

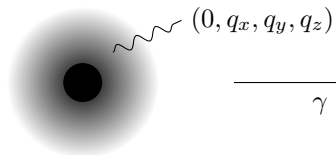
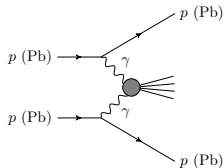
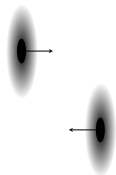
# Equivalent photons approximation: survival factor

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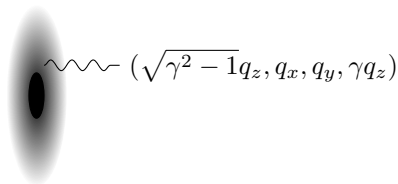
Physical Institute of the Russian Academy of Sciences  
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Moscow, Russia

# Ultrapерipheral collisions



$\xrightarrow{\gamma \gg 1}$



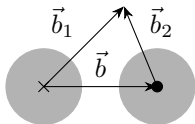
Photon virtuality:  $-q^2 = q_x^2 + q_y^2 + q_z^2 \ll (\gamma q_z)^2 \equiv \omega^2$

## UPC spectrum and cross section

With the sizes of the colliding particles neglected:

$$n(\omega) = \frac{2Z^2\alpha}{\pi\omega} \int_0^\infty \left[ \frac{F(\sqrt{q_\perp^2 + (\omega/\gamma)^2})}{q_\perp^2 + (\omega/\gamma)^2} \right]^2 q_\perp^3 dq_\perp,$$

$$\sigma(AB \rightarrow ABX) = \int_0^\infty d\omega_1 \int_0^\infty d\omega_2 \sigma(\gamma\gamma \rightarrow X) n_A(\omega_1) n_B(\omega_2).$$



With the sizes of the colliding particles taken into account:

$$n(\omega) = \int n(b, \omega) d^2b, \quad n(b, \omega) = \frac{Z^2\alpha}{\pi^2\omega} \left[ \int_0^\infty \frac{F(\sqrt{q_\perp^2 + (\omega/\gamma)^2})}{q_\perp^2 + (\omega/\gamma)^2} J_1(bq_\perp) q_\perp^2 dq_\perp \right]^2,$$

$$\sigma(AB \rightarrow ABX)$$

$$= \int_0^\infty d\omega_1 \int_0^\infty d\omega_2 \sigma(\gamma\gamma \rightarrow X) \int d^2b_1 \int d^2b_2 n_A(b_1, \omega_1) n_B(b_2, \omega_2) P_{AB}(|\vec{b}_1 - \vec{b}_2|).$$

$F(q)$  is the electromagnetic form factor,  $P_{AB}(b)$  is the probability for the colliding particles to survive after the collision with the impact parameter  $b$ .

## Photon-photon luminosity and survival factor

Let  $x = \omega_1/\omega_2$ ,  $s = 4\omega_1\omega_2$ . Then

$$\sigma(AB \rightarrow ABX) = \int_0^\infty \sigma(\gamma\gamma \rightarrow X) \cdot \frac{dL_{AB}}{ds} ds,$$

where

$$\frac{dL_{AB}}{ds} = \int_0^\infty \frac{dx}{8x} \int d^2b_1 \int d^2b_2 n_A \left( b_1, \sqrt{\frac{sx}{4}} \right) n_B \left( b_2, \sqrt{\frac{s}{4x}} \right) P_{AB}(|\vec{b}_1 - \vec{b}_2|)$$

is the  $\gamma\gamma$  luminosity in the collision of the particles  $A$  and  $B$ . When  $P_{AB}(b) \equiv 1$ ,

$$\left. \frac{dL_{AB}}{ds} \right|_{P=1} = \int_0^\infty \frac{dx}{8x} n_A \left( \sqrt{\frac{sx}{4}} \right) n_B \left( \sqrt{\frac{s}{4x}} \right).$$

The survival factor:

$$S_{AB} = \frac{dL_{AB}/ds}{dL_{AB}/ds|_{P=1}}.$$

$\sqrt{s}$  is the invariant mass of the system produced.

Proton form factor:

$$F_p(q) = \frac{1}{(1 + (q/\Lambda)^2)^2} \left[ 1 + \frac{(\mu_p - 1)\tau}{1 + \tau} \right] \approx \frac{1}{(1 + (q/\Lambda)^2)^2}$$

$\Lambda = 0.84$  GeV,  $\mu_p = 2.73$  is the proton magnetic moment,  $\tau = q^2/4m_p^2$ . Since  $q^2 \lesssim \Lambda_{\text{QCD}}^2 \ll 4m_p^2$ ,  $\tau \ll 1$ .

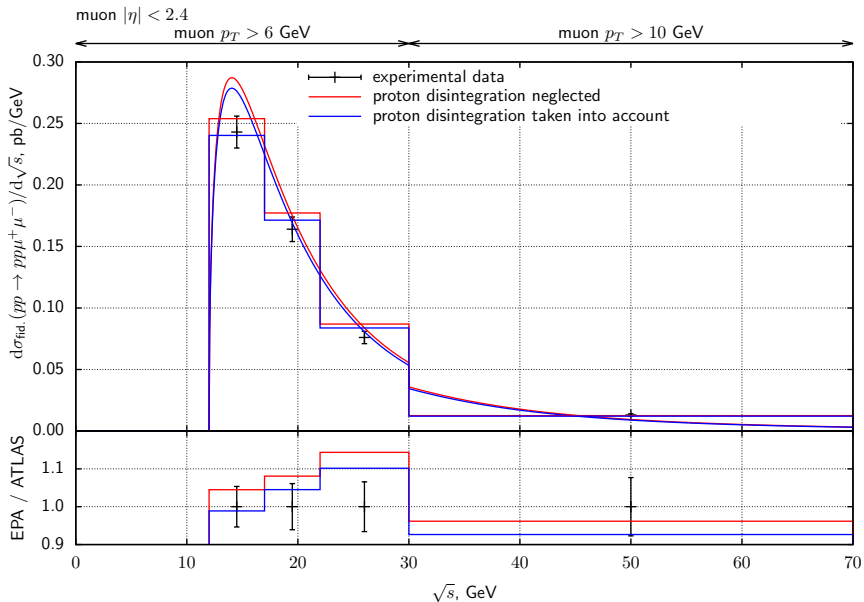
Proton EPA spectrum (dipole approximation):

$$n_2(b, \omega) = \frac{\alpha}{\pi^2 \omega} \left[ \frac{\omega}{\gamma} K_1 \left( \frac{b\omega}{\gamma} \right) - \sqrt{\Lambda^2 + \left( \frac{\omega}{\gamma} \right)^2} K_1 \left( b \sqrt{\Lambda^2 + \left( \frac{\omega}{\gamma} \right)^2} \right) - \frac{b\Lambda^2}{2} K_0 \left( b \sqrt{\Lambda^2 + \left( \frac{\omega}{\gamma} \right)^2} \right) \right]^2.$$

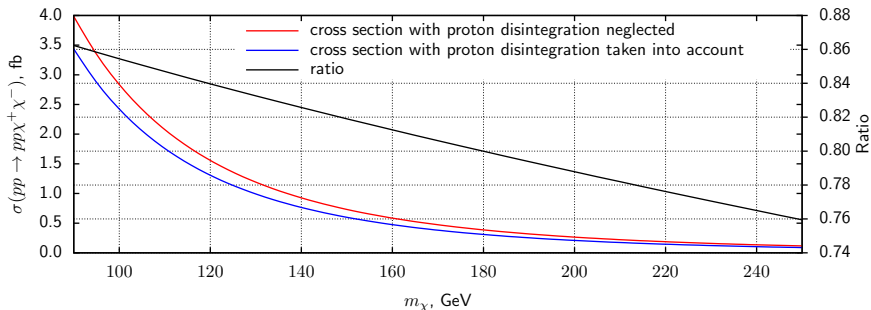
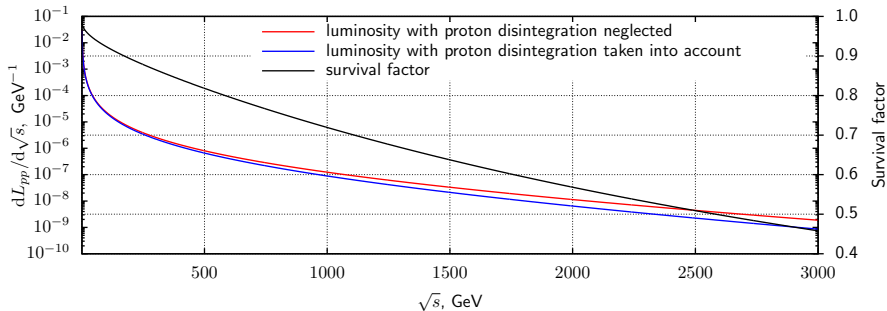
Proton survival probability [hep-ph/0608271, 1408.5778]:

$$P_{pp}(b) = \left( 1 - e^{-\frac{b^2}{2B}} \right)^2, \quad B = 19.7 \text{ GeV}^2.$$

# $pp \rightarrow pp \mu \mu$ [1708.04053]



# $pp$ survival factor



# Pb Pb UPC

Heavy nucleus form factor (Fourier-Bessel decomposition):

$$F_{\text{Pb}}(q) = \frac{\sin qR \sum_{k=1}^N \frac{(-1)^k a_k}{(k\pi)^2 - (qR)^2}}{qR \sum_{k=1}^N \frac{(-1)^k a_k}{(k\pi)^2}}.$$

Monopole approximation:

$$F_1(q) = \frac{1}{1 + (q/\Lambda)^2}, \quad \Lambda = 50 \text{ MeV [1995] or } 80 \text{ MeV [1987]}.$$

Monopole EPA spectrum:

$$n_1(b, \omega) = \frac{Z^2 \alpha}{\pi^2 \omega} \left[ \frac{\omega}{\gamma} K_1 \left( \frac{b\omega}{\gamma} \right) - \sqrt{\Lambda^2 + \left( \frac{\omega}{\gamma} \right)^2} K_1 \left( b \sqrt{\Lambda^2 + \left( \frac{\omega}{\gamma} \right)^2} \right) \right]^2.$$

Pb survival probability [nucl-ex/0302016, 1607.03838]

$$P_{\text{Pb}}(b) = \exp \left[ -\sigma_{NN} \iint T(b_1) T(b_2) \delta^{(2)}(\vec{b} - \vec{b}_1 + \vec{b}_2) d^2 b_1 d^2 b_2 \right],$$

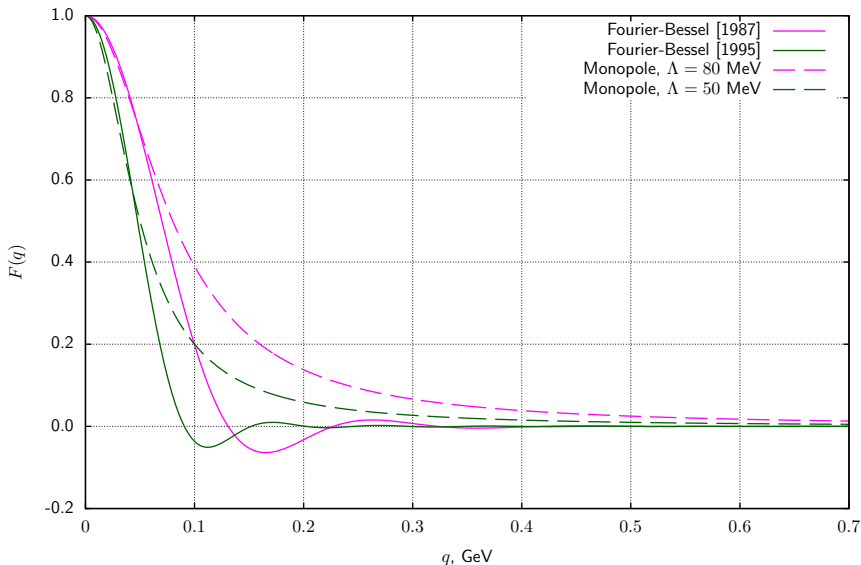
$$T(b) = \int_{-\infty}^{\infty} \rho_N(\sqrt{b^2 + z^2}) dz$$



# $^{208}\text{Pb}$ form factor

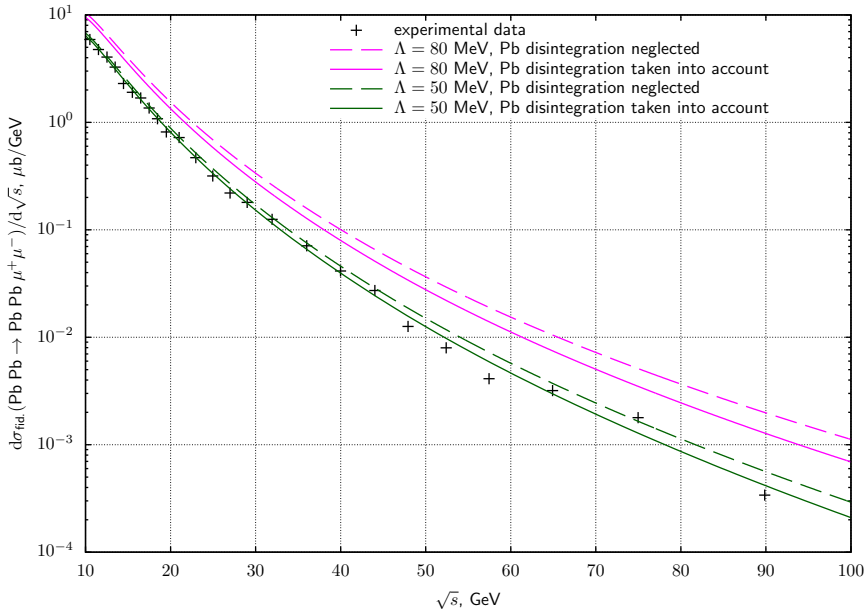
[1987]: At.Data and Nucl.Data Tabl. 36, 495 (1987)

[1995]: At.Data and Nucl.Data Tabl. 60, 177 (1995)

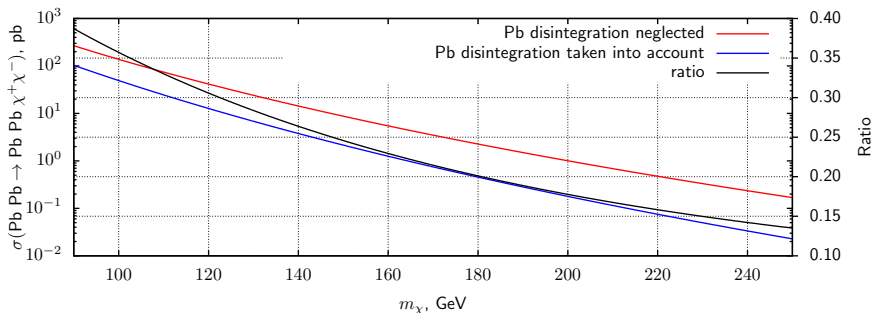
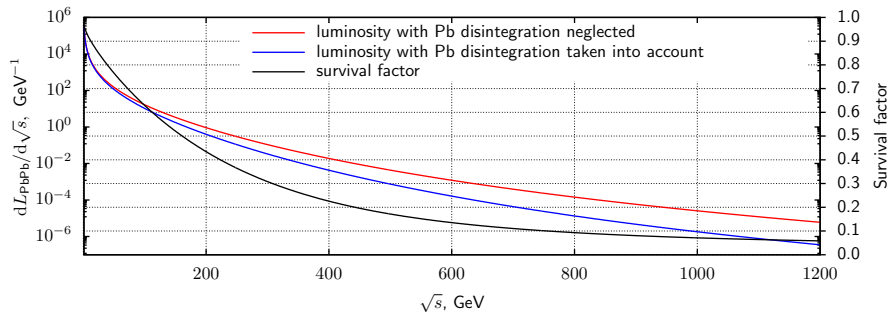


# PbPb $\rightarrow$ PbPb $\mu\mu$ [ATLAS-CONF-2016-025]

muon  $|\eta| < 2.4$ ,  $p_T > 4$  GeV



# $^{208}\text{Pb}$ survival factor

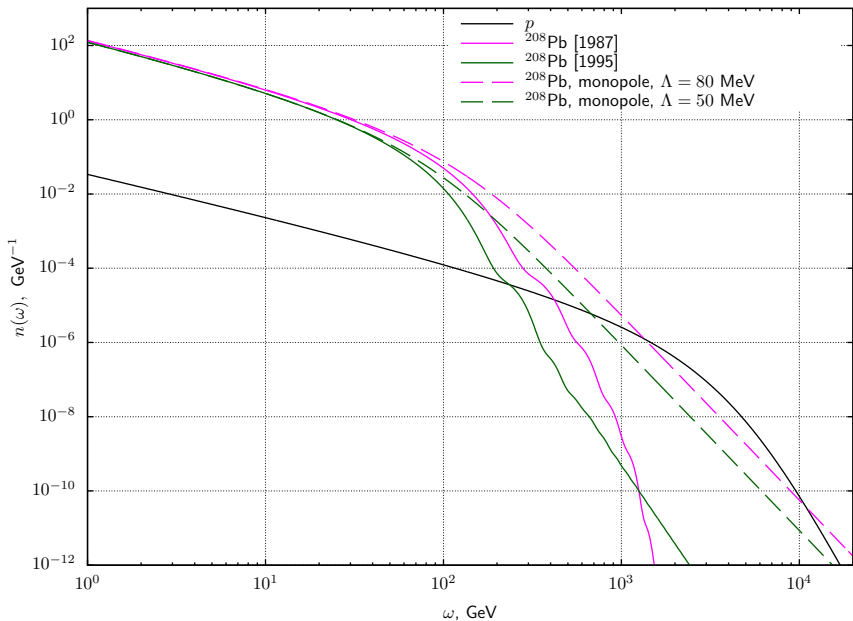


# Conclusion

- ▶ Protons:
  - ▶ The spectrum can be calculated analytically.
  - ▶ EPA approximation describes well the experimental data.
  - ▶ Taking into account the survival factor reduces by  $\sim 5\%$  the cross section for the production of light particles (muons) and by  $\sim 20\%$  the cross section for the production of particles with the mass 200 GeV.
- ▶  $^{208}\text{Pb}$ :
  - ▶ So far the spectrum can only be calculated for the monopole approximation.
  - ▶ EPA approximation depends strongly on the nucleus form factor.
  - ▶ Calculation with the older form factor overshoots the experiment by 50%; the newer form factor fits the experiment perfectly but has some inconsistencies (see backup). This is the subject of future investigations.
  - ▶ Taking into account the survival factor reduces by  $\sim 6\%$  the cross section for the production of light particles (muons) and over 80% the cross section for the production of particles with mass 200 GeV.

Thank you for your attention

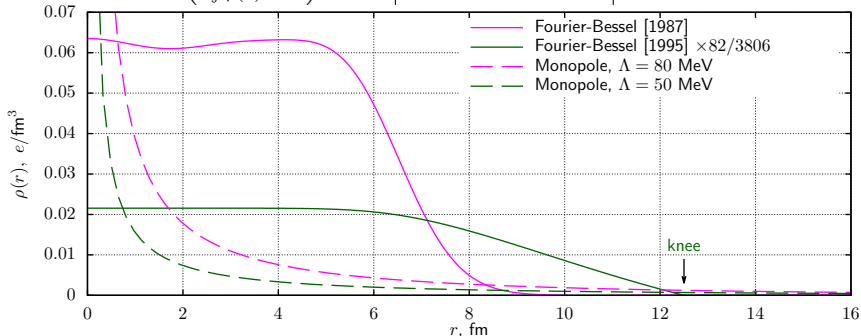
# EPA spectra



# <sup>208</sup>Pb charge density

$$\rho(r) = \sum_{k=1}^N a_k j_0(k\pi r/R) \theta(R-r), \quad F(q) = \frac{\int \rho(r) e^{i\vec{q}\vec{r}} d^3r}{\int \rho(r) d^3r}$$

	[1987]		[1995]	
	Quoted	Actual	Quoted	Actual
$Z = \int \rho(r) d^3r$	82	82.03	82	3806
$\langle r^2 \rangle^{1/2} = \left( \frac{\int \rho(r) r^2 d^3r}{\int \rho(r) d^3r} \right)^{1/2}, \text{ fm}$	5.499(1)	5.501	5.4785	8.0850



[1987]: At.Data and Nucl.Data Tabl. 36, 495 (1987)

[1995]: At.Data and Nucl.Data Tabl. 60, 177 (1995)