

Status of Light Sterile Neutrinos

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Standard Three Neutrino Mixing

- ▶ Flavor Neutrinos: ν_e, ν_μ, ν_τ produced in Weak Interactions
- ▶ Massive Neutrinos: ν_1, ν_2, ν_3 propagate from Source to Detector
- ▶ Neutrino Mixing: a Flavor Neutrino is a **superposition** of Massive Neutrinos

$$\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \\ |\nu_\tau\rangle \end{pmatrix} = \begin{pmatrix} U_{e1}^* & U_{e2}^* & U_{e3}^* \\ U_{\mu 1}^* & U_{\mu 2}^* & U_{\mu 3}^* \\ U_{\tau 1}^* & U_{\tau 2}^* & U_{\tau 3}^* \end{pmatrix} \begin{pmatrix} |\nu_1\rangle \\ |\nu_2\rangle \\ |\nu_3\rangle \end{pmatrix}$$

- ▶ U is the 3×3 unitary **Neutrino Mixing Matrix**

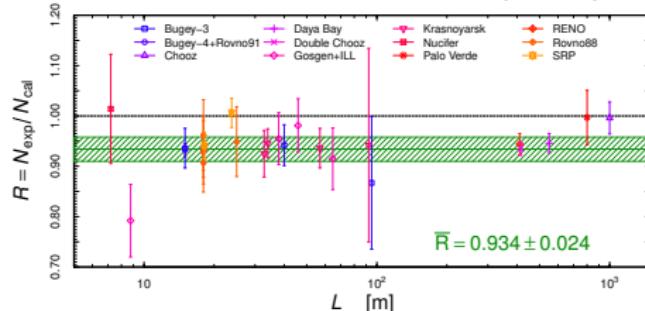
- ▶ $P_{\nu_\alpha \rightarrow \nu_\beta}(L) = \sum_{k,j} U_{\beta k} U_{\alpha k}^* U_{\beta j}^* U_{\alpha j} \exp\left(-i \frac{\Delta m_{kj}^2 L}{2E}\right)$ ($\alpha, \beta = e, \mu, \tau$)
- ▶ The oscillation probabilities depend on

U (osc. amplitude) and $\Delta m_{kj}^2 \equiv m_k^2 - m_j^2$ (osc. phase)

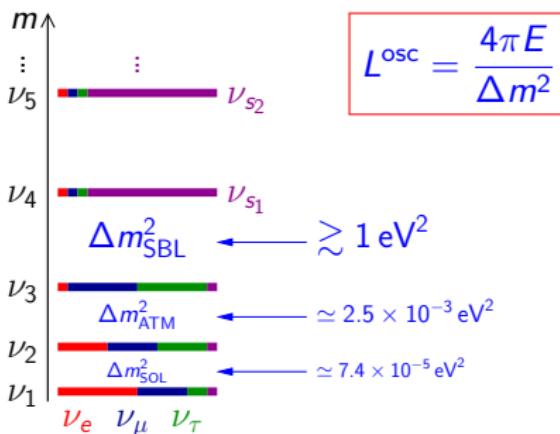
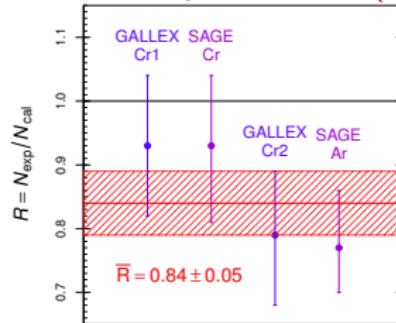
- ▶ In the standard framework of three-neutrino mixing there are two independent Δm^2 's:
 - ▶ $\Delta m_{\text{SOL}}^2 = \Delta m_{21}^2 \simeq 7.4 \times 10^{-5} \text{ eV}^2$ Solar Mass Splitting
 - ▶ $\Delta m_{\text{ATM}}^2 \simeq |\Delta m_{31}^2| \simeq 2.5 \times 10^{-3} \text{ eV}^2$ Atmospheric Mass Splitting
- ▶ The solar and atmospheric mass splittings generate oscillations that are detectable at the distances
 - ▶ $L_{\text{SOL}}^{\text{osc}} \gtrsim \frac{E_\nu}{\Delta m_{\text{SOL}}^2} \approx 50 \text{ km} \frac{E_\nu}{\text{MeV}}$
 - ▶ $L_{\text{ATM}}^{\text{osc}} \gtrsim \frac{E_\nu}{\Delta m_{\text{ATM}}^2} \approx 1 \text{ km} \frac{E_\nu}{\text{MeV}}$
- ▶ The solar and atmospheric mass splittings cannot explain flavor neutrino transitions at shorter distances.

Short-Baseline Neutrino Oscillation Anomalies

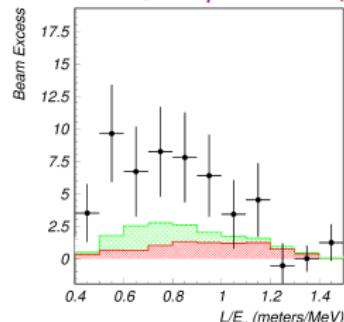
Reactor Anomaly: $\bar{\nu}_e \rightarrow \bar{\nu}_x$ ($\sim 3\sigma$)



Gallium Anomaly: $\nu_e \rightarrow \nu_x$ ($\sim 3\sigma$)



LSND Anomaly: $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ ($\sim 4\sigma$)



Minimal perturbation of 3ν mixing: effective 3+1 with $|U_{e4}|, |U_{\mu 4}|, |U_{\tau 4}| \ll 1$

Effective 3+1 SBL Oscillation Probabilities

Appearance ($\alpha \neq \beta$)

$$P_{\nu_\alpha \rightarrow \nu_\beta}^{\text{SBL}} \simeq \sin^2 2\vartheta_{\alpha\beta} \sin^2 \left(\frac{\Delta m_{41}^2 L}{4E} \right)$$

$$\sin^2 2\vartheta_{\alpha\beta} = 4|U_{\alpha 4}|^2 |U_{\beta 4}|^2$$

Disappearance

$$P_{\nu_\alpha \rightarrow \nu_\alpha}^{\text{SBL}} \simeq 1 - \sin^2 2\vartheta_{\alpha\alpha} \sin^2 \left(\frac{\Delta m_{41}^2 L}{4E} \right)$$

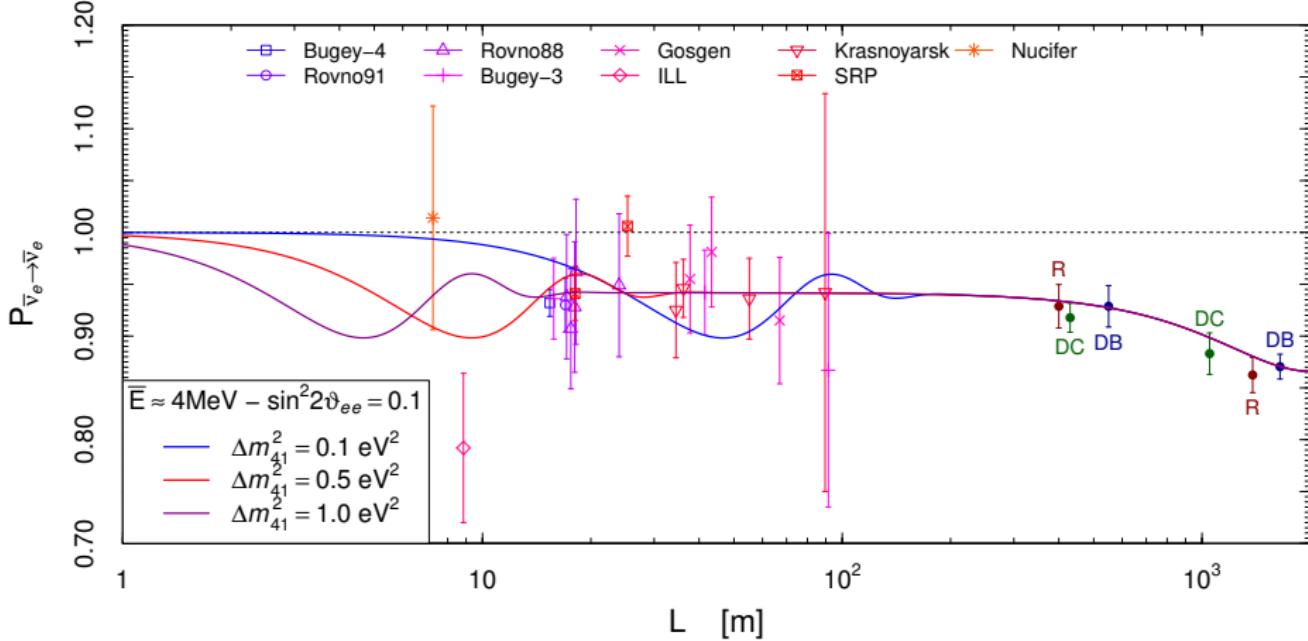
$$\sin^2 2\vartheta_{\alpha\alpha} = 4|U_{\alpha 4}|^2 (1 - |U_{\alpha 4}|^2)$$

$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} & U_{\mu 4} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} & U_{\tau 4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{pmatrix}_{\text{SBL}}$$

- ▶ 6 mixing angles
- ▶ 3 Dirac CP phases
- ▶ 3 Majorana CP phases

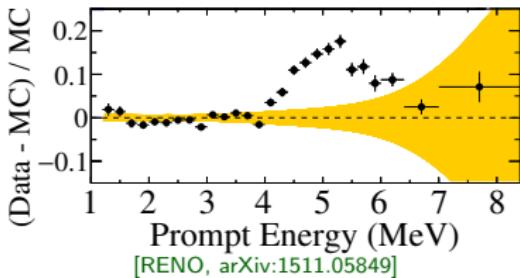
- ▶ $\Delta m_{\text{SBL}}^2 = \Delta m_{41}^2 \simeq \Delta m_{42}^2 \simeq \Delta m_{43}^2$
- ▶ CP violation is not observable in SBL experiments!
- ▶ Observable in LBL accelerator exp. sensitive to Δm_{ATM}^2 [de Gouvea et al, arXiv:1412.1479, arXiv:1507.03986, arXiv:1605.09376; Palazzo et al, arXiv:1412.7524, arXiv:1509.03148; Kayser et al, arXiv:1508.06275, arXiv:1607.02152] and solar exp. sensitive to Δm_{SOL}^2 [Long, Li, Giunti, arXiv:1304.2207]

Short-Baseline Reactor Neutrino Oscillations

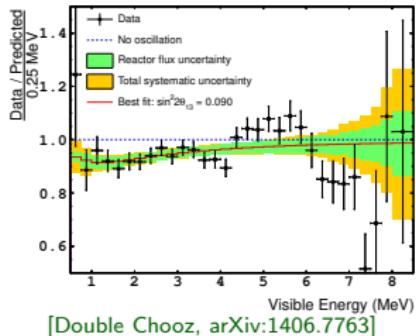


- $\Delta m_{\text{SBL}}^2 \gtrsim 0.5 \text{ eV}^2 \gg \Delta m_{\text{ATM}}^2$
- SBL oscillations are averaged at the Daya Bay, RENO, and Double Chooz near detectors \implies no spectral distortion
- The reactor antineutrino anomaly is model dependent (depends on the theoretical reactor neutrino flux calculation; is it reliable?).

Reactor Antineutrino 5 MeV Bump

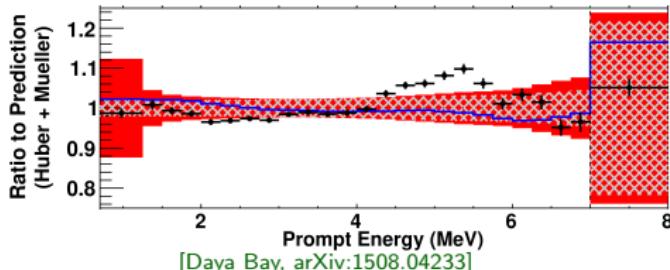


- ▶ Cannot be explained by neutrino oscillations (SBL oscillations are averaged in RENO, DC, DB).



- ▶ If it is due to a theoretical miscalculation of the spectrum, it can have opposite effects on the anomaly:

[see: Berryman, Huber, arXiv:1909.09267]



- ▶ If it is a 4-6 MeV excess it increases the anomaly: new HKSS flux calculation

[Hayen, Kostensalo, Severijns, Suhonen, arXiv:1908.08302]

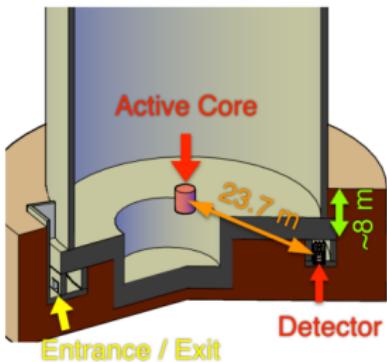
- ▶ If it is a 1-4 MeV suppression it decreases the anomaly: new EF flux calculation

[Estienne, Fallot, et al, arXiv:1904.09358]

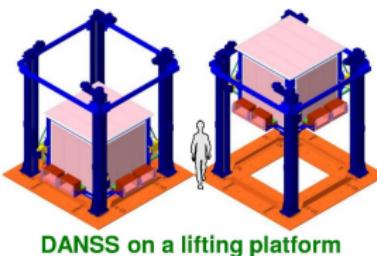
Model Indep. Measurements of Reactor ν Osc.

Ratios of spectra at different distances

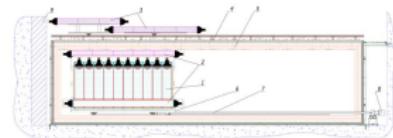
NEOS



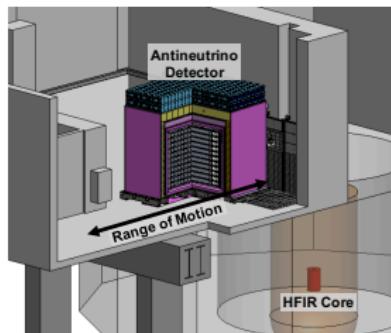
DANSS



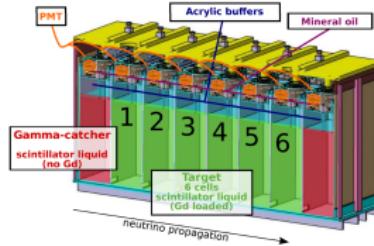
Neutrino-4



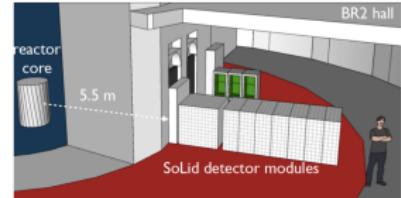
PROSPECT



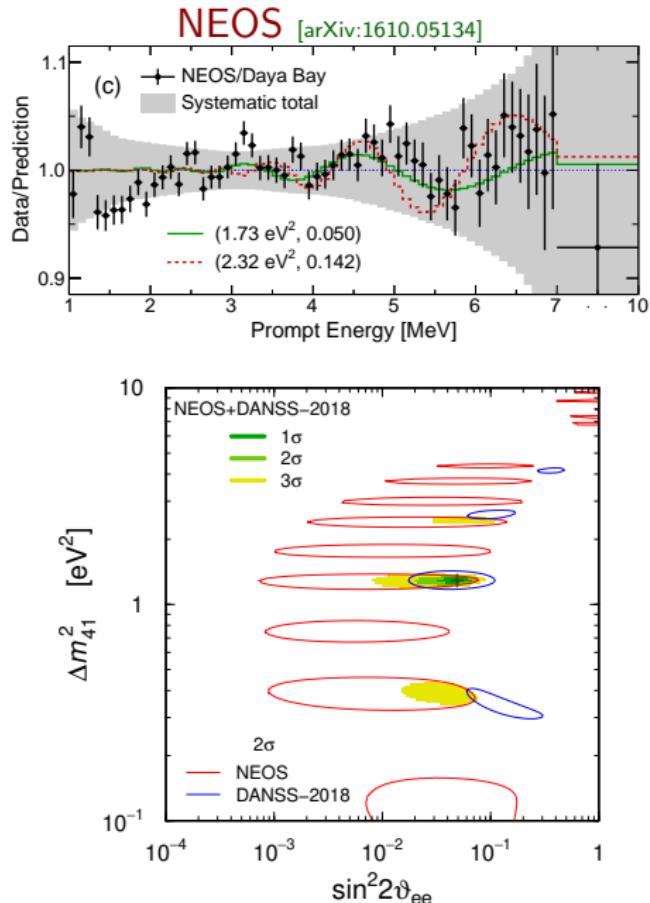
STEREO



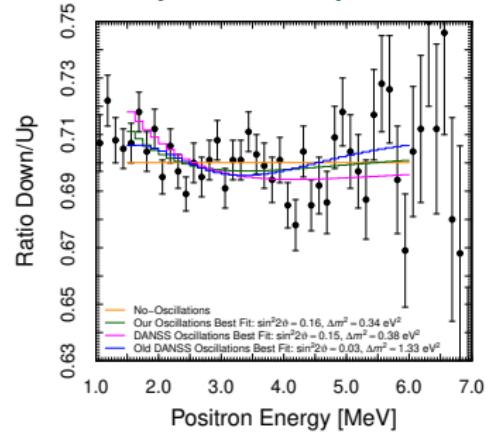
SoLid



2018 Results



DANSS-2018
[arXiv:1804.04046]



2018 model independent indication
in favor of SBL oscillations

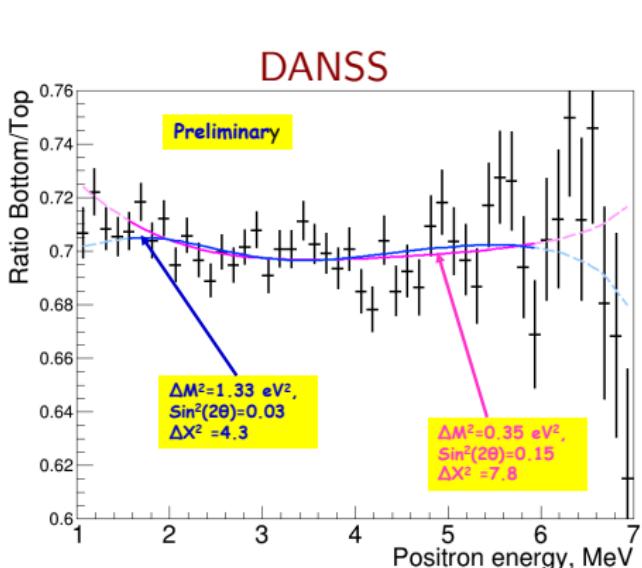
NEOS: $\sim 2.0\sigma$

DANSS-2018: $\sim 2.2\sigma$

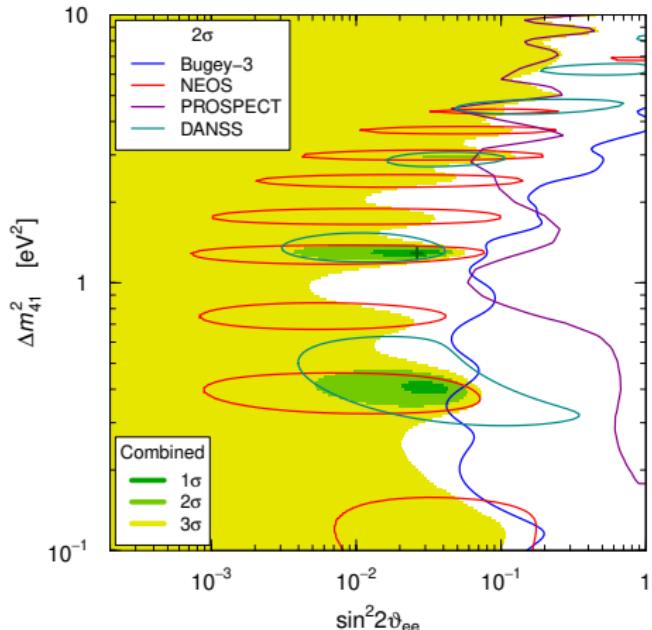
Combined: $\sim 3.5\sigma$

[Gariazzo, Giunti, Laveder, Li, arXiv:1801.06467]
[Dentler, Hernandez-Cabezudo, Kopp, Machado, Maltoni,
Martinez-Soler, Schwetz, arXiv:1803.10661]

2019 Results



[Danilov @ EPS-HEP 2019, arXiv:1911.10140]

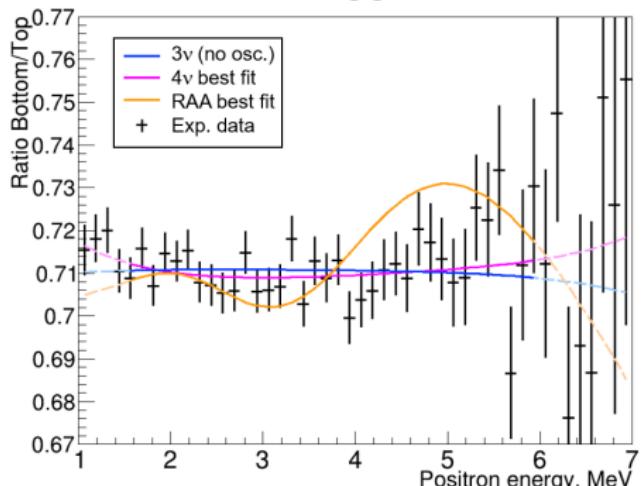


[Giunti, Y.F. Li, Y.Y. Zhang, arXiv:1912.12956]

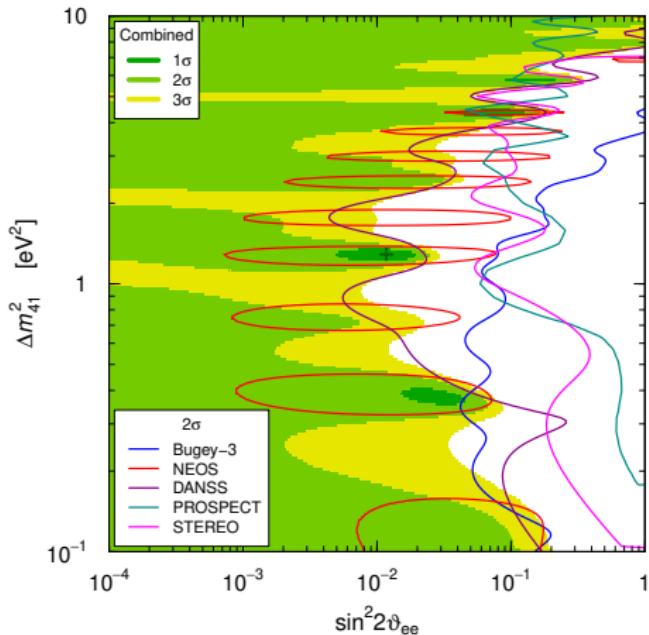
- The agreement between NEOS and DANSS diminished.

2020 Results

DANSS



[Shitov @ Neutrino 2020]

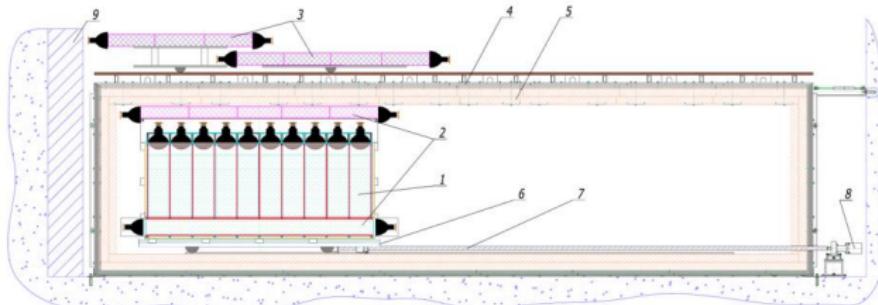


- ▶ No indication of oscillations from DANNS data.
- ▶ In practice these reactor spectral ratios give upper bound on

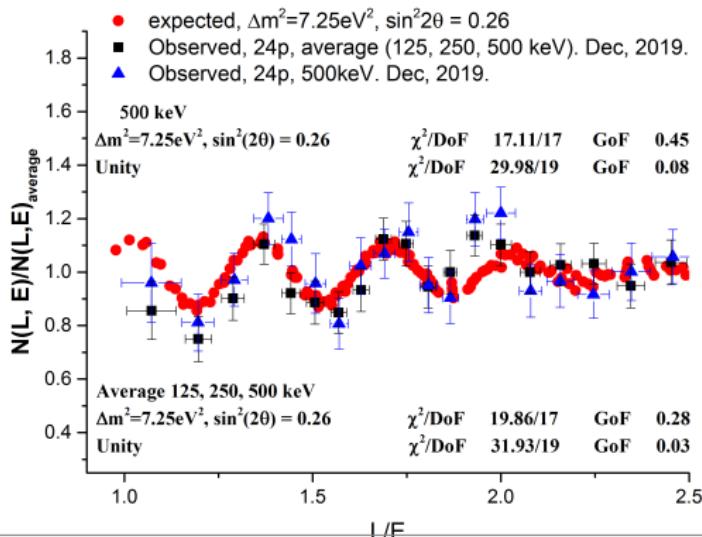
$$\sin^2 2\vartheta_{ee} \simeq 4|U_{e4}|^2$$

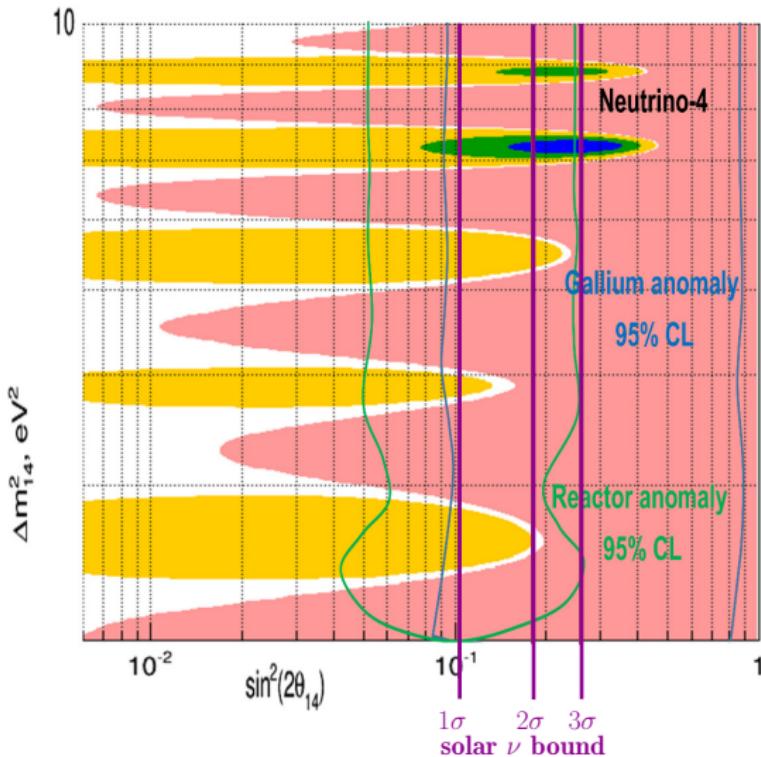
Neutrino-4

[arXiv:1708.00421, arXiv:1809.10561, arXiv:2003.03199, arXiv:2005.05301, arXiv:2006.13639]



- ▶ SM-3 reactor at Dimitrovgrad, Russia
- ▶ Pure ^{235}U $\bar{\nu}_e$ flux
- ▶ $L = 6 - 12 \text{ m}$





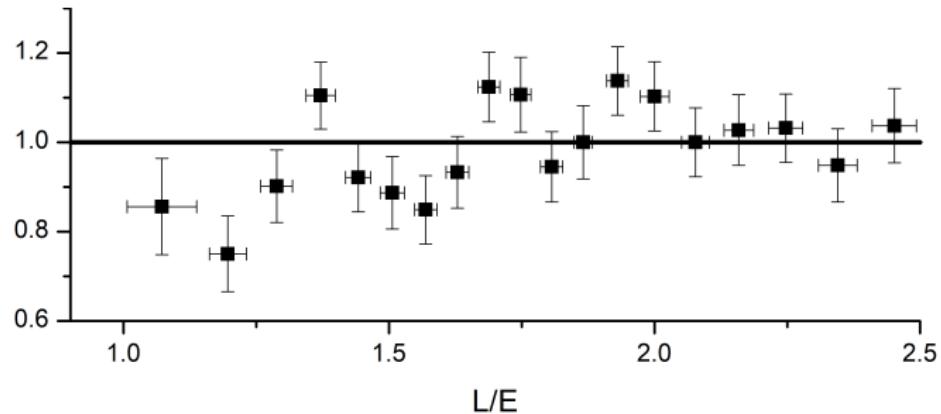
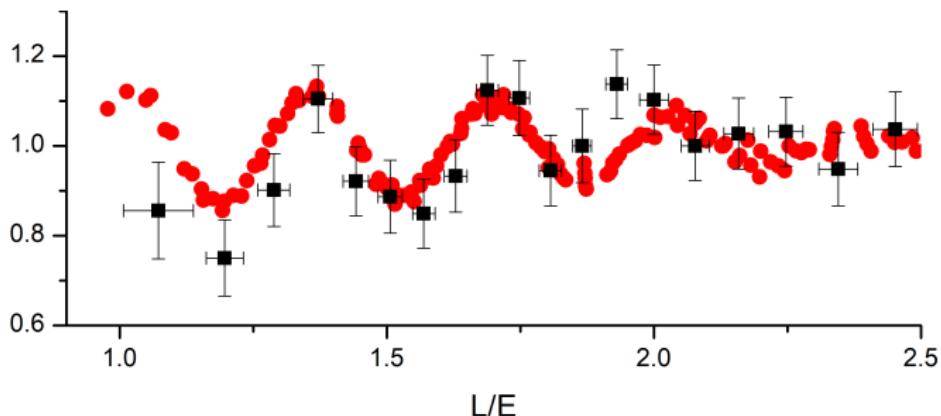
- ▶ Neutrino-4 best fit:
 $\sin^2 2\vartheta_{ee} = 0.26$
 $\Delta m_{41}^2 = 7.25 \text{ eV}^2$
- ▶ Very large mixing!
- ▶ Not a small perturbation of 3ν mixing.
- ▶ Tension with solar neutrino bound.

[Palazzo, arXiv:1105.1705, arXiv:1201.4280]

[Giunti, Laveder, Li, Liu, Long, arXiv:1210.5715]

[Gariazzo, Giunti, Laveder, Li, arXiv:1703.00860]

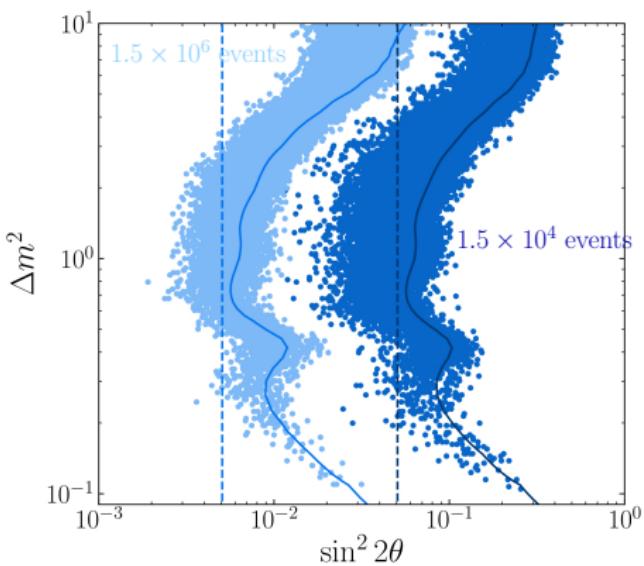
Oscillations or Fluctuations?



Deviations from χ^2 Distribution (Wilks' Theorem)

[Agostini, Neumair, arXiv:1906.11854; Silaeva, Sinev, arXiv:2001.10752; Giunti, arXiv:2004.07577]
[PROSPECT+STEREO, arXiv:2006.13147; Coloma, Huber, Schwetz, arXiv:2008.06083]

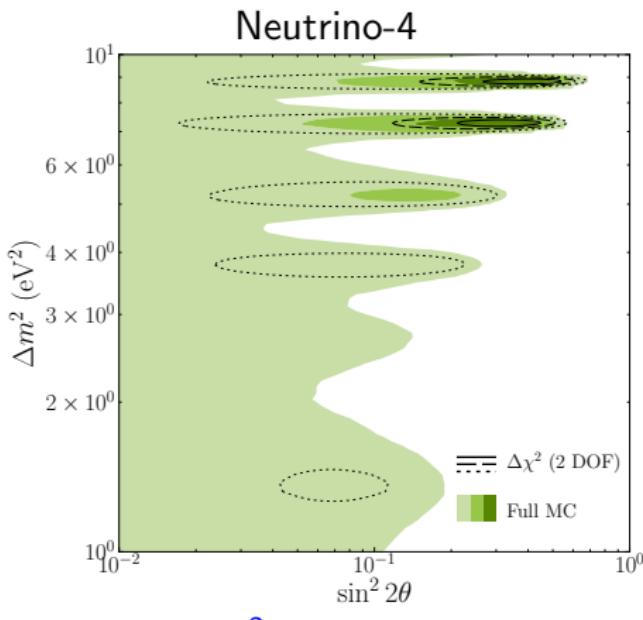
Even in the **absence of real oscillations**, binned data can often be **fitted better by oscillations** that reproduce the statistical fluctuations of the bins.



[Coloma, Huber, Schwetz, arXiv:2008.06083]

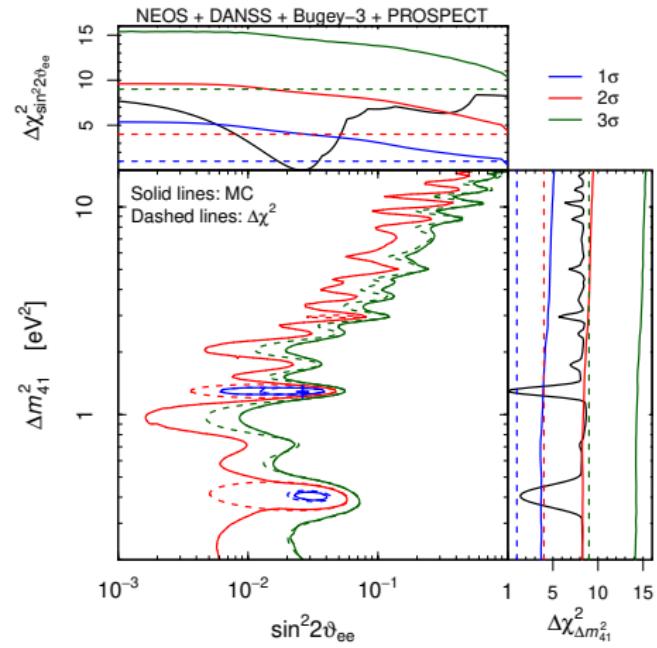
- ▶ Numerical simulations of a reactor neutrino experiment experiment: $\bar{\nu}_e$ disappearance.
- ▶ Location of the **best-fit points** of 20,000 pseudo-experiments **simulated under the no-oscillation hypothesis**.
- ▶ Vertical lines: expected analytical $\langle \sin^2 2\theta \rangle$ from a toy model.
- ▶ Solid curves: sensitivity at 95% CL assuming that Wilks' theorem holds.

MC evaluation of test statistic distribution



$3.2\sigma (\Delta \chi^2) \rightarrow 2.6\sigma (\text{MC})$

[Coloma, Huber, Schwetz, arXiv:2008.06083]



$2.4\sigma (\Delta \chi^2) \rightarrow 1.8\sigma (\text{MC})$

[Giunti, arXiv:2004.07577]

- MC calculations are unfortunately difficult and require a lot of computer time.
- They must be completely redone for each combination of experiments.

- ▶ The MC evaluation of test statistic distribution **decreases** the statistical significance of the indications in favor of oscillations.
- ▶ Nevertheless, the indications must be checked by other experiments.
- ▶ We do not want to miss a chance to discover sterile neutrinos and physics beyond the Standard Model.
- ▶ This is valid for Neutrino-4 as well as for any other indication.
- ▶ It would be very interesting if the Neutrino-4 results are confirmed by other experiments, opening an unexpected new scenario.

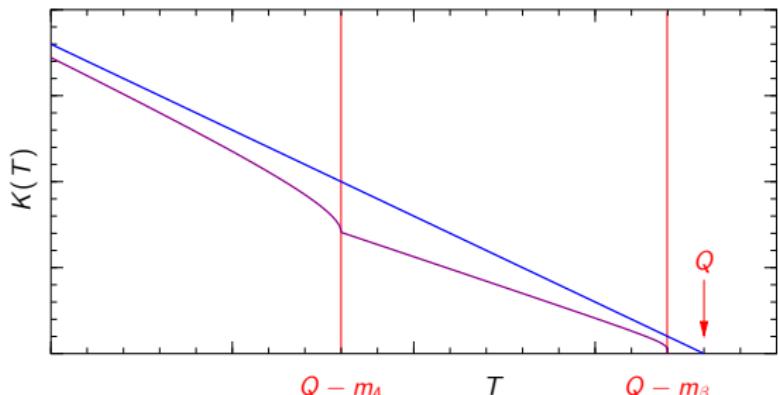
Robust kinematical probe of $\nu_e - \nu_s$ mixing

Tritium Beta-Decay: ${}^3\text{H} \rightarrow {}^3\text{He} + e^- + \bar{\nu}_e$

$$\frac{d\Gamma}{dT} = \frac{(\cos\vartheta_C G_F)^2}{2\pi^3} |\mathcal{M}|^2 F(E) p E K^2(T)$$

$$\frac{K^2(T)}{Q - T} = \sum_k |U_{ek}|^2 \sqrt{(Q - T)^2 - m_k^2} \theta(Q - T - m_k)$$

$$m_4 \gg m_{1,2,3} \Rightarrow \simeq (1 - |U_{e4}|^2) \sqrt{(Q - T)^2 - m_\beta^2} \theta(Q - T - m_\beta) \\ + |U_{e4}|^2 \sqrt{(Q - T)^2 - m_4^2} \theta(Q - T - m_4)$$



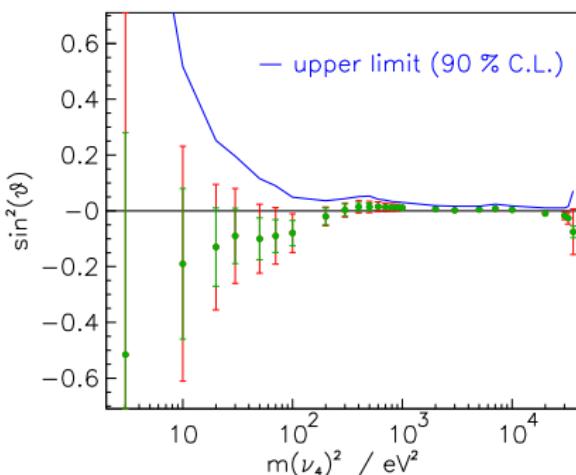
$$Q = M_{^3\text{H}} - M_{^3\text{He}} - m_e \\ = 18.58 \text{ keV}$$

$$m_\beta^2 = \sum_{k=1}^3 |U_{ek}|^2 m_k^2$$

Mainz and Troitsk Limit on $\Delta m_{41}^2 \simeq m_4^2$

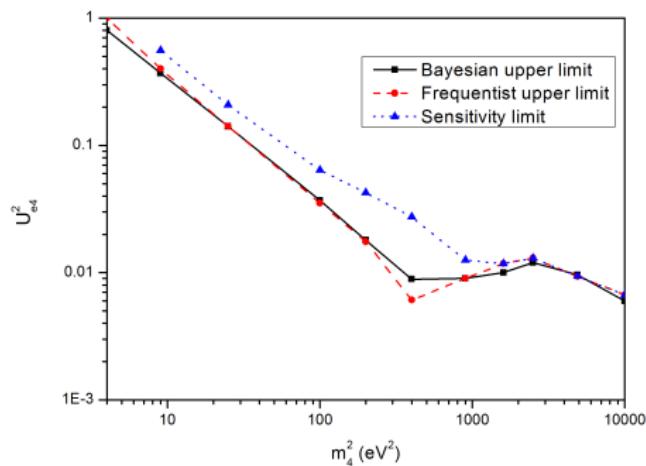
$$m_4 \gg m_{1,2,3} \implies \Delta m_{41}^2 \equiv m_4^2 - m_1^2 \simeq m_4^2$$

Mainz



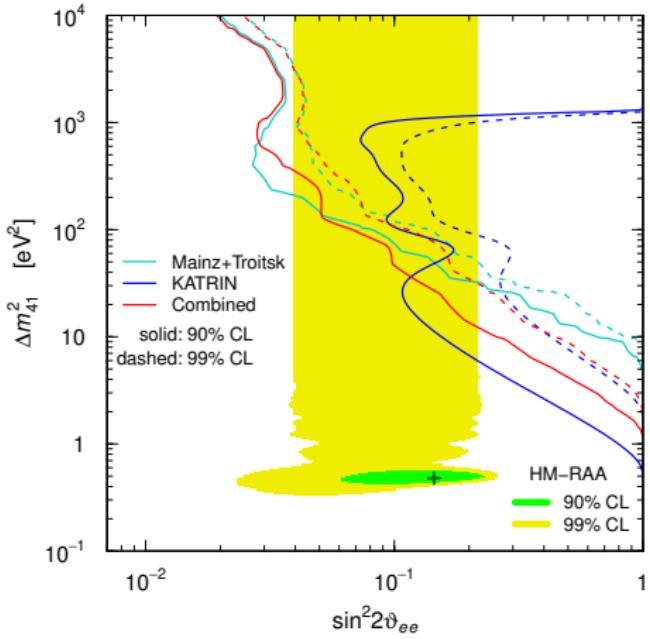
[Kraus, Singer, Valerius, Weinheimer, arXiv:1210.4194]

Troitsk



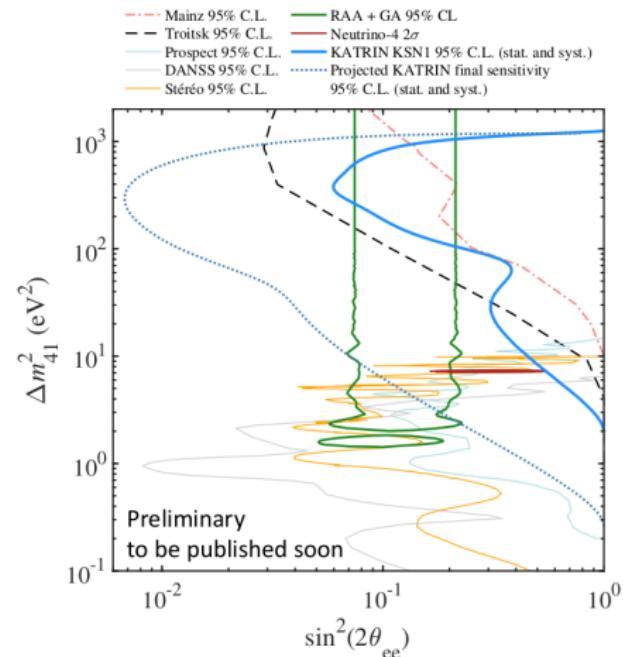
[Belesev et al, arXiv:1307.5687]

Bound from first KATRIN data



[Giunti, Y.F. Li, Y.Y. Zhang, arXiv:1912.12956]

$$\Delta m_{41}^2 \simeq m_4^2$$



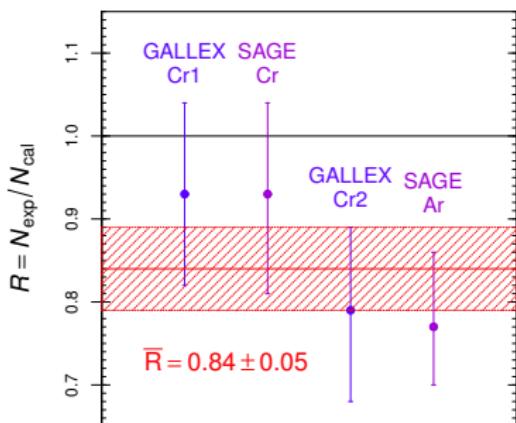
[KATRIN @ Neutrino 2020]

Gallium Anomaly

Gallium Radioactive Source Experiments: GALLEX and SAGE

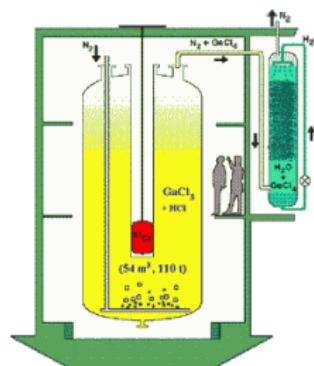


Test of Solar ν_e Detection:



$$\langle L \rangle_{\text{GALLEX}} = 1.9 \text{ m} \quad \langle L \rangle_{\text{SAGE}} = 0.6 \text{ m}$$

$$\Delta m_{\text{SBL}}^2 \gtrsim 1 \text{ eV}^2 \gg \Delta m_{\text{ATM}}^2$$

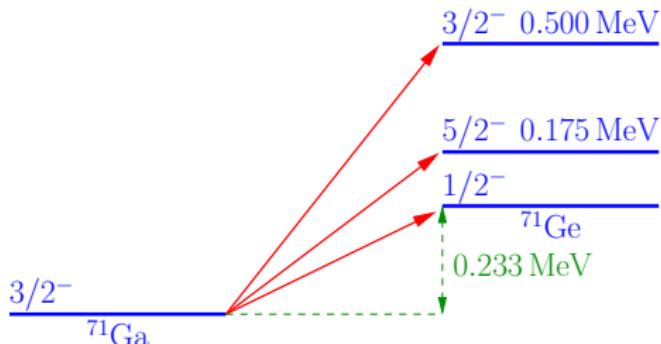


$\approx 2.9\sigma$ deficit

[SAGE, nucl-ex/0512041, arXiv:0901.2200;
Laveder et al, Nucl.Phys.Proc.Suppl. 168 (2007) 344,
hep-ph/0610352, arXiv:0711.4222,
arXiv:1006.3244]

- Deficit could be due to an overestimate of
 $\sigma(\nu_e + {}^{71}\text{Ga} \rightarrow {}^{71}\text{Ge} + e^-)$

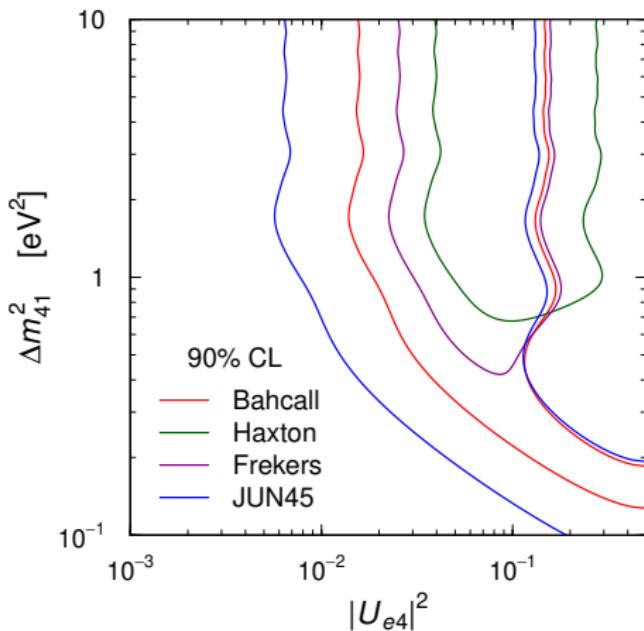
- First calculation: Bahcall, PRC 56 (1997) 3391



- $\sigma_{\text{G.S.}}$ from $T_{1/2}({}^{71}\text{Ge}) = 11.43 \pm 0.03$ days [Hampel, Remsberg, PRC 31 (1985) 666]

$$\sigma_{\text{G.S.}}({}^{51}\text{Cr}) = 55.3 \times 10^{-46} \text{ cm}^2 (1 \pm 0.004)_{3\sigma}$$

- $\sigma({}^{51}\text{Cr}) = \sigma_{\text{G.S.}}({}^{51}\text{Cr}) \left(1 + 0.669 \frac{\text{BGT}_{175}}{\text{BGT}_{\text{G.S.}}} + 0.220 \frac{\text{BGT}_{500}}{\text{BGT}_{\text{G.S.}}} \right)$
- The contribution of excited states is only $\sim 5\%$, but it is crucial for the size of the Gallium anomaly!



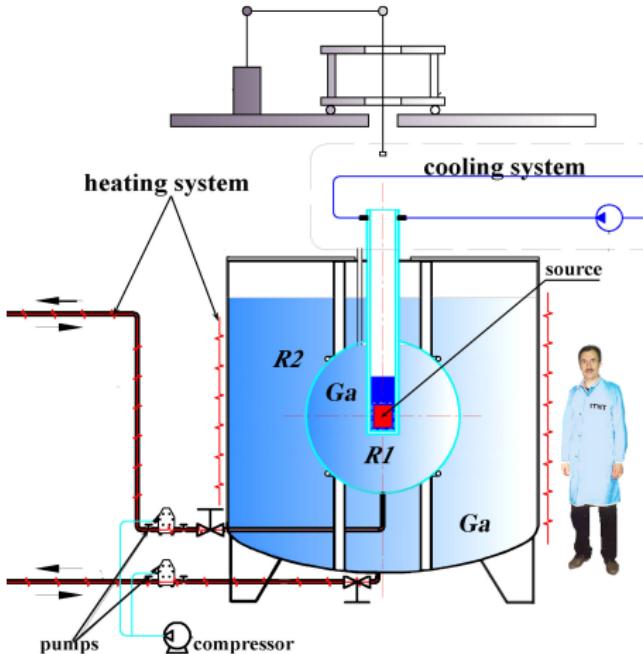
Cross sections in units of 10^{-45} cm^2 :

	$\sigma(^{51}\text{Cr})$	$\sigma(^{37}\text{Ar})$
Bahcall	5.81 ± 0.16	7.00 ± 0.21
Haxton	6.39 ± 0.65	7.72 ± 0.81
Frekers	5.92 ± 0.11	7.15 ± 0.14
JUN45	5.67 ± 0.06	6.80 ± 0.08

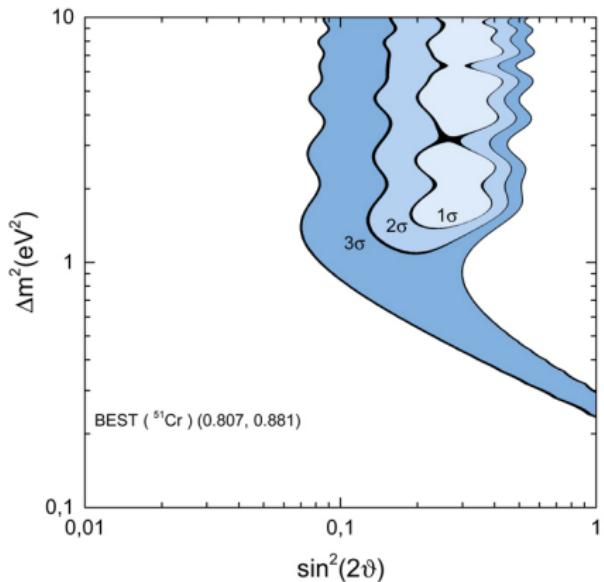
[Kostensalo, Suhonen, Giunti, Srivastava, arXiv:1906.10980]

[arXiv:1006.2103, arXiv:1602.03826, arXiv:1710.06326, arXiv:1807.02977, arXiv:1905.07437]

Direct test of the Gallium anomaly with ^{51}Cr source.

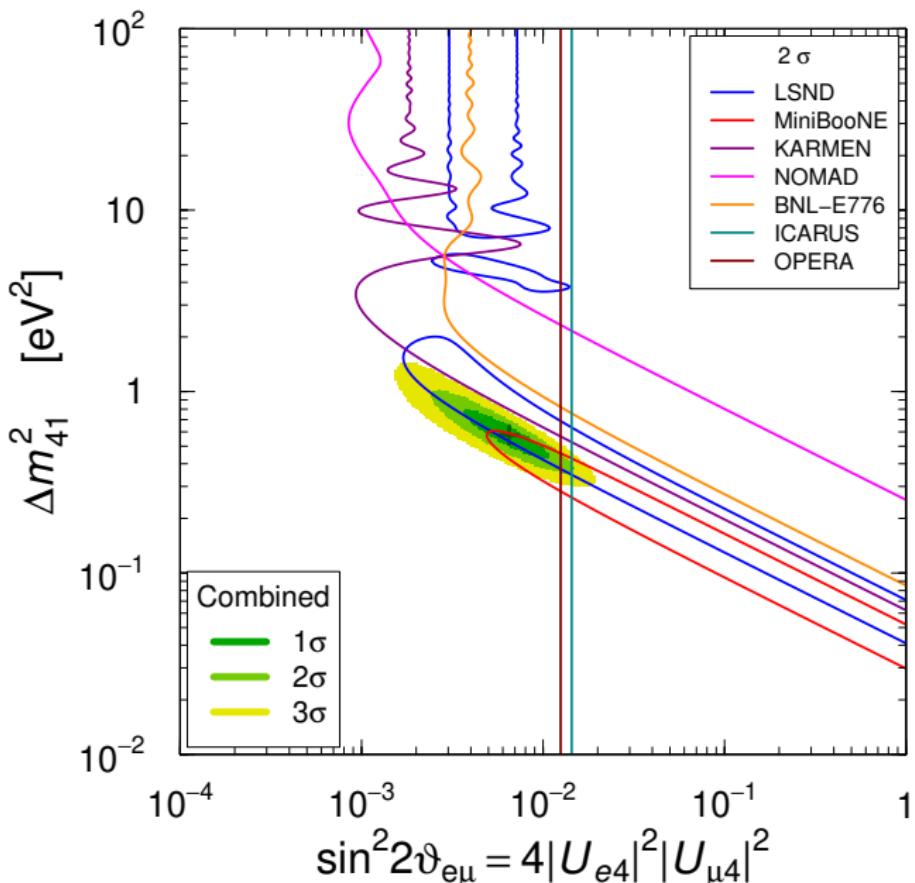


$$R_1 = 0.66 \text{ m}, \quad R_2 = 1.096 \text{ m}$$

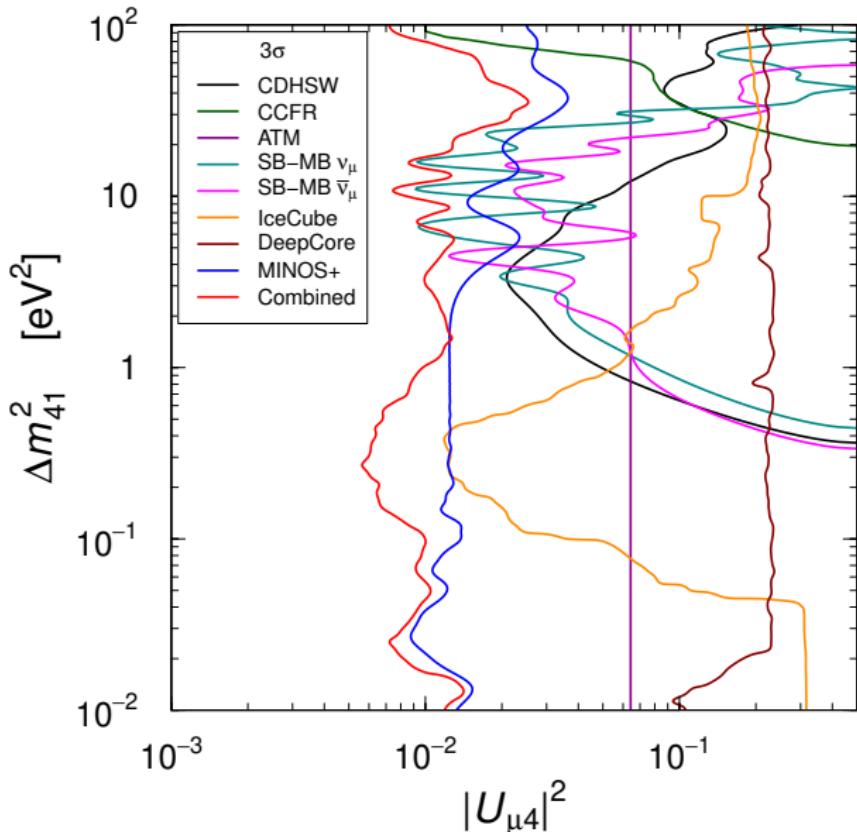


Allowed regions of oscillation parameters if the result of the BEST experiment corresponds to the best fit point for combining the SAGE + GALLEX. The numbers in parentheses indicate the most probable ratios R of observed-to-expected without sterile neutrinos germanium atoms in the two vessels.

$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ and $\nu_\mu \rightarrow \nu_e$ Appearance



ν_μ and $\bar{\nu}_\mu$ Disappearance



[Gariazzo, Giunti, Li, Ternes, Zhang, in preparation]

3+1: Appearance vs Disappearance

► SBL Oscillation parameters: Δm_{41}^2 $|U_{e4}|^2$ $|U_{\mu 4}|^2$ ($|U_{\tau 4}|^2$)

► Amplitude of ν_e disappearance:

$$\sin^2 2\vartheta_{ee} = 4|U_{e4}|^2 (1 - |U_{e4}|^2) \simeq 4|U_{e4}|^2$$

► Amplitude of ν_μ disappearance:

$$\sin^2 2\vartheta_{\mu\mu} = 4|U_{\mu 4}|^2 (1 - |U_{\mu 4}|^2) \simeq 4|U_{\mu 4}|^2$$

► Amplitude of $\nu_\mu \rightarrow \nu_e$ appearance:

$$\sin^2 2\vartheta_{e\mu} = 4|U_{e4}|^2 |U_{\mu 4}|^2 \simeq \frac{1}{4} \sin^2 2\vartheta_{ee} \sin^2 2\vartheta_{\mu\mu}$$

quadratically suppressed for small $|U_{e4}|^2$ and $|U_{\mu 4}|^2$



Appearance-Disappearance Tension

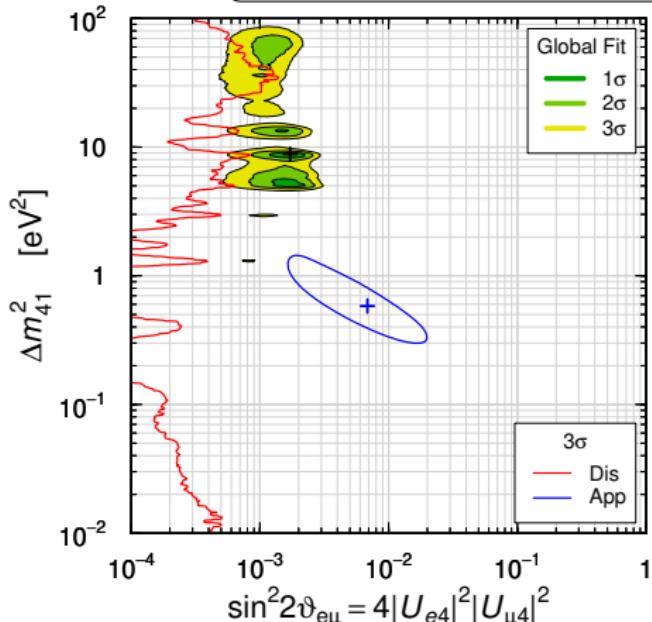
[Okada, Yasuda, hep-ph/9606411; Bilenky, CG, Grimus, hep-ph/9607372]

Global Appearance-Disappearance Tension

$$\nu_e \text{ DIS}$$
$$\sin^2 2\vartheta_{ee} \simeq 4|U_{e4}|^2$$

$$\nu_\mu \text{ DIS}$$
$$\sin^2 2\vartheta_{\mu\mu} \simeq 4|U_{\mu 4}|^2$$

$$\nu_\mu \rightarrow \nu_e \text{ APP}$$
$$\sin^2 2\vartheta_{e\mu} = 4|U_{e4}|^2|U_{\mu 4}|^2 \simeq \frac{1}{4} \sin^2 2\vartheta_{ee} \sin^2 2\vartheta_{\mu\mu}$$



► $\nu_\mu \rightarrow \nu_e$ is quadratically suppressed!

► Global Fit:

$$\chi^2/\text{NDF} = 843.6/794$$

$$\text{GoF} = 11\%$$

$$\chi^2_{\text{PG}}/\text{NDF}_{\text{PG}} = 46.7/2$$

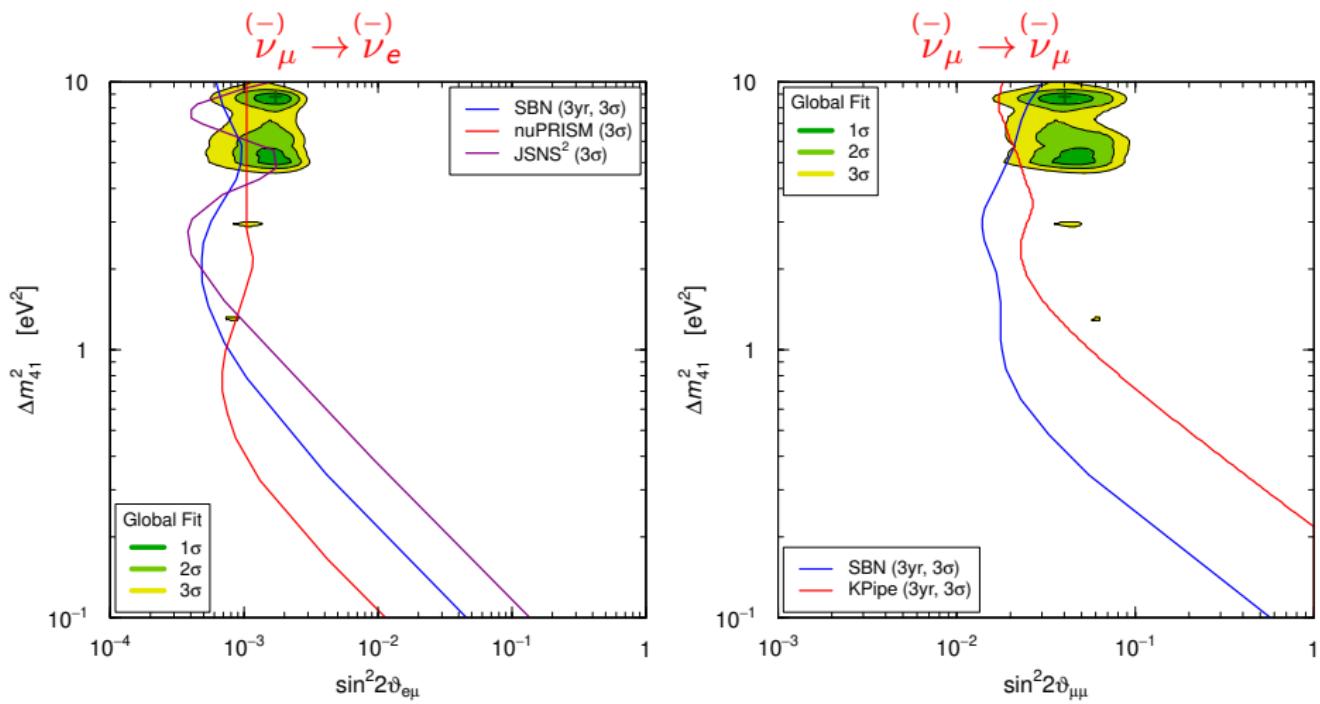
$$\text{GoF}_{\text{PG}} = 7 \times 10^{-11} \leftarrow \text{:(}$$

► Similar tension in

$$3+2, \quad 3+3, \quad \dots, \quad 3+N_s$$

[Giunti, Zavanin, arXiv:1508.03172]

New Dedicated Experiments



Conclusions

- ▶ Neutrinos can be powerful messengers of new physics beyond the SM.
- ▶ The existence of light sterile neutrinos beyond the SM is indicated by the Reactor, Gallium and LSND anomalies.
- ▶ Experimental results are confusing, pointing in different directions.
- ▶ Therefore, there is no definitive conclusion yet.
- ▶ The search must be continued with enthusiasm, because a positive outcome would yield a huge reward.
- ▶ Oscillation experiments suffer of misleading oscillatory fit of statistical fluctuations of the data. Difficult MC evaluation of test statistic distribution is needed to obtain reliable confidence levels.
- ▶ Robust kinematical probe of $\nu_e - \nu_s$ mixing with β decay (KATRIN, Project 8, ...) and electron-capture (ECHO, HOLMES, ...) experiments.