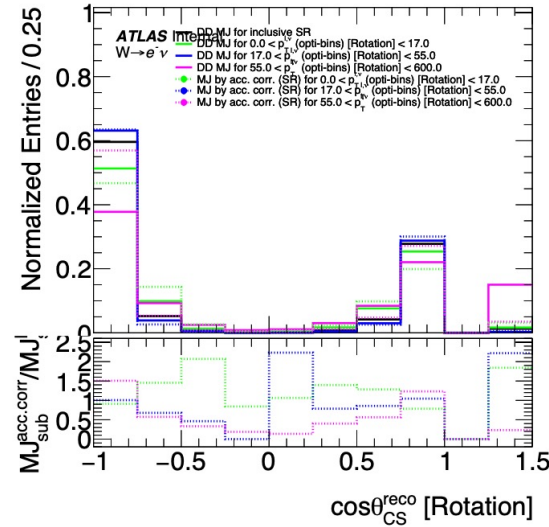
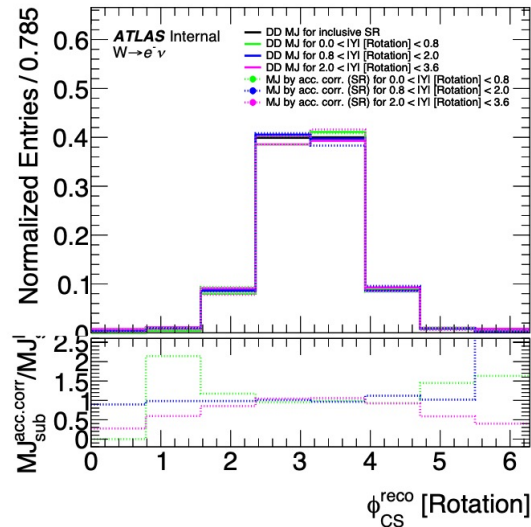
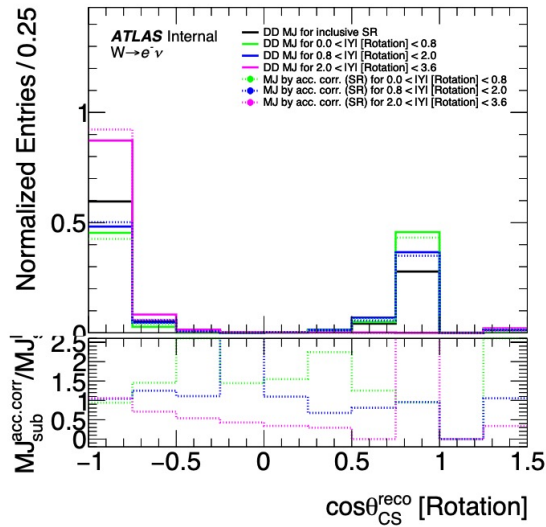
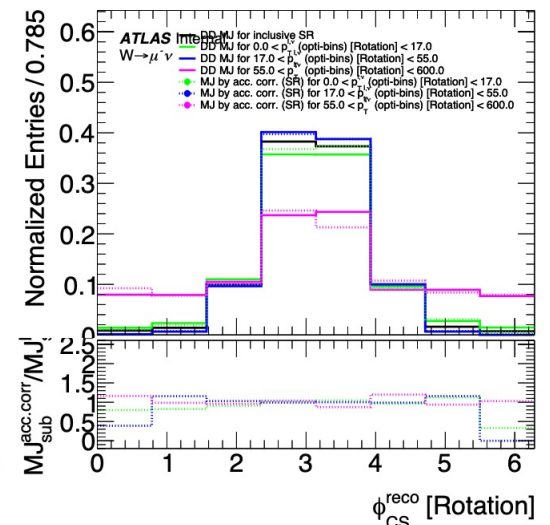
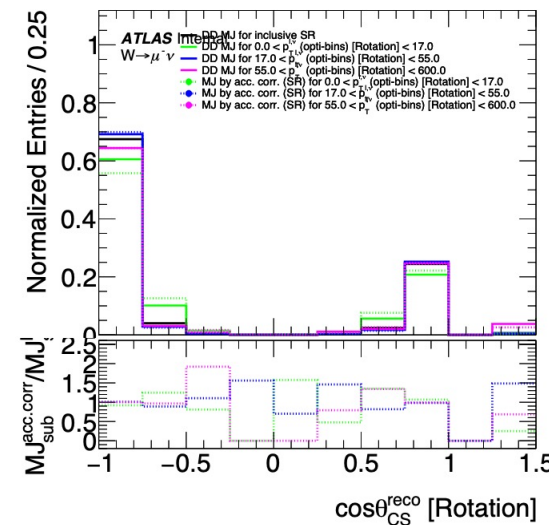
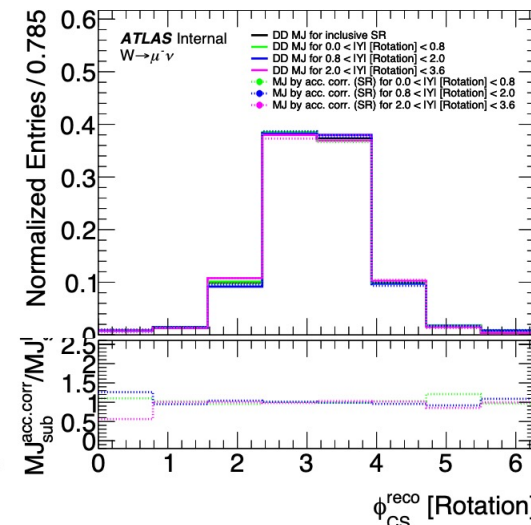
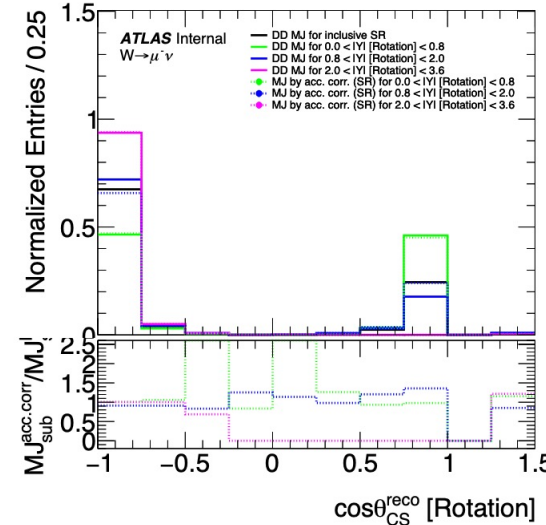
$$A_j = h_{MC}^{bin_j} / h_{MC}^{SR} \quad h_{MJ}^{bin_j} = h_{MJ,data}^{SR} \times A_j$$



Acc. corr. function: closure test

$$p_T^{\ell, \nu} : [0, 17, 55, 600];$$

$$|Y| : [0, 0.8, 2.0, 3.6]$$

 $W^- \rightarrow e^- \bar{\nu}$

 $W^- \rightarrow \mu^- \bar{\nu}$


MJ shape as function of $|Y|$ and $p_T^{\ell,\nu}$

- There are strong dependence of the $\cos\theta_{CS}$ MJ template shape as function $|Y|$ and ϕ_{CS} as function of $p_T^{\ell,\nu}$
- Acceptance correction functions were build using QCD MC samples to correct MJ data-driven template derived in the SR to the given $|Y|$ or $p_T^{\ell,\nu}$ slice:

$$h_{MJ}^{bin_i} = h_{MJ}^{SR} \cdot h_{QCDMC}^{bin_i} / h_{QCDMC}^{SR}$$

Problem:

- Comparing to MJ template given by AccCocc in the last $p_T^{\ell,\nu}$ bin for e^- with the DD MJ, last one predicts higher MJ contribution to the no solution bin + overall shape differs a lot.
- **DD MJ template derived from SR might be "too far" from DD MJ template derived for $p_T^{\ell,\nu} > 55$ GeV bin.**
- We might not see this effect in muons due to the fact this channel is "cleaner".
- Integrated distributions for Y coarse bins looks good due to small MJ fraction for the last coarse $p_T^{\ell,\nu}$ bin.

Solution:

- Instead of using SR use the coarse bins: $h_{QCDMC}^{CoarseBin_k}$ and $h_{MJ}^{CoarseBin_k}$.

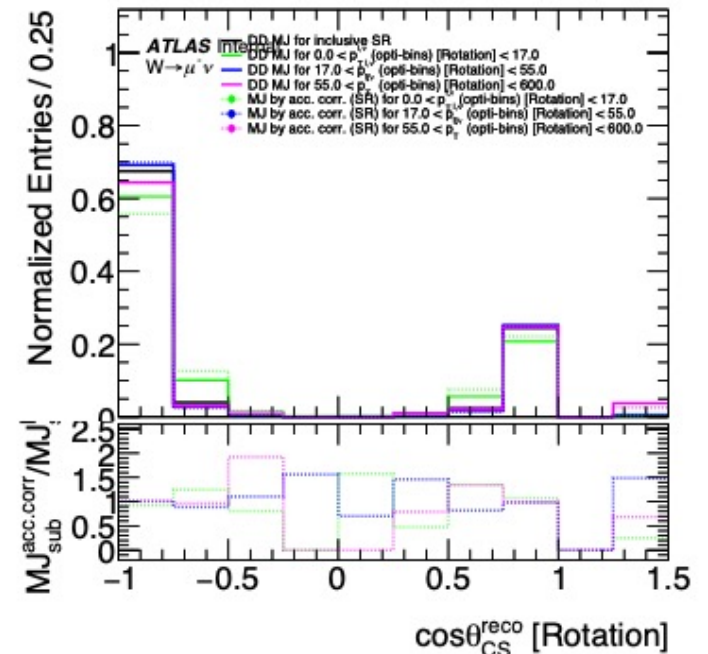
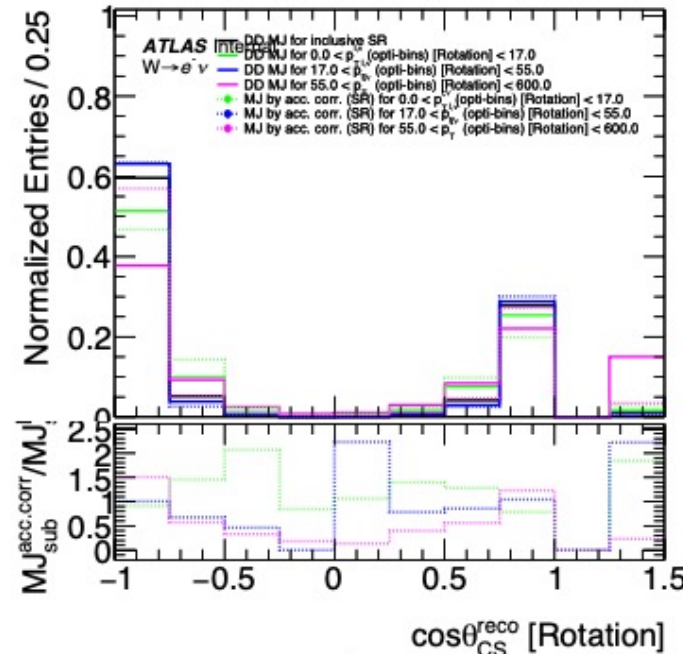
$$p_T^{\ell,\nu} : [0, 17, 55, 600];$$

$$|Y| : [0, 0.8, 2.0, 3.6]$$

Closure test for Data-driven MJ and Acc.Corr MJ derived for 3 individual regions in $p_T^{\ell,\nu}$:

$$W^- \rightarrow e^- \nu$$

$$W^- \rightarrow \mu^- \nu$$



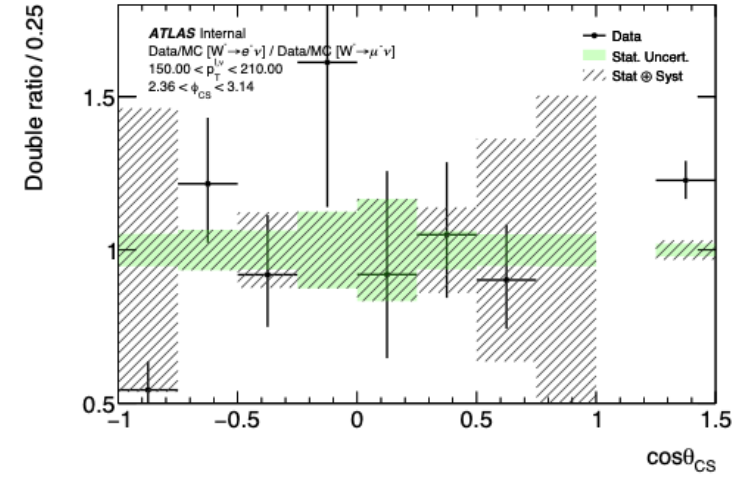
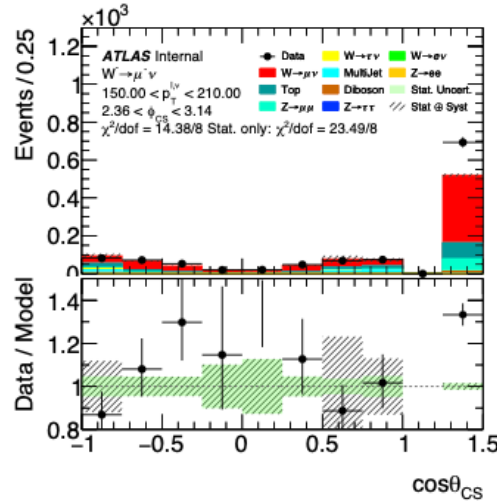
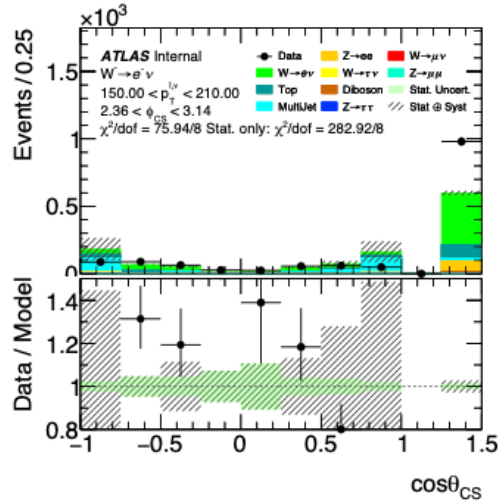
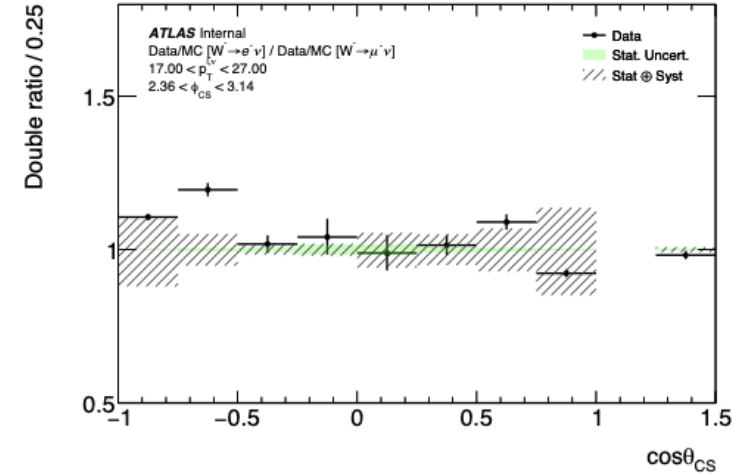
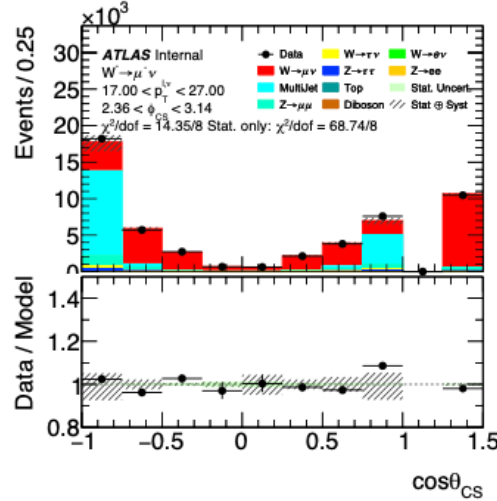
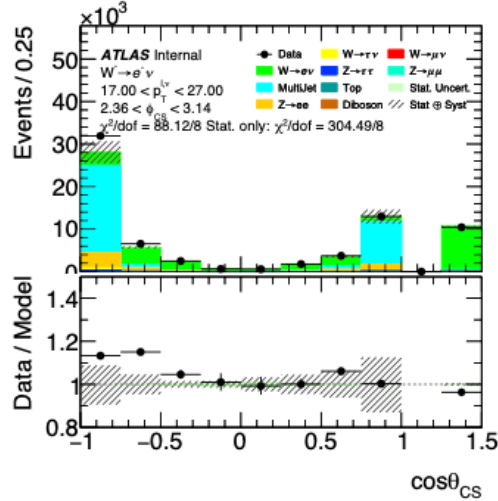
Some W-Ai control plots and e/μ compatibility

Compare 2 $p_T^{\ell,\nu}$ bins: $17 < p_T^{\ell,\nu} < 27$ GeV and $150 < p_T^{\ell,\nu} < 210$ GeV

Before

$17 < p_T^{\ell,\nu} < 27$ GeV

$150 < p_T^{\ell,\nu} < 210$ GeV

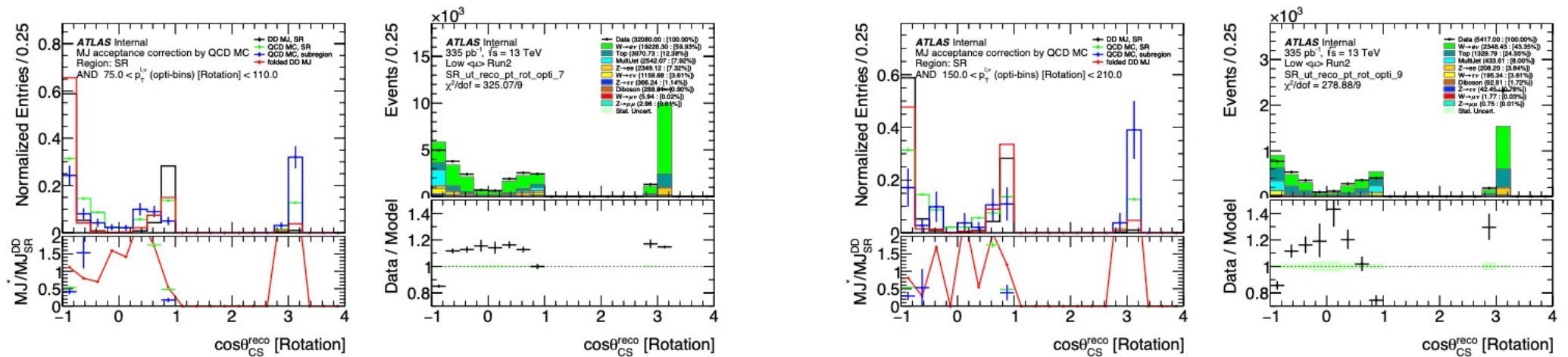


Note: discrepancy in the no solution bins in $p_{TW} > 100$ GeV is from absence of the PTRW

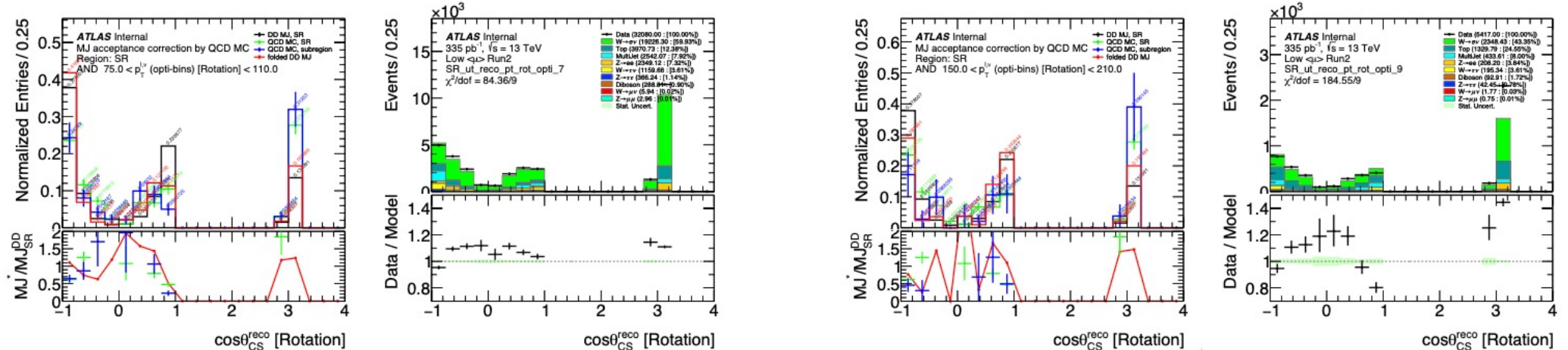
Test with 1D $\cos\theta_{CS}$ integrated over ϕ_{CS} in some $p_T^{\ell,\nu}$ bins

Instead of using SR, use the coarse bins: $h_{QCDMC}^{CoarseBin_k}$ and $h_{MJ}^{CoarseBin_k}$.

Before



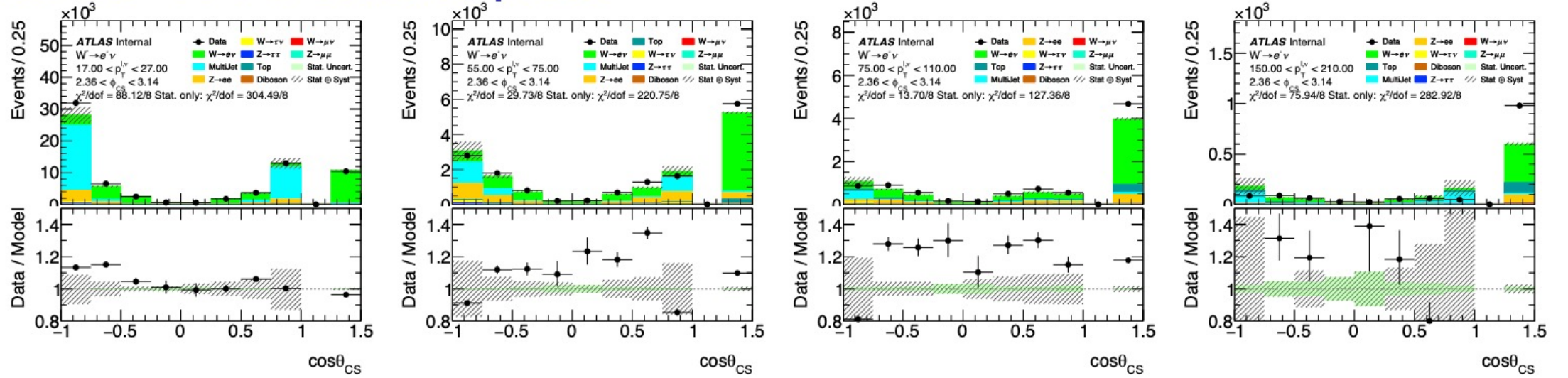
After



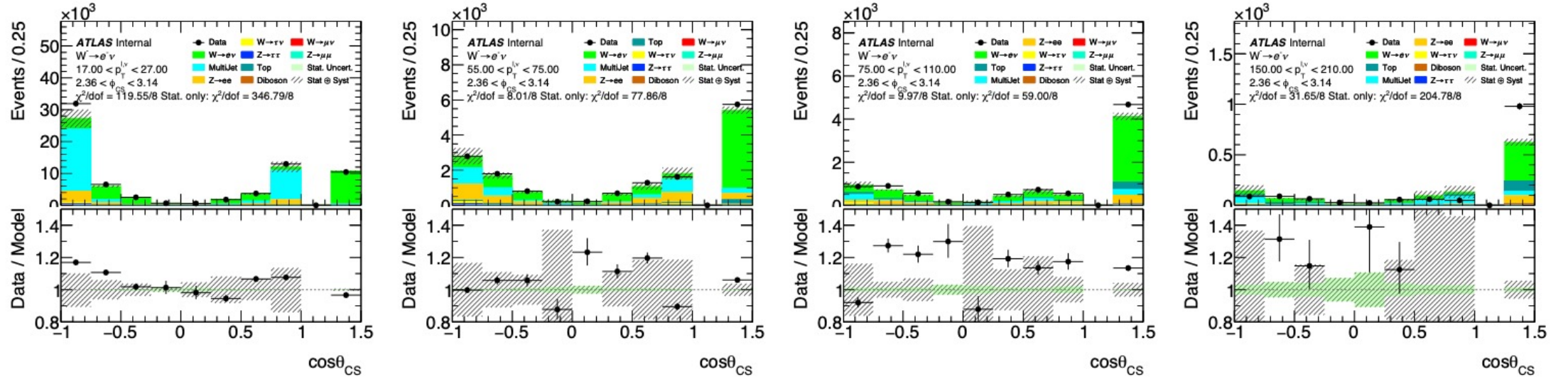
Note: discrepancy in the no solution bins in $p_{TW} > 100 \text{ GeV}$ is from absence of the PTRW

Test with 3D W-Ai control plots

Before



After

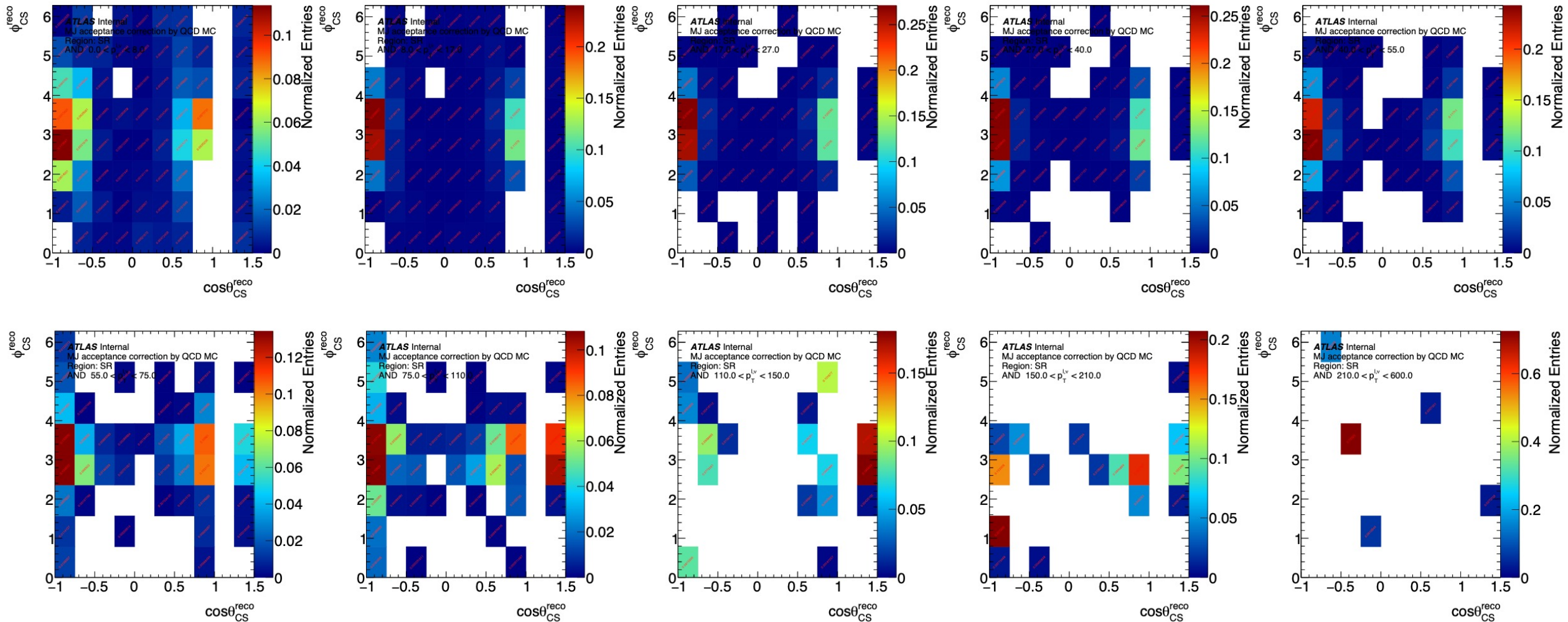


Only some $p_T^{\ell,\nu}$ coarse bins and ϕ_{CS} slices are shown.

Note: discrepancy in the no solution bins in $p_{TW} > 100$ GeV is from absence of the PTRW

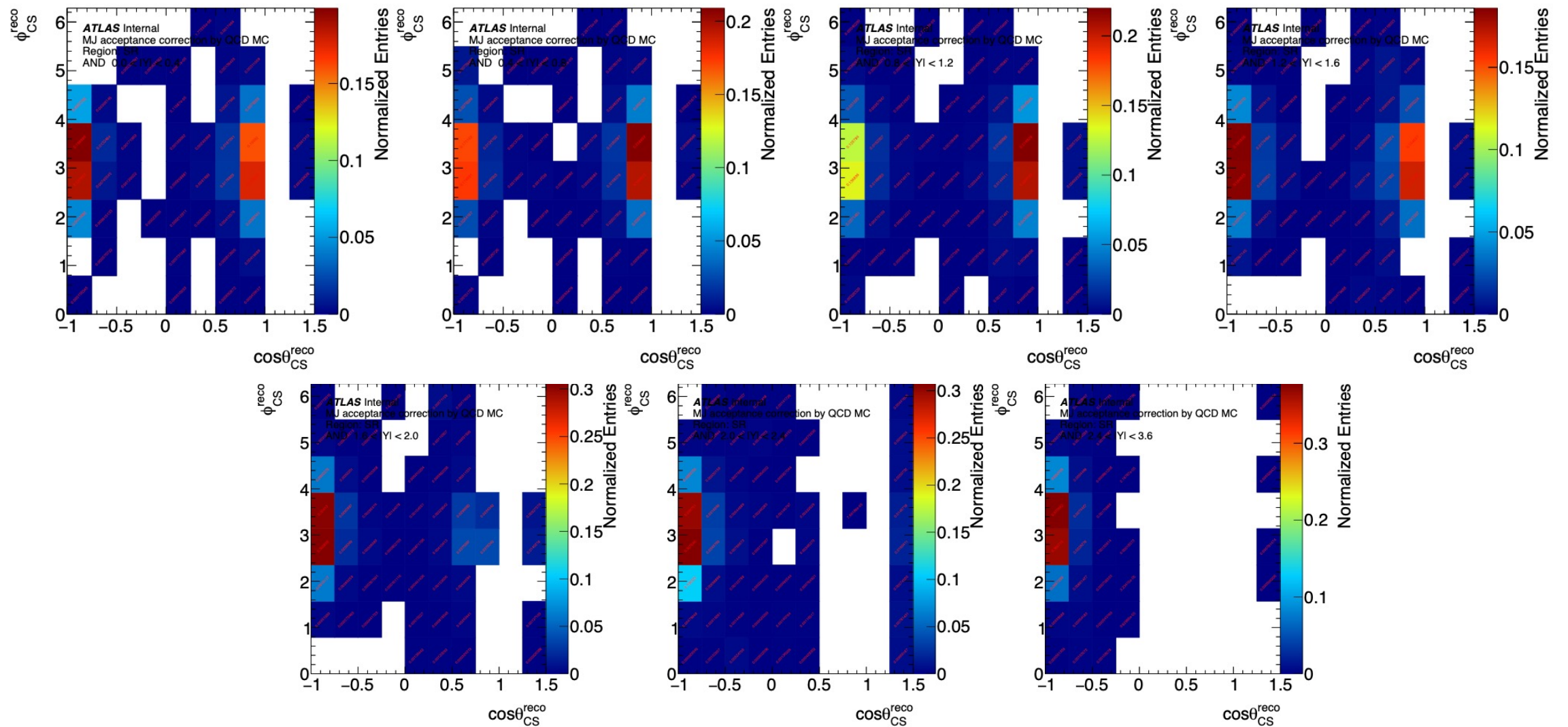
MJ templates for W-Ai analysis in $p_T^{\ell,\nu}$ binning

$$W^- \rightarrow e^- \bar{\nu}$$



$$p_T^{\ell,\nu} : [0, 8, 17, 27, 40, 55, 75, 110, 150, 210, 600];$$

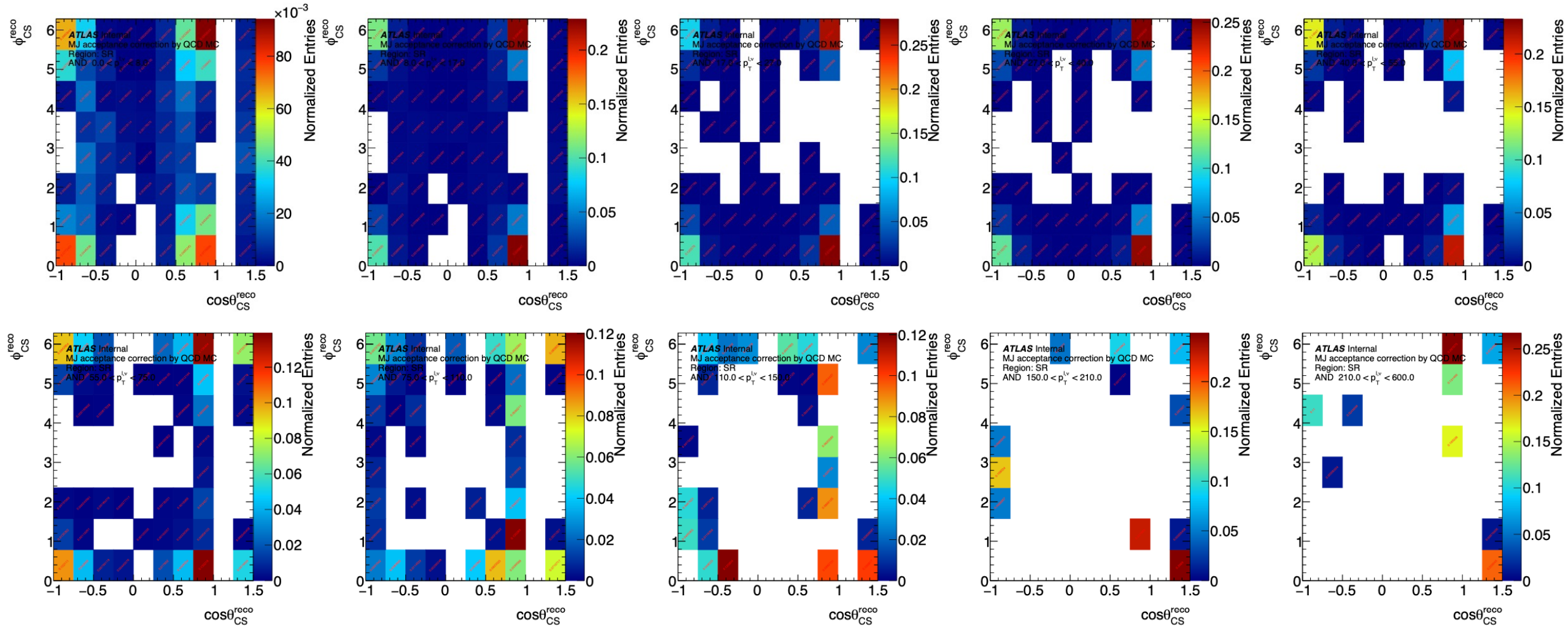
MJ templates for W-Ai analysis in $|Y|$ binning $W^- \rightarrow e^- \bar{\nu}$



$|Y| : [0, 0.4, 0.8, 1.2, 1.6, 2.0, 2.4, 3.6]$

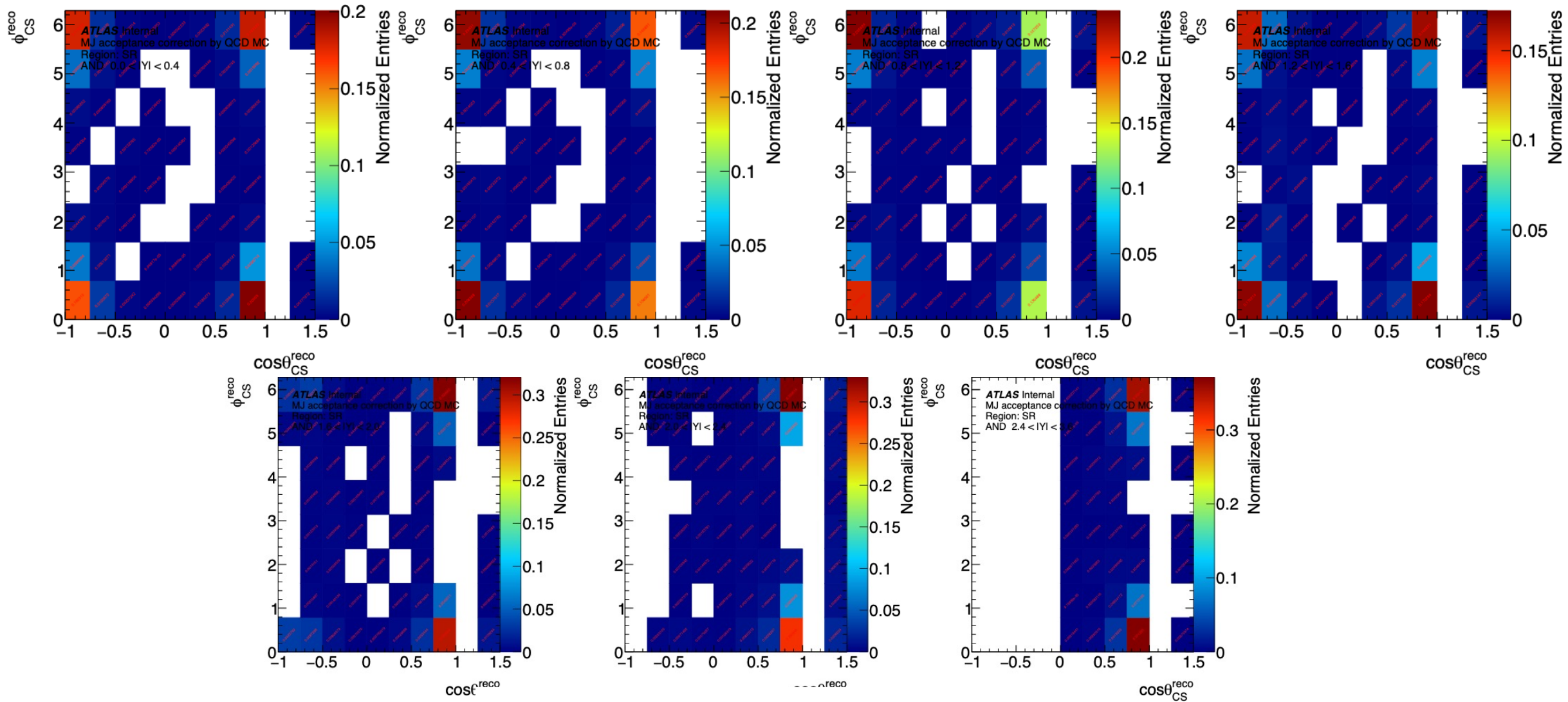
MJ templates for W-Ai analysis in $p_T^{\ell,\nu}$ binning

$$W^- \rightarrow \mu^- \bar{\nu}$$



$$p_T^{\ell,\nu} : [0, 8, 17, 27, 40, 55, 75, 110, 150, 210, 600];$$

MJ templates for W-Ai analysis in $|Y|$ binning $W^- \rightarrow \mu^- \bar{\nu}$



$|Y| : [0, 0.4, 0.8, 1.2, 1.6, 2.0, 2.4, 3.6]$

MJ systematics naming convention

Ludovica proposed a [new naming convention](#).

Since xTR produces output for each p_T^W and Y^W bin, I dropping channel and bins numbering from the naming in the WS. Still, this naming convention should be easy to restore later in AiDY.

Name	# NP in AiDY	Comments
mj_yield_cor	1	MJ yield extrapolation (bin-by-bin correlated, two-sided).
mj_shape	3 for p_T^W / 3 for Y^W	MJ shape of the given $\cos\theta_{CS} :: \phi_{CS}$ distribution in the coarse bin (bin-by bin uncorrelated, two-sided)
mj_acc_cor_closure	3 for p_T^W / 3 for Y^W	Apply difference between DD template and MC corrected template as acc corr sys uncertainty. This syst. is correlated within p_T^W and Y^W coarse bin. (bin-by-bin uncorrelated, Q: two-sided?)
mj_acc_cor_stat	10 for p_T^W / 7 for Y^W	Basically stat power of the QCD MC sample in the coarse and slicing bin used to fold MJ DD from coarse bin to the subregion. (bin-by-bin uncorrelated, two-sided)
mj_acc_cor_sys	10 for p_T^W / 7 for Y^W	Apply systematic as 100% on the acceptance correction it-slef. (bin-by-bin uncorrelated, Q: two-sided?)

Note: The "mj_yield_uncor" systematic is the MJ shape unc. in the SR for p_T^W (or Y^W) distribution. It affects yield normalisation in the given slice bin and it should be associated with 10 for p_T^W / 7 for Y^W NP. Impact of this syst. is huge, especially in the hard p_T^W region. We had agreed in July 2022 **not to use** "mj_yield_uncor" systematic in the fit.

Coarse bins: $p_T^{\ell,\nu}$: 0 – 17 GeV, 17 – 55 GeV, 55 – 600 GeV; $|Y|$: 0 – 0.8, 0.8 – 2.0, 2.0 – 3.6.

Systematic uncertainties breakdown

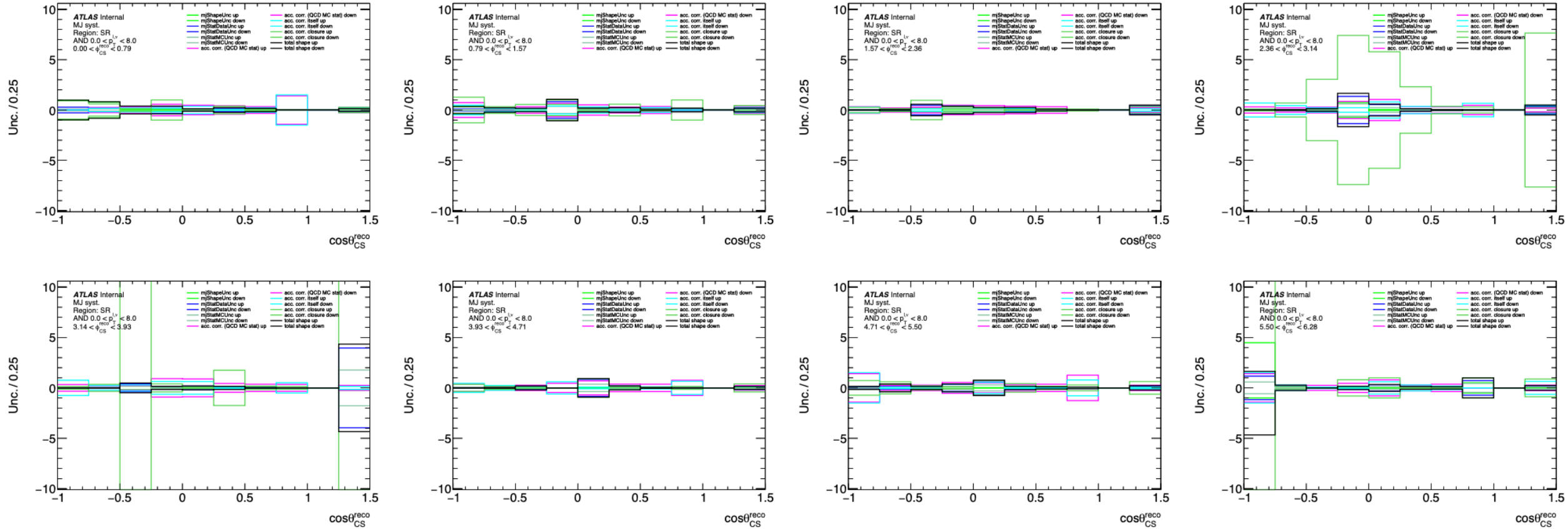


Figure 237: Multi-jet background shape systematics breakdown for $\cos \theta_{CS}$ as slices of ϕ_{CS} for $0 < p_T^{t,v} < 8$ GeV bin for $W^- \rightarrow e^- \bar{\nu}$ channel.

Systematic uncertainties breakdown

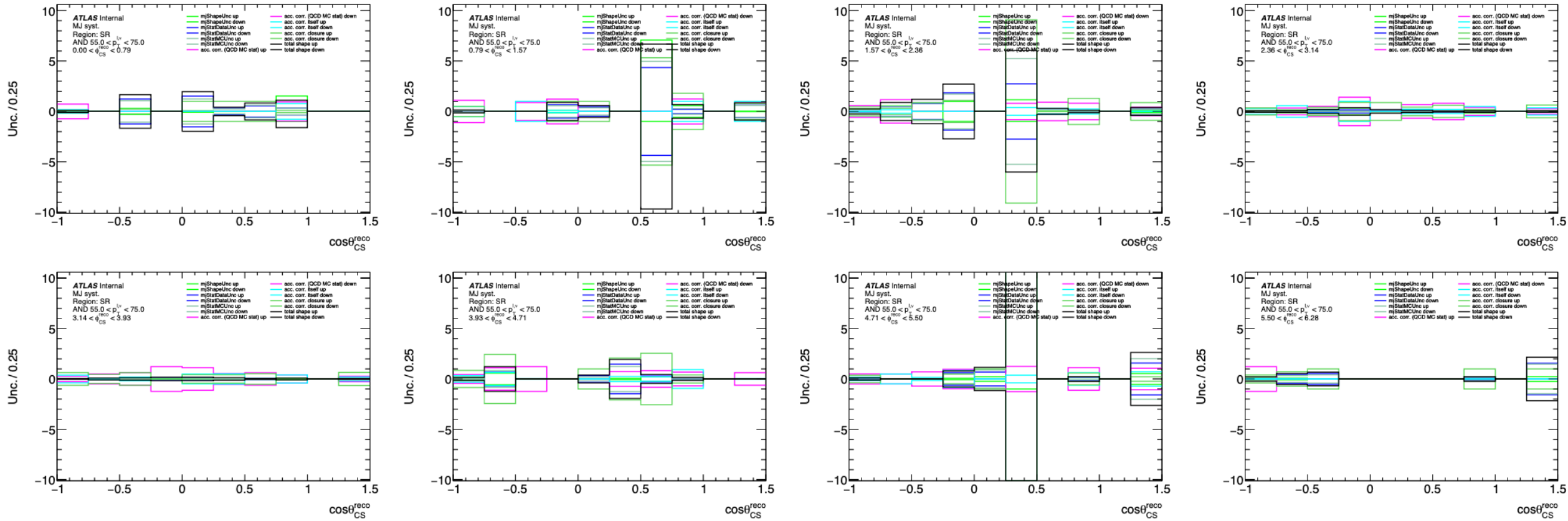


Figure 242: Multi-jet background shape systematics breakdown for $\cos \theta_{CS}$ as slices of ϕ_{CS} for $55 < p_T^{\ell, \nu} < 75$ GeV bin for $W^- \rightarrow e^- \bar{\nu}$ channel.

Systematic uncertainties breakdown

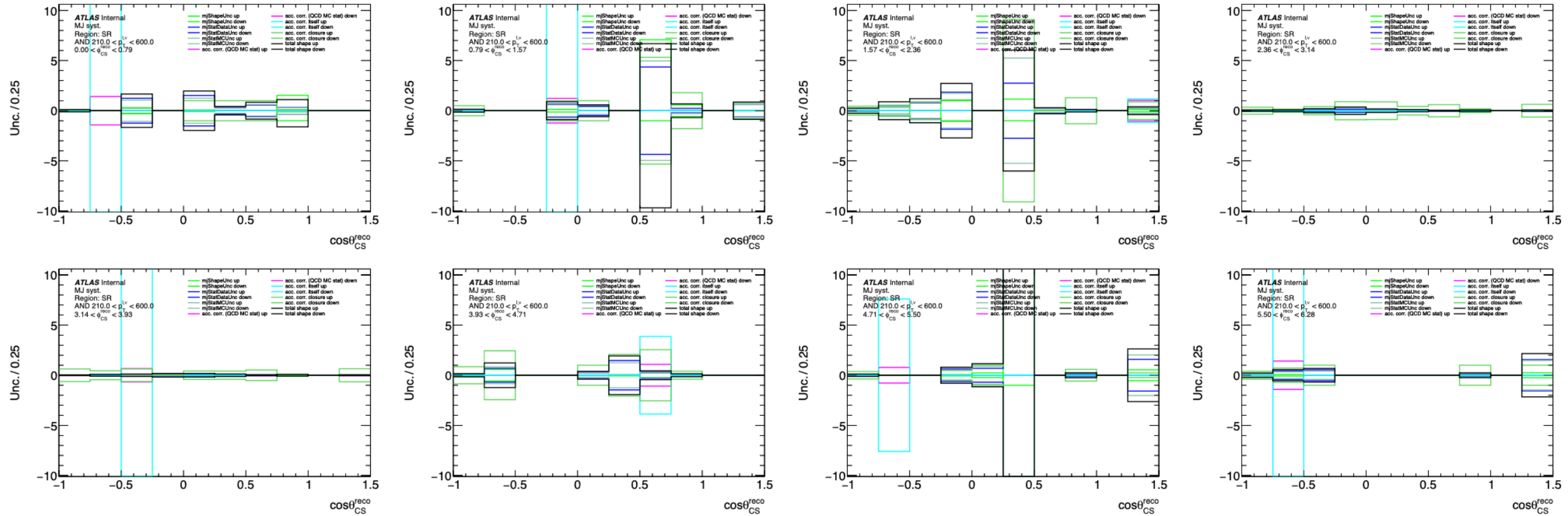
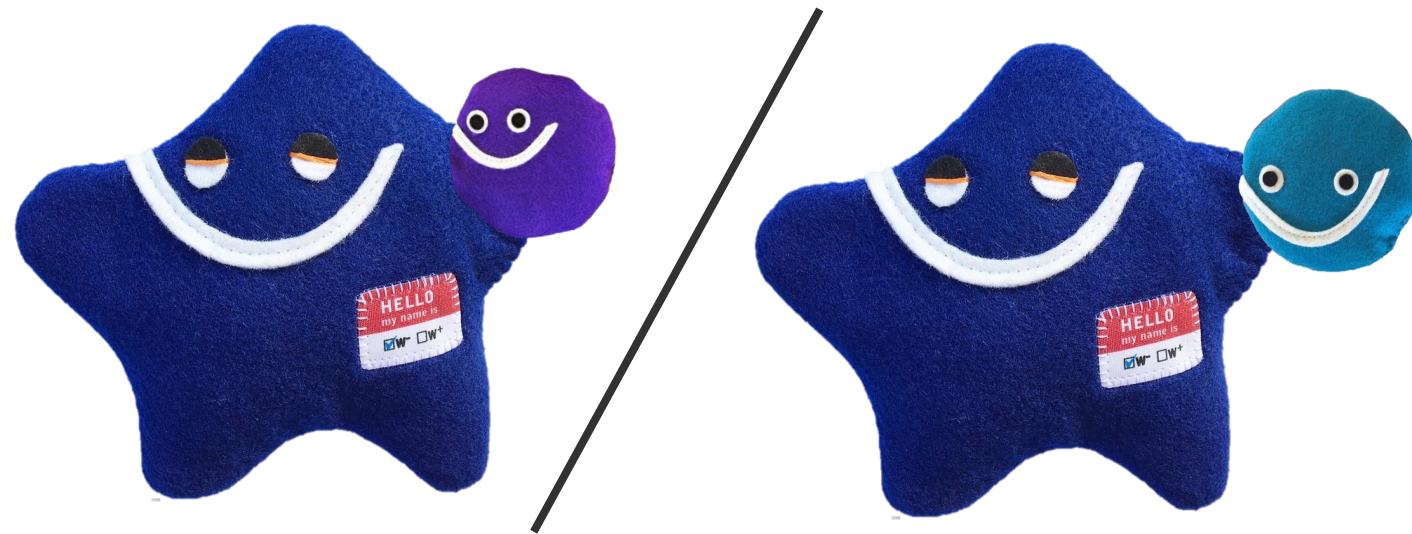


Figure 246: Multi-jet background shape systematics breakdown for $\cos \theta_{CS}$ as slices of ϕ_{CS} for $210 < p_T^{\ell, \nu} < 600$ GeV bin for $W^- \rightarrow e^- \bar{\nu}$ channel.

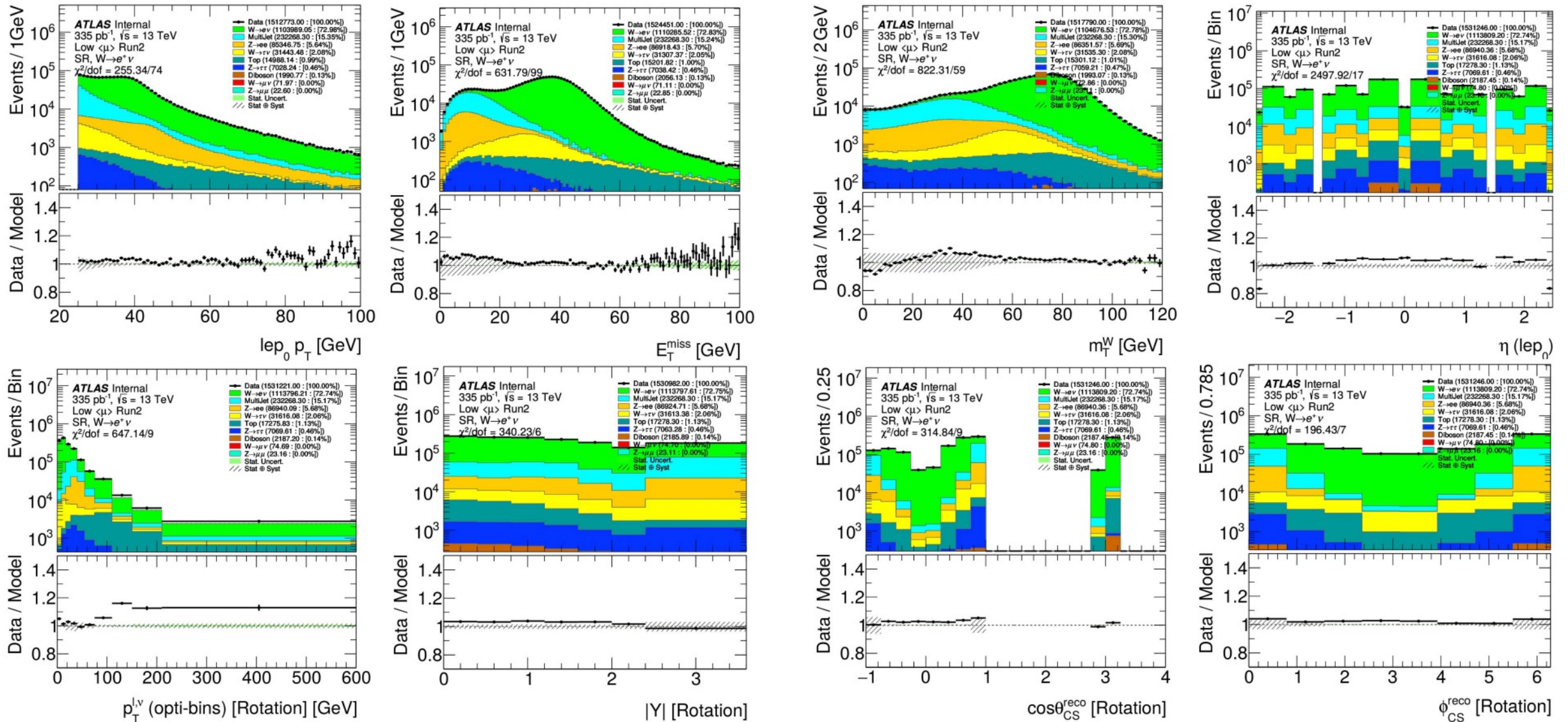
Ongoing tasks / problems

- **Fit doesn't work with real data (→ Alex)**
 - Works fine with Asmov and pseudo-data. For real data could not find minimum.
 - Perform cross-check using reco level only (Daniil)
 - Fit stability
 - MJ systematics constrains
- **New NTuple production (→ Daniil)**
 - Apply new PTRW for $p_{TW} > 100$ GeV
 - Apply additional ID and iso SF systematics
 - Simplify systematics production
- **PDF studies (→ Grigorii)**
 - Theoretical predictions for multiple PDF sets and compare PDF uncertainties
 - CT10NLO, CT18NNLO, NNPDF4.0NNLO, MMHT20

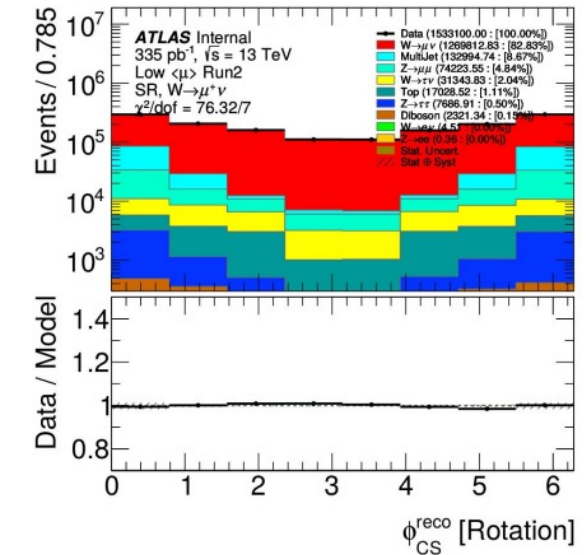
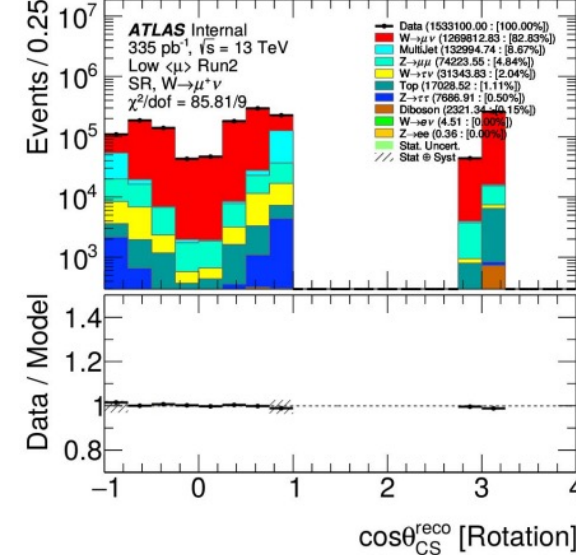
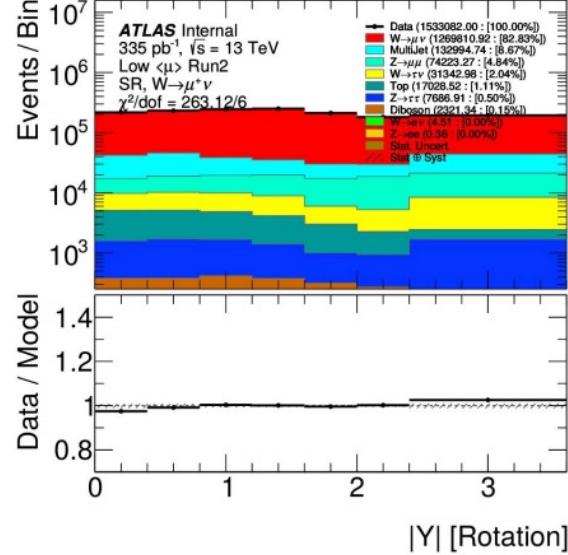
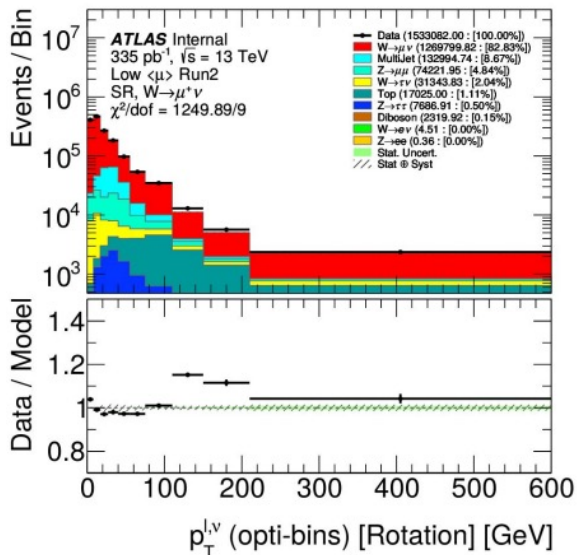
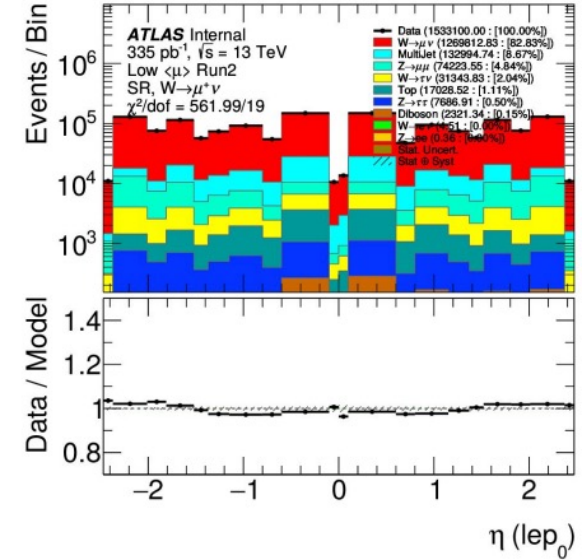
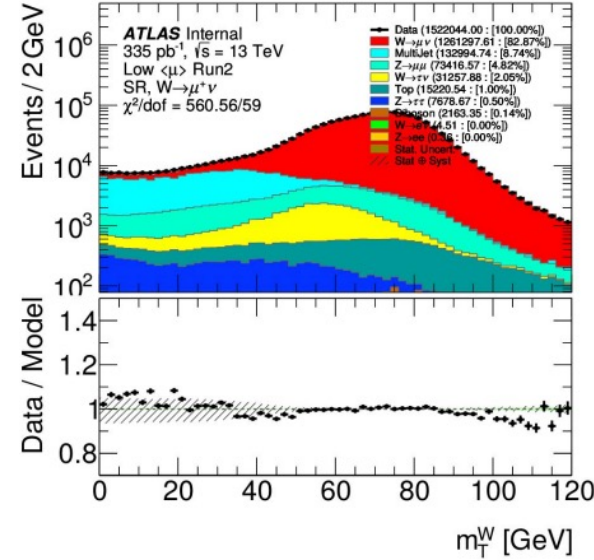
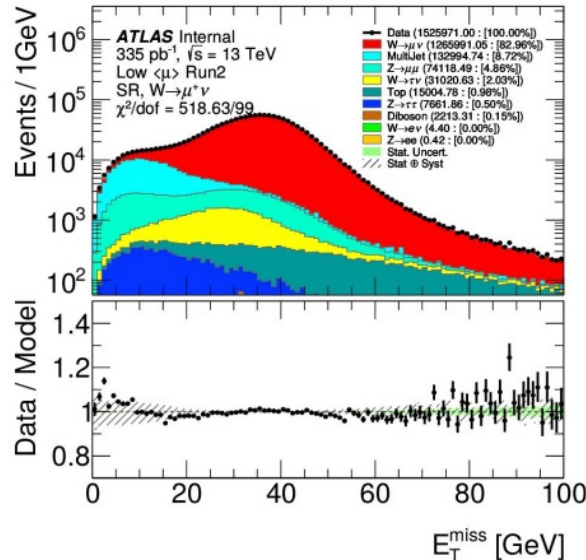
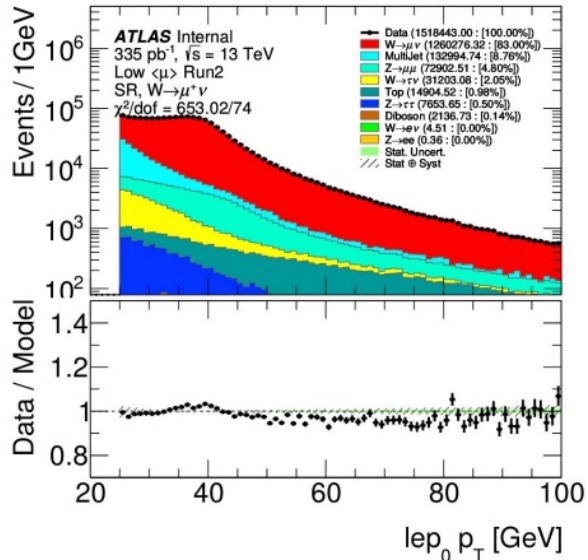


Thanks for attention!

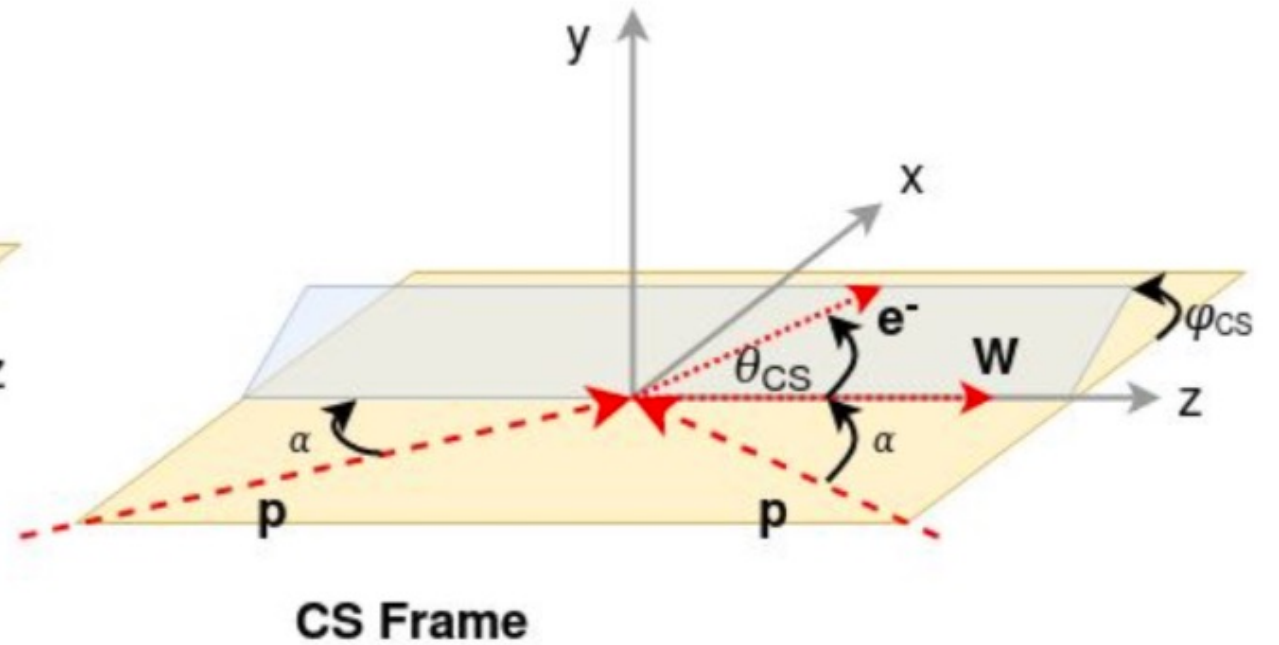
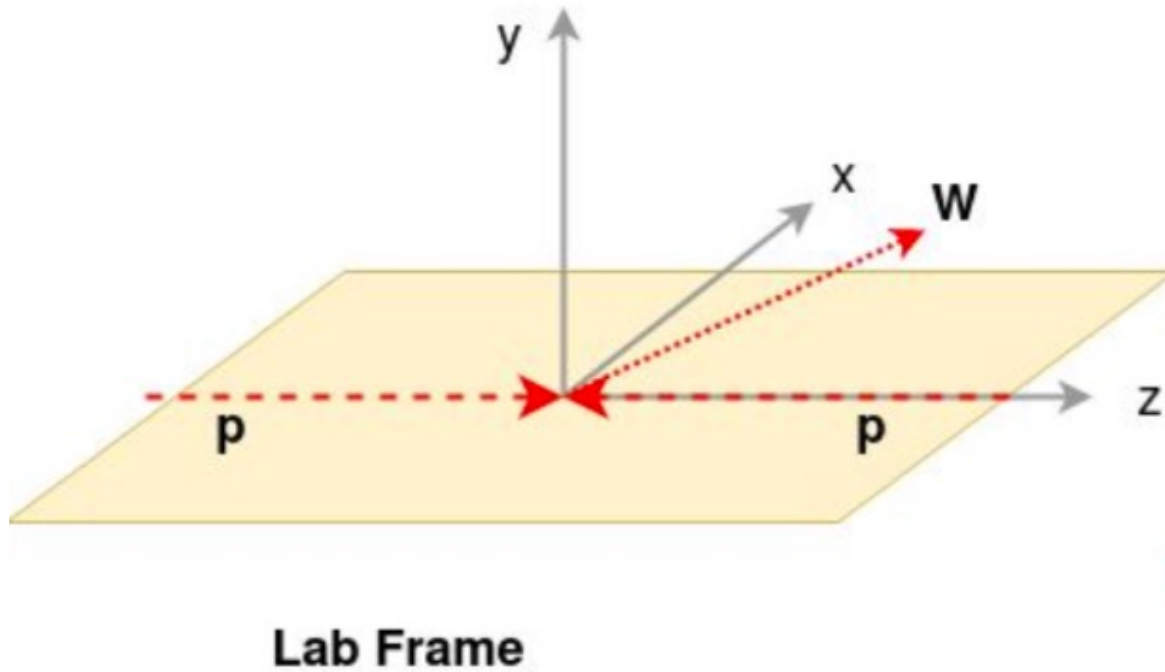
Control Plots in the Signal Region for $W^+ \rightarrow e^+ \nu$



Control Plots in the Signal Region for $W^+ \rightarrow \mu^+ \nu$



Collins Soper Frame

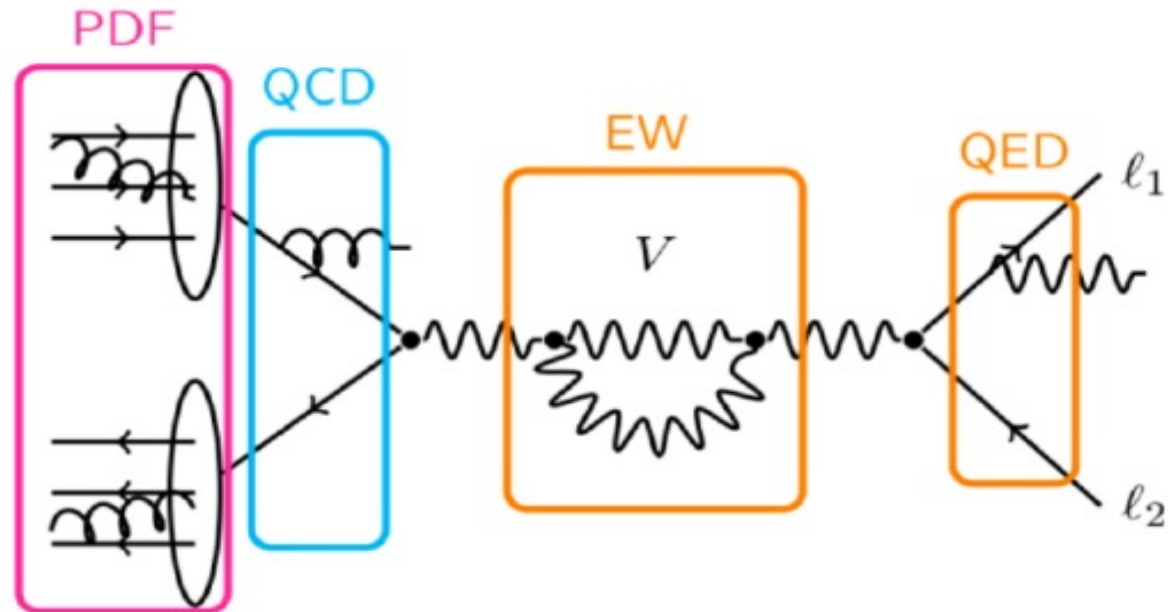


- Special W boson rest frame where angles are defined by lepton and proton kinematics.
- Define using “negative” lepton, so for W^+ case this is the neutrino

Polynomial Information

\mathbf{A}_i	\mathbf{P}_i	$\mathbf{Y}_1^m(\theta, \phi)$	Coupling	Non-Zero
A_0	$1/2 (1 - 3 \cos^2 \theta)$	Y_2^0		$O(\alpha_S^1)$
A_1	$\sin 2\theta \cos \phi$	$(Y_2^{-1} - Y_2^1)$	$(v_\ell^2 + a_\ell^2) (v_q^2 + a_q^2)$	$O(\alpha_S^1)$
A_2	$1/2 \sin^2 \theta \cos 2\phi$	$(Y_2^{-2} + Y_2^2)$		$O(\alpha_S^1)$
A_3	$\sin \theta \cos \phi$	$(Y_1^{-1} - Y_1^1)$	$v_\ell a_\ell v_q a_q$	$O(\alpha_S^1)$
A_4	$\cos \theta$	Y_1^0		$O(\alpha_S^0)$
A_5	$\sin^2 \theta \sin 2\phi$	$(Y_2^{-2} - Y_2^2)$	$(v_\ell^2 + a_\ell^2) (v_q a_q)$	$O(\alpha_S^2)$
A_6	$\sin 2\theta \sin \phi$	$(Y_2^{-1} + Y_2^1)$		$O(\alpha_S^2)$
A_7	$\sin \theta \sin \phi$	$(Y_1^{-1} + Y_1^1)$	$(v_\ell a_\ell) (v_q^2 + a_q^2)$	$O(\alpha_S^2)$

Physics Modelling



Angular Coefficients

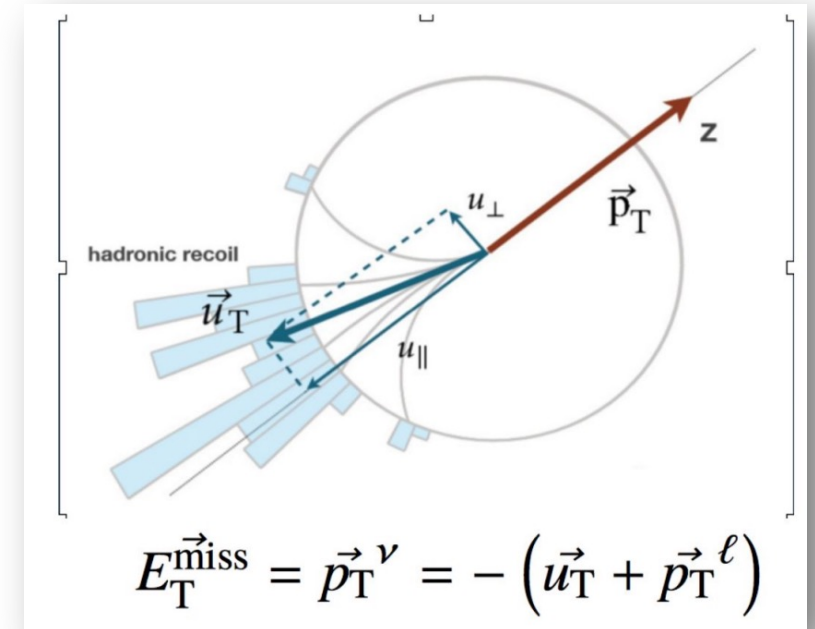
$$\frac{d\sigma}{dp_T^W dm^W dy^W d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_T^W dm^W dy^W} \left\{ \begin{aligned} & \left(1 + \cos^2\theta\right) + \frac{1}{2}A_0 \left(1 - 3\cos^2\theta\right) \\ & + A_1 \sin 2\theta \cos\phi + \frac{1}{2}A_2 \sin^2\theta \cos 2\phi \\ & + A_3 \sin\theta \cos\phi + A_4 \cos\theta + A_5 \sin^2\theta \sin 2\phi \\ & + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \end{aligned} \right\}.$$

Non-zero at LO
 Non-zero at NLO
 Non-zero at NNLO

- Full angular cross-section parameterization
- QCD production dynamics fully contained in A_i coefficients, while decay kinematics fully contained in coupled angular polynomials

Neutrino reconstruction

- Angular definitions require fully constructed neutrino
- **Hadronic recoil:**
 - Vectorial sum of all transverse momenta of ISR objects
 - ATLAS: uses PFlow objects (neutral+charged)
- **Solving for Neutrino p_z :**
 - Take mass constraint resulting in quadratic equation
 - $(q_l + q_\nu)^\mu (q_l + q_\nu)_\mu = m_W^2$
 - $m_T < m_W$ gives 2 solutions
 - one chosen at random but can statistically resolve correct distributions
 - No real solution when $m_T > m_W$, still can use events for φ_{CS} information
- Once the mass constraint is chosen φ_{CS} is solved while the 2 solutions corresponds to a sign ambiguity in $\cos\theta_{CS}$



$\cos\theta_{CS}$ ambiguity (1)

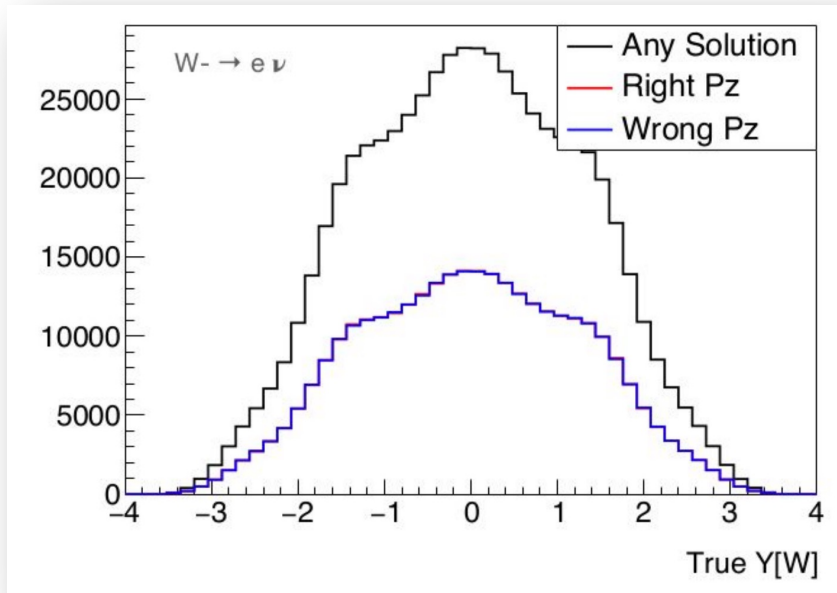
- How can we statistically resolve $\cos\theta_{CS}$?

$$\cos\theta = 2 \frac{(l^+ \bar{l}^- - l^- \bar{l}^+) \cdot \text{sign}(y_{\ell\ell})}{m_{\ell\ell} \sqrt{m_{\ell\ell}^2 + \vec{p}_T^2}}$$

$l^\pm = E \pm p_z$ for the lepton/neutrino

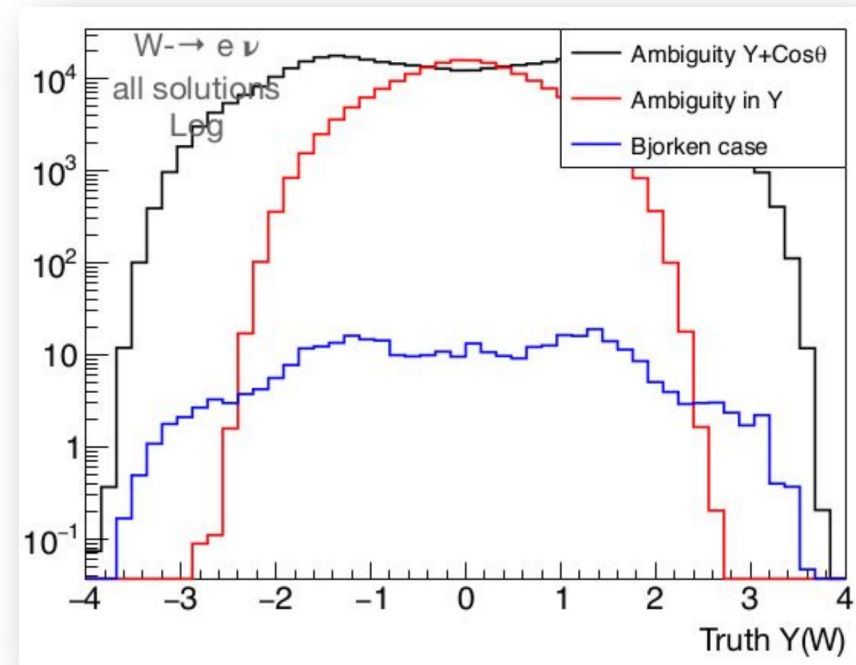
$\bar{l}^\pm = E \pm p_z$ for the antilepton/antineutrino

- Only choose correct p_z 50% of the time, incorrect p_z can still be used



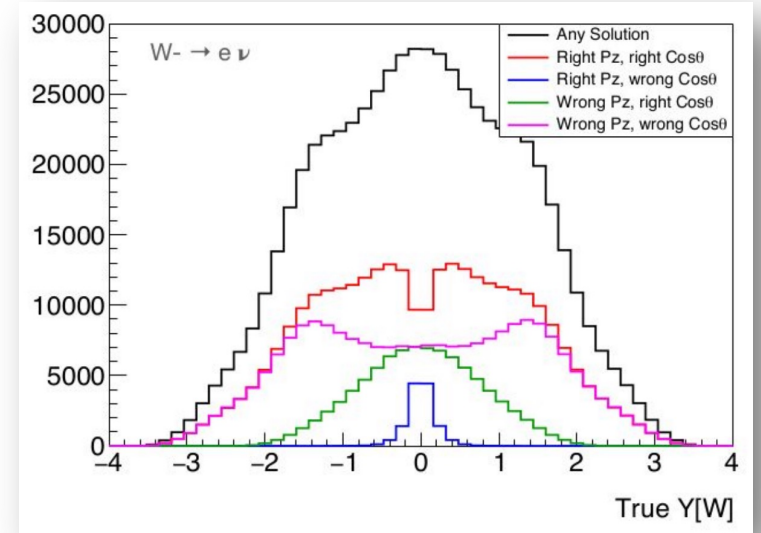
- Can solve it in 2 cases:

- 1) Solutions where y_{ll}^1 and y_{ll}^2 have opposite signs
 - sign flips in orange and blue cancel
- 2) One of the solutions violates Bjorken condition $x < 1$
 - neutrino has more p_z than beam energy \rightarrow unphysical



$\cos\theta_{CS}$ ambiguity (2)

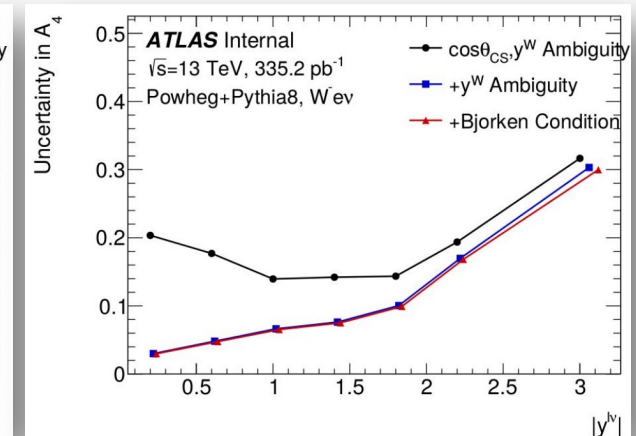
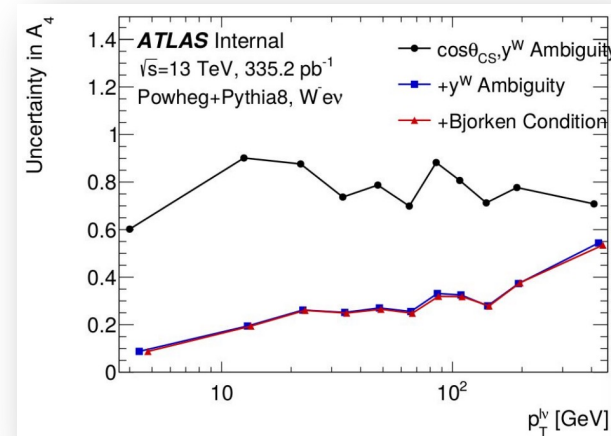
- Red are events where we picked right p_z to get right sign of $\cos\theta_{CS}$
- Green are events where we picked wrong p_z but y_U flipped so we get correct $\cos\theta_{CS}$, these are the events where we get our sensitivity.



1) Both $\cos\theta_{CS}$ and y_U have sign ambiguity so right 50%, no A_4 sensitivity

2) Adding to 1): only y^W has ambiguity which cancels with $\cos\theta_{CS} \rightarrow$ right 100% (where we gain our sensitivity)

3) Adding to 2): One solution doesn't satisfy Bjorken condition ($x < 1$) \rightarrow right 100% but pretty rare

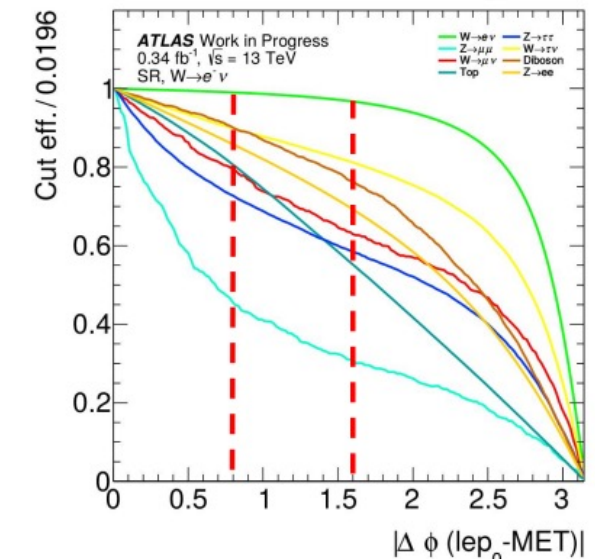
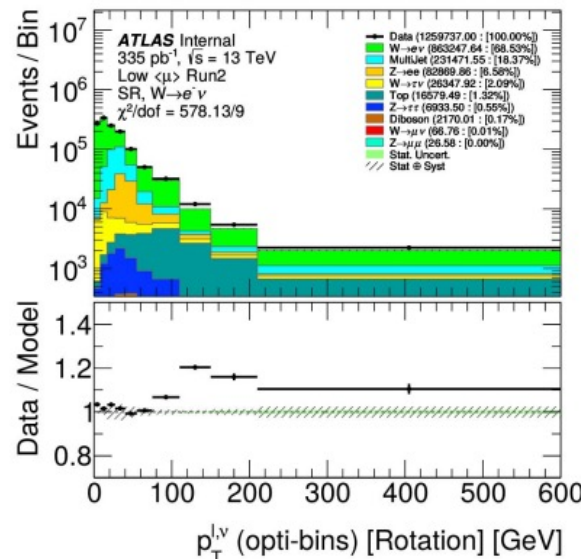
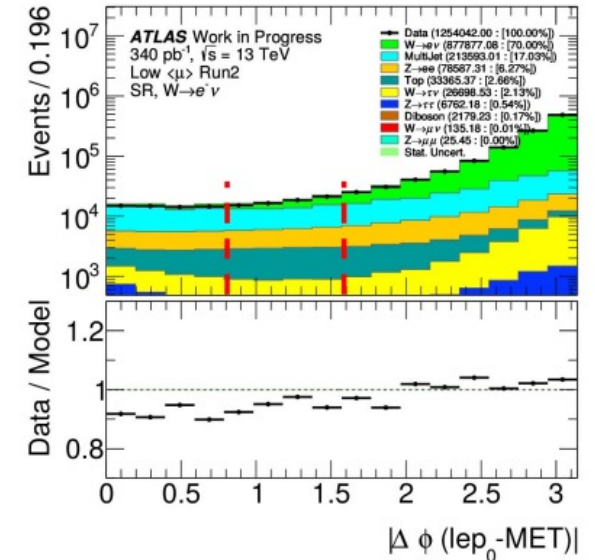
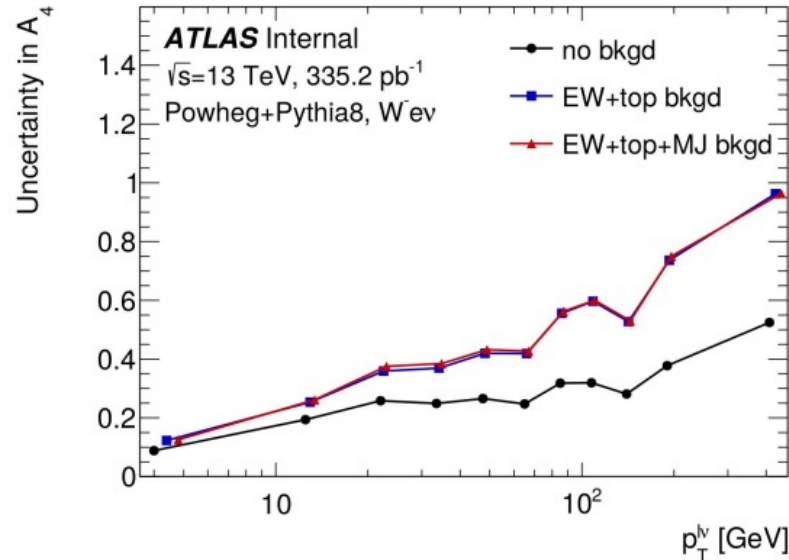


Top Background (1)

When first adding background, the increase in uncertainty was larger than expected especially at high $p_T^{l,\nu}$.

Due to larger top background contributions.

Preliminarily we thought a $|\Delta\phi| > \pi/4$ cut should help reduce the top background but doesn't as shown on the next slide.

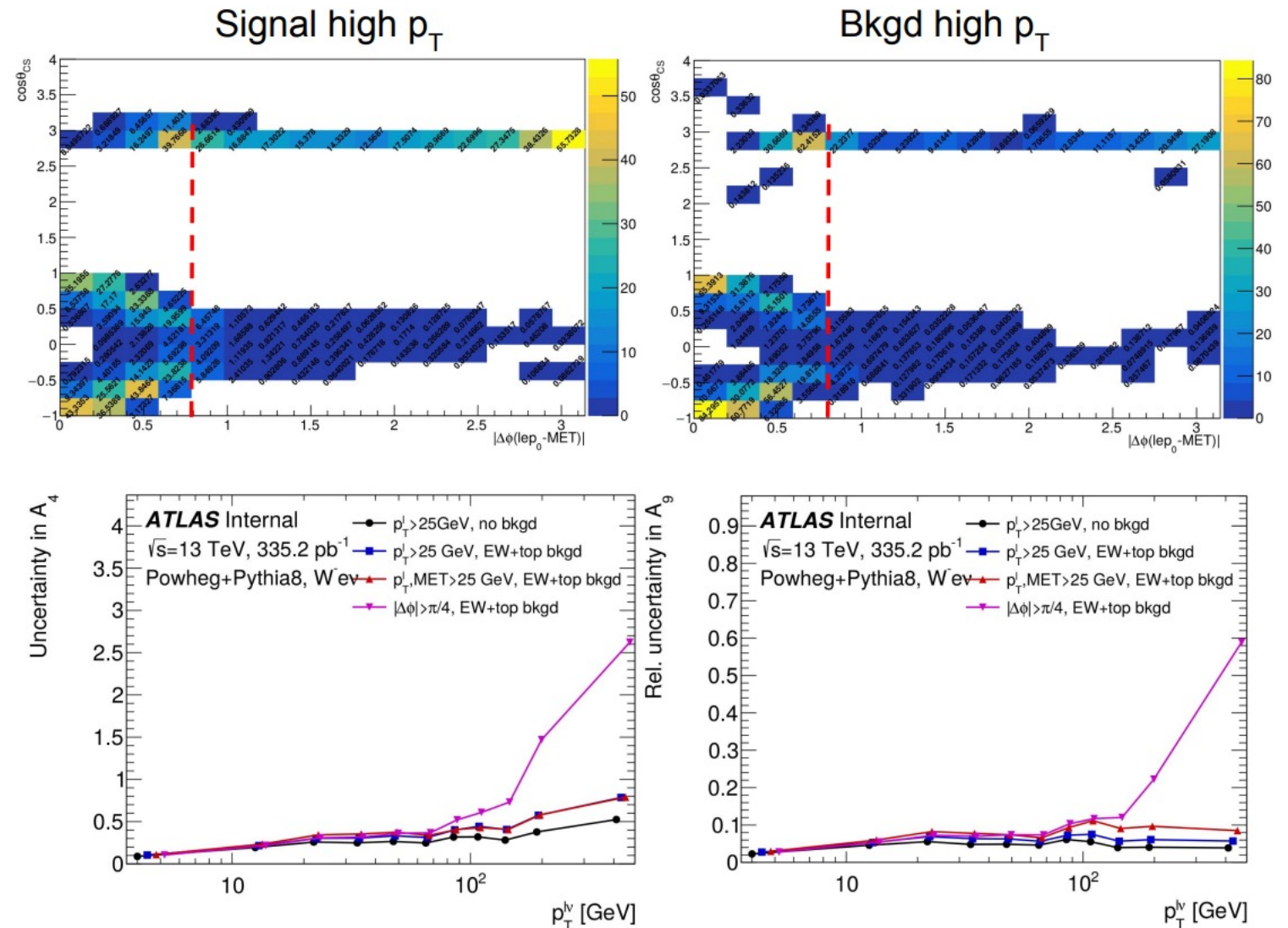


Top Background (2)

Both signal and background at high p_T lose all $\cos\theta_{CS}$ shape information and causing large uncertainty.

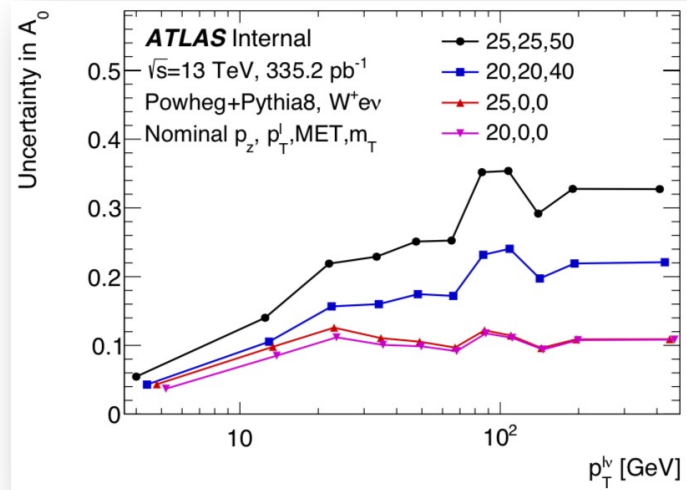
Addition of MET cut reduces background but also signal too much.

From our studies the top background seems irreducible.

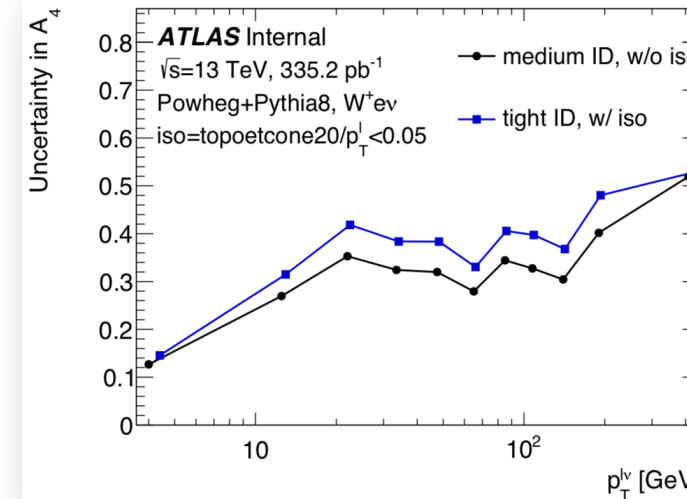
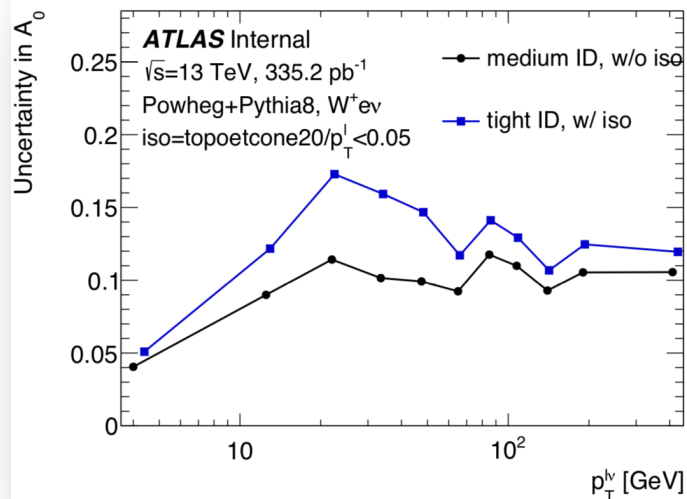
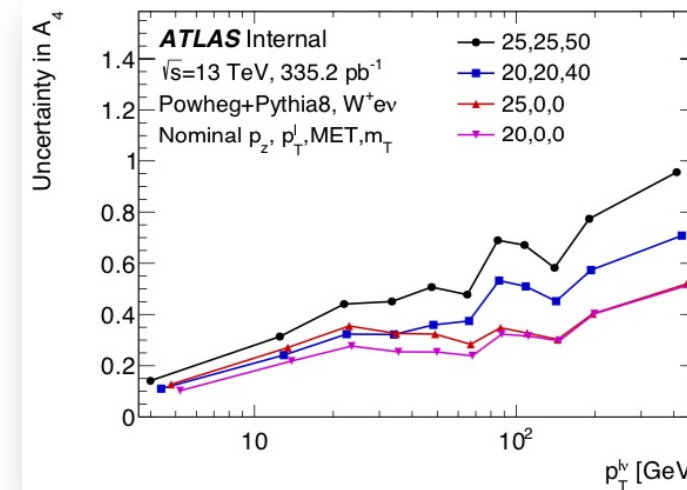


Event selection: Sensitivity Comparison

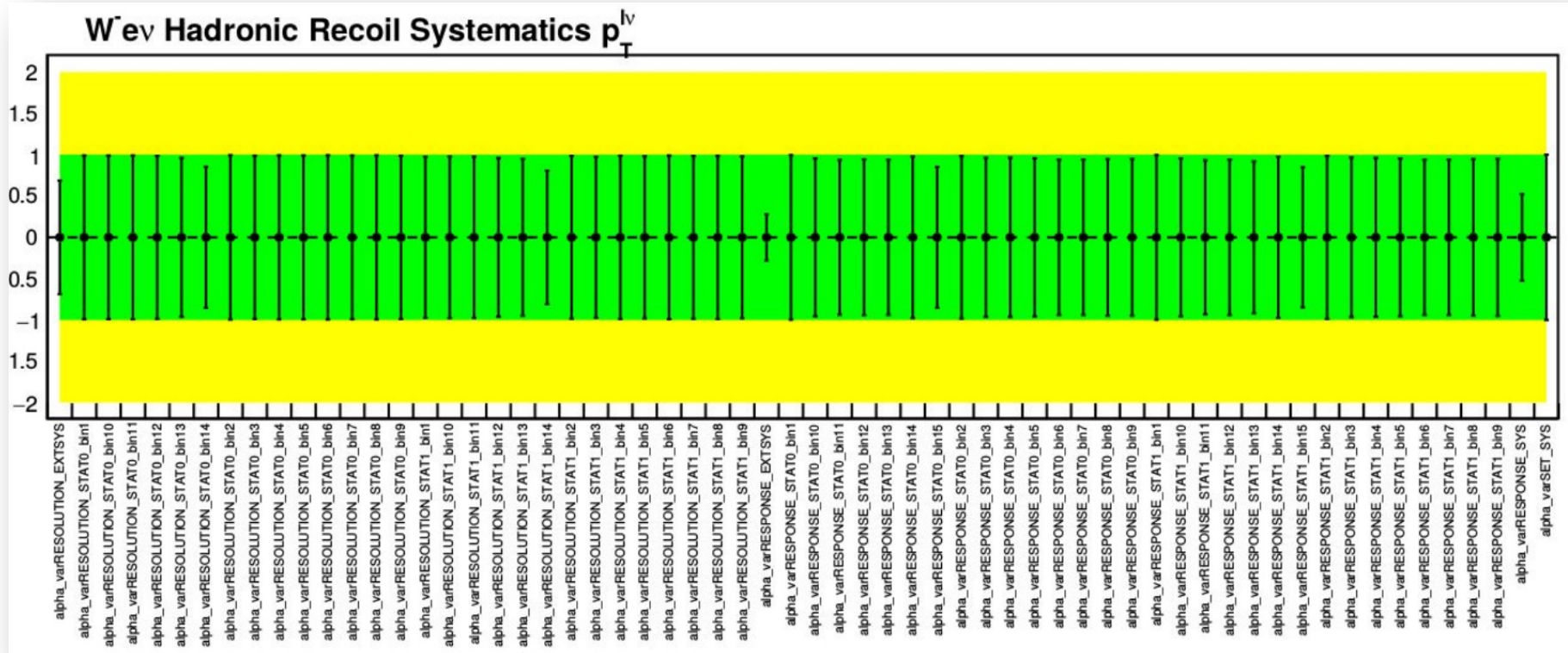
A_0



A_4

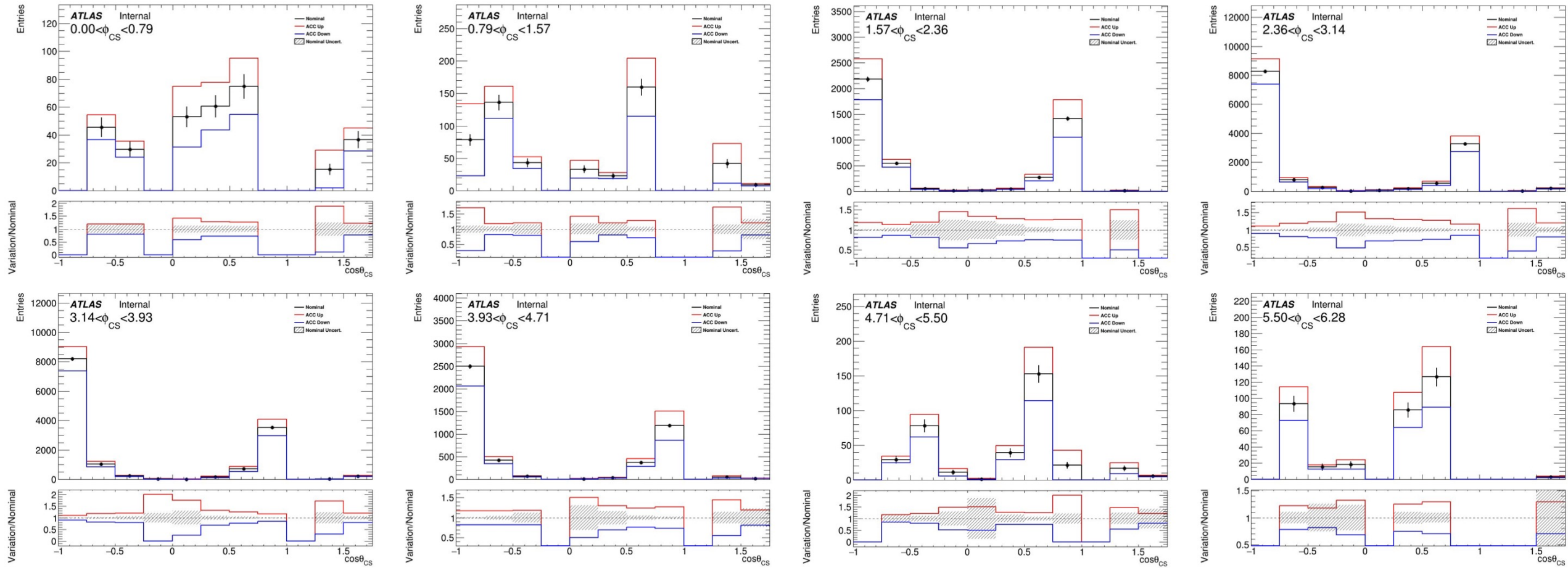


Systematics

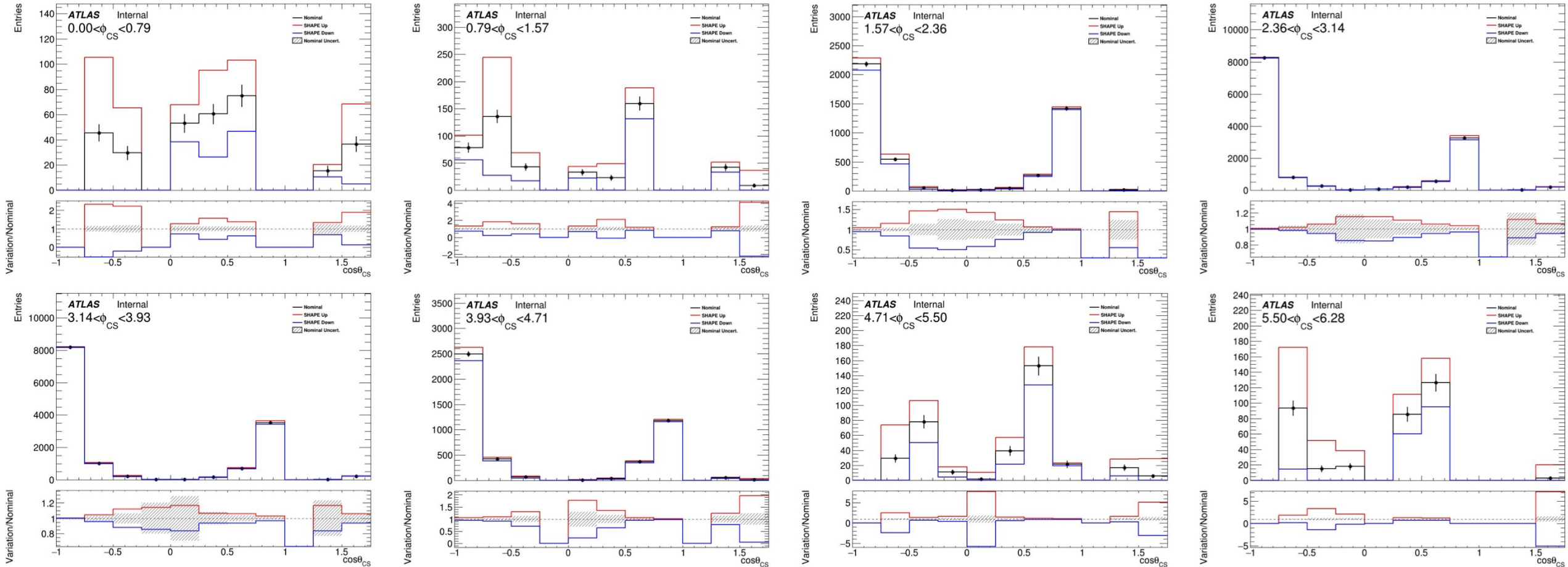


- Have tested 2 largest expected systematics (Hadronic Recoil and Background) independently
- MJ NPs heavily constrained, still being understood.

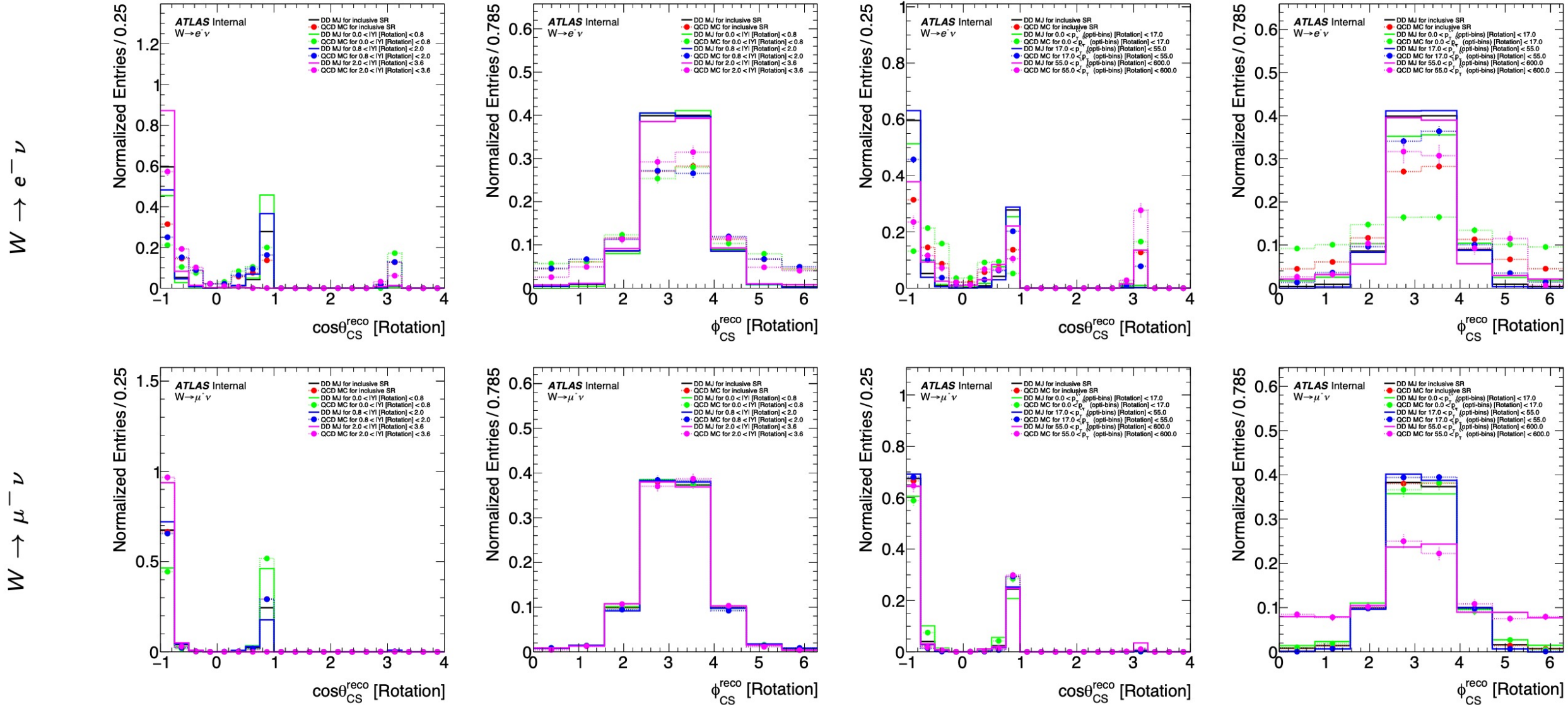
MJ Acceptance Systematic ($8 < p_T < 17$ GeV)



MJ Shape Systematic ($8 < p_T < 17$ GeV)

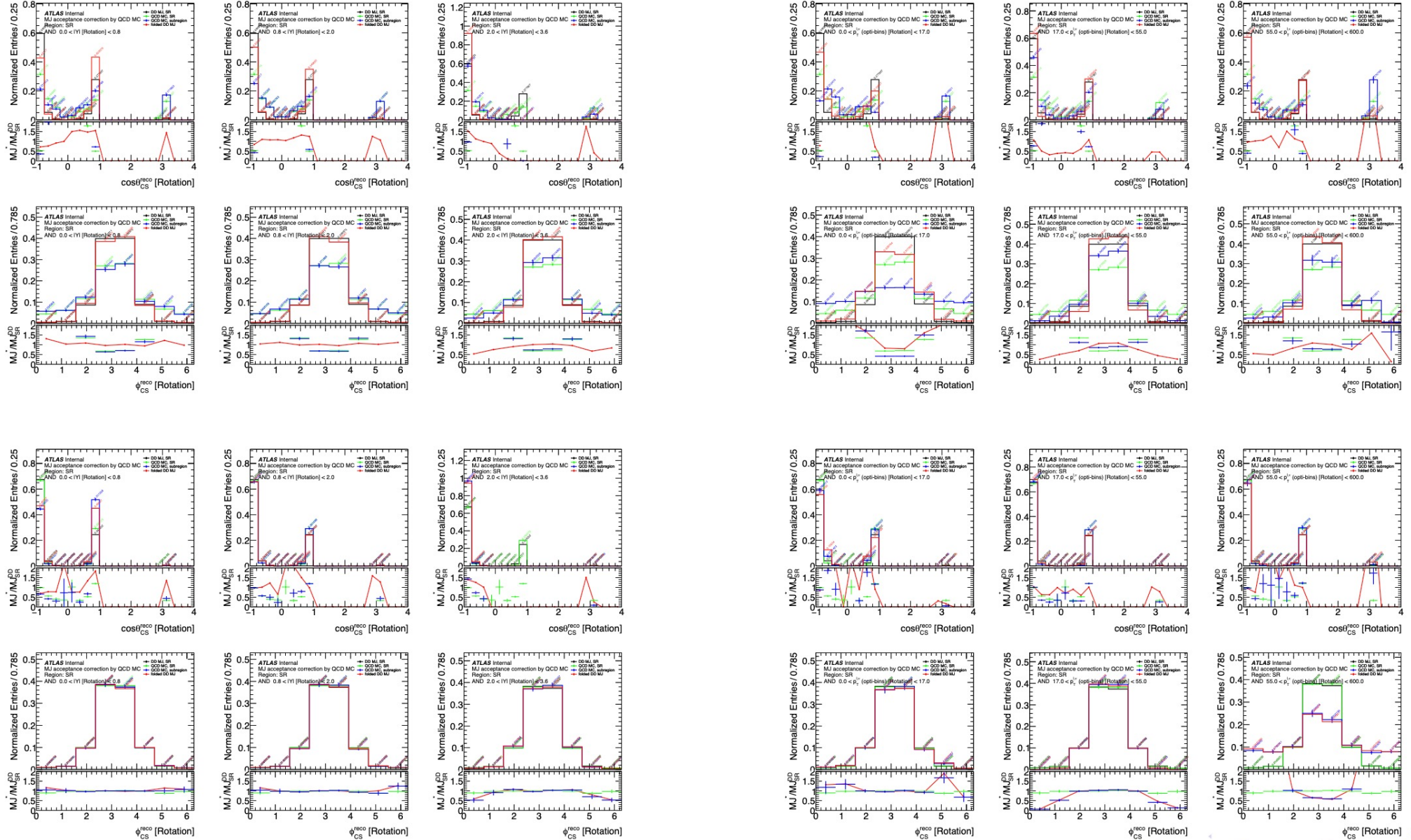


MJ templates derived for pure DD and predicted by QCD MC



Debugging: MJ templates derived by accCorr

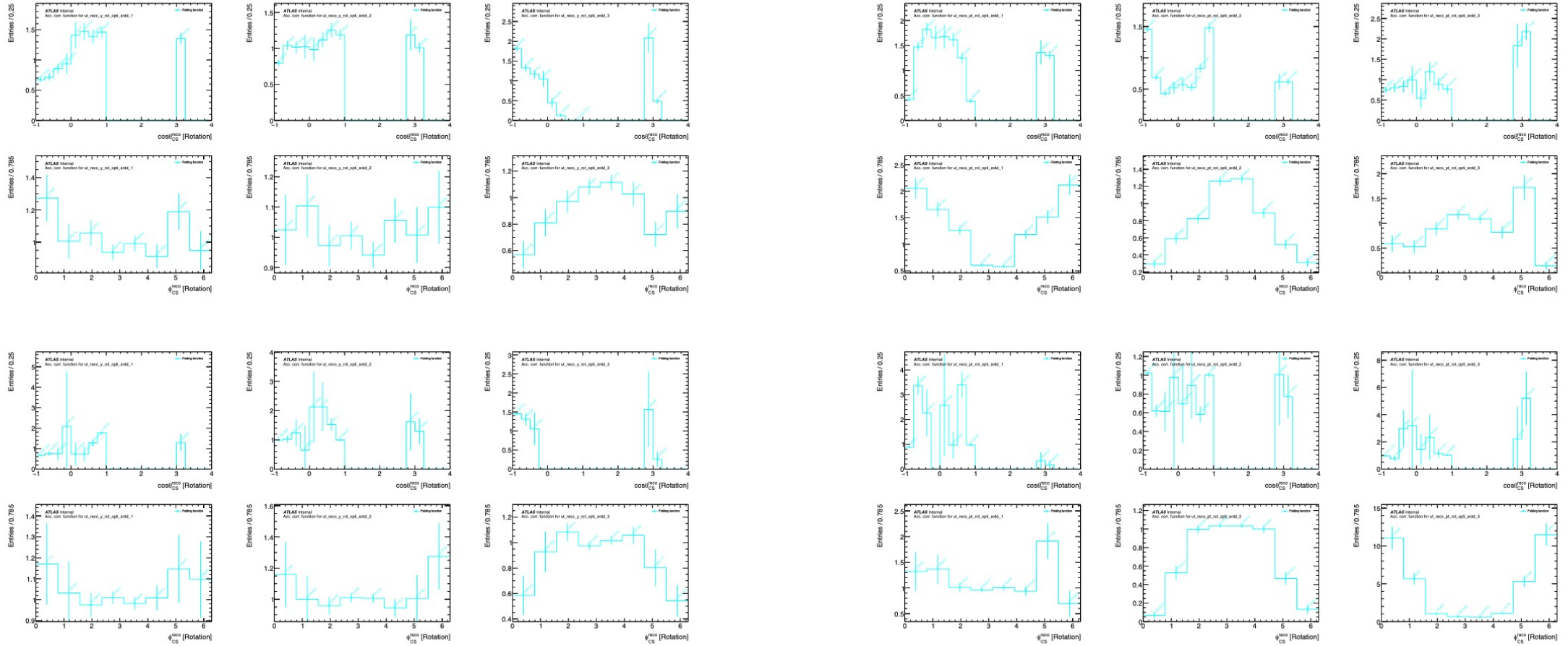
$W \rightarrow e^- \nu$



$W \rightarrow \mu^- \nu$

Debugging: accCorr functions

$W \rightarrow e^- \nu$



$W \rightarrow \mu^- \nu$

Acc Correction for 3D templates

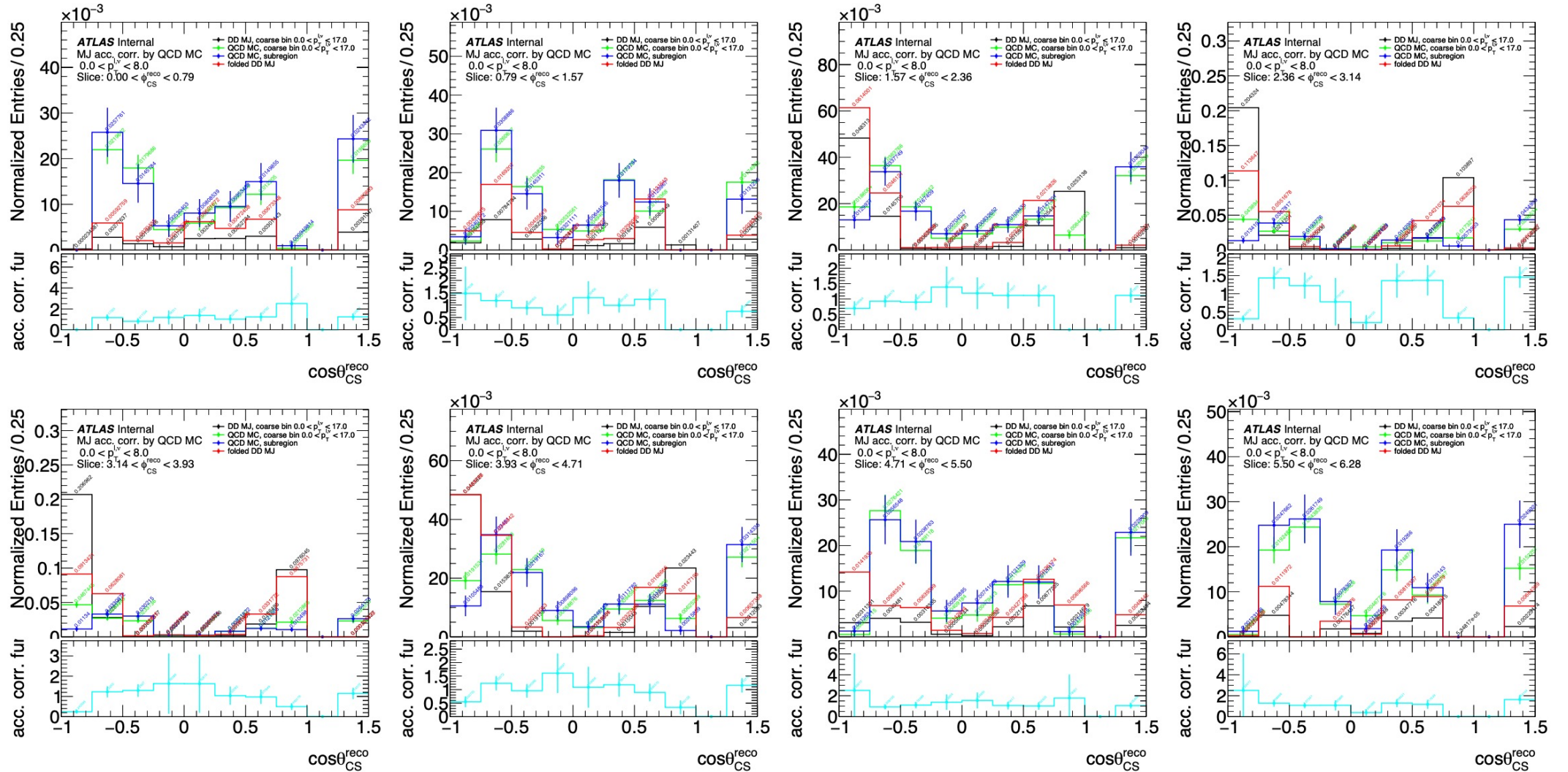


Figure 161: Multi-jet background template acceptance correction functions for $\cos\theta_{CS}$ as slices of ϕ_{CS} for $0 < p_T^{\ell,\nu} < 8$ GeV bin for $W^- \rightarrow e^- \bar{\nu}$ channel. Distributions are normalized over ϕ_{CS} slices. Error bands represent statistical uncertainties.

Acc Correction for 3D templates

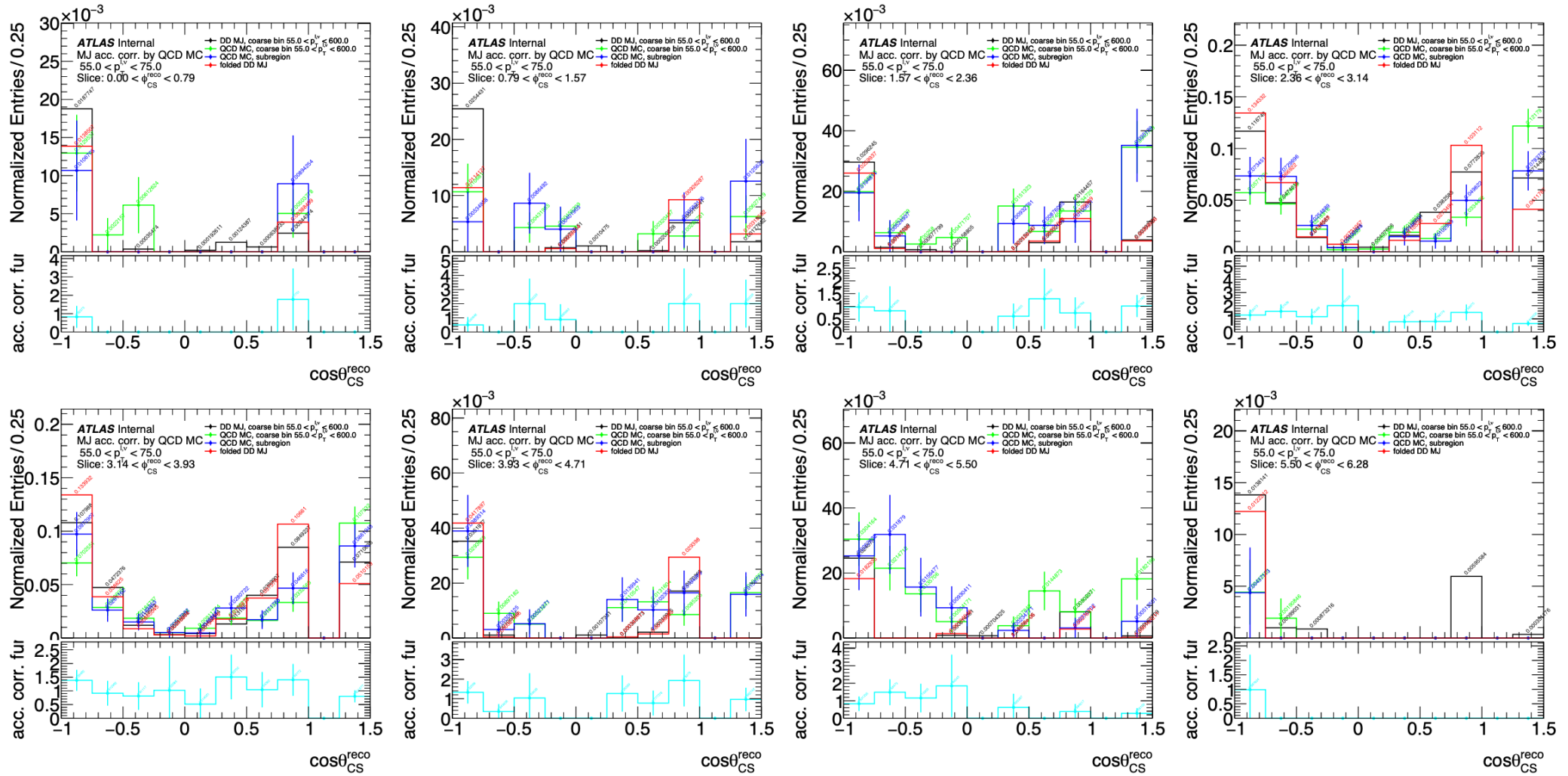


Figure 166: Multi-jet background template acceptance correction functions for $\cos\theta_{CS}$ as slices of ϕ_{CS} for $55 < p_T^{\ell,\nu} < 75$ GeV bin for $W^- \rightarrow e^- \bar{\nu}$ channel. Distributions are normalized over ϕ_{CS} slices. Error bands represents statistical uncertainties.

Acc Correction for 3D templates

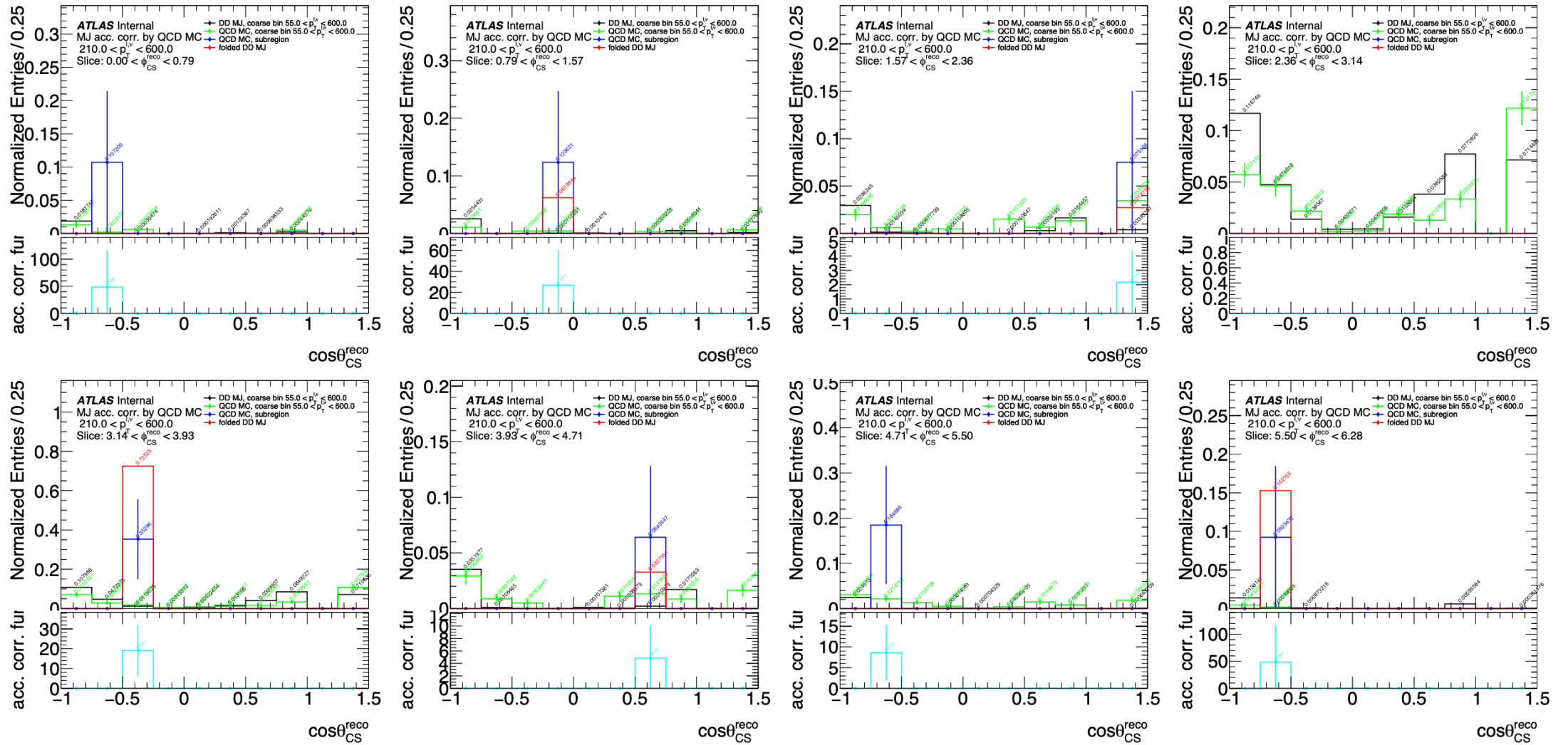


Figure 170: Multi-jet background template acceptance correction functions for $\cos\theta_{CS}$ as slices of ϕ_{CS} for $210 < p_T^{\ell,\nu} < 600$ GeV bin for $W^- \rightarrow e^- \bar{\nu}$ channel. Distributions are normalized over ϕ_{CS} slices. Error bands represents statistical uncertainties.