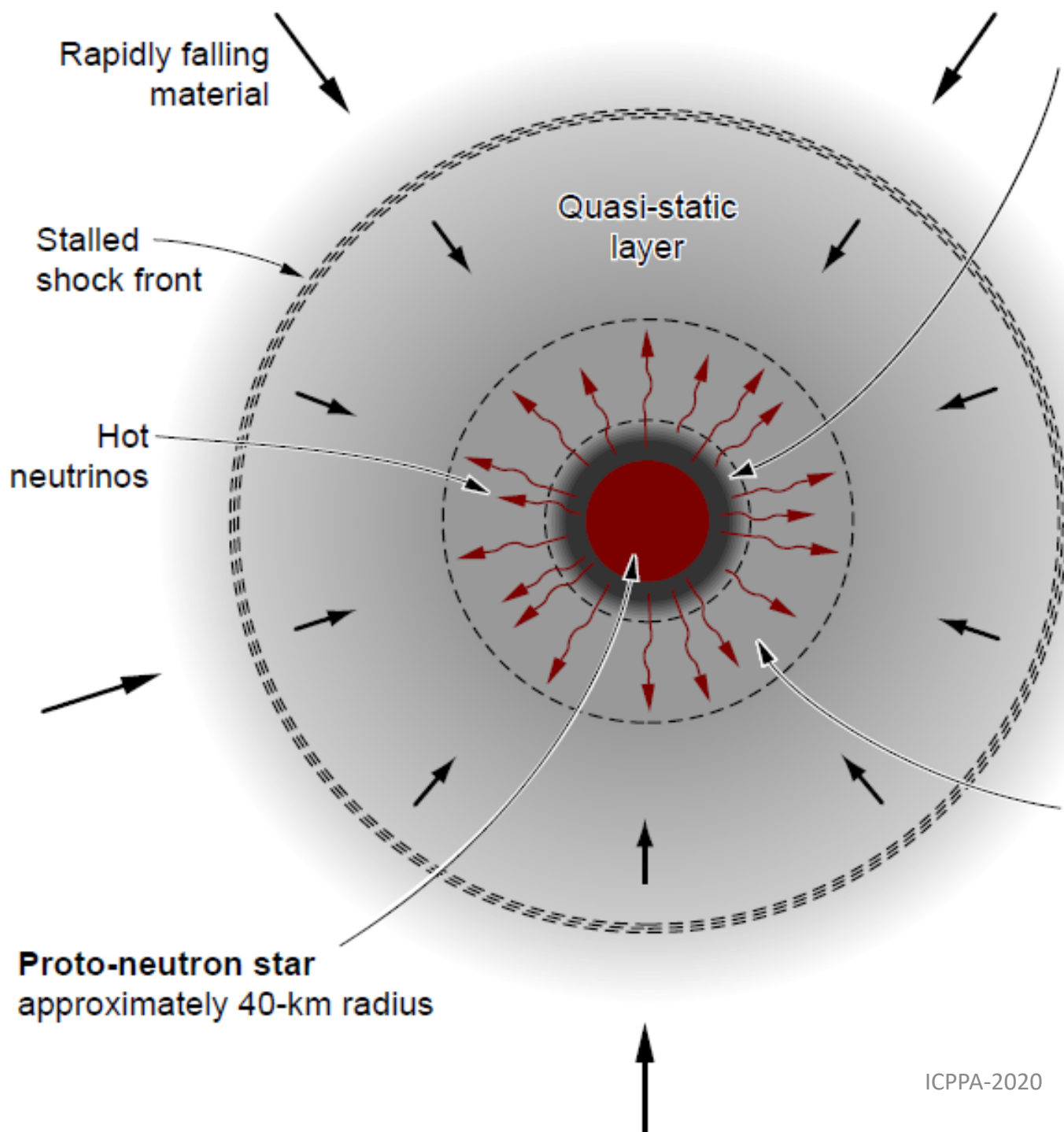


Incorporating the Heisenberg and Pauli principles into the kinetic approach to neutrino oscillations

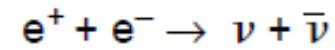
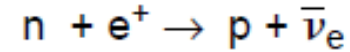
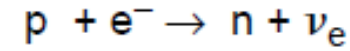
Alexander Kartavtsev

P. G. Demidov Yaroslavl State University, Russia

Based on ArXiv [2007.13736](https://arxiv.org/abs/2007.13736)

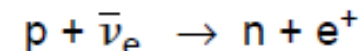
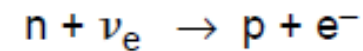


Region cooled by neutrinos
40- to 50-km radius;
neutrinos are produced by the following reactions:



$\sim 10^{57}$ neutrinos,
 ~ 10 MeV
99% gravitational energy

Region heated by neutrinos
50- to 100-km radius;
neutrinos are absorbed by the following reactions:



Collective neutrino oscillations

[Raffelt & Tamborra](#)

Phys. Rev. D **82**, 125004

“decoherence by dephasing of many neutrinos”

[Akhmedov & Mirrizi](#)

nuclphysb.2016.02.011

“decoherence by wave packet separation”

Outline

The effect of

decoherence by wave packet separation

can be incorporated into the

kinetic approach to neutrino oscillations

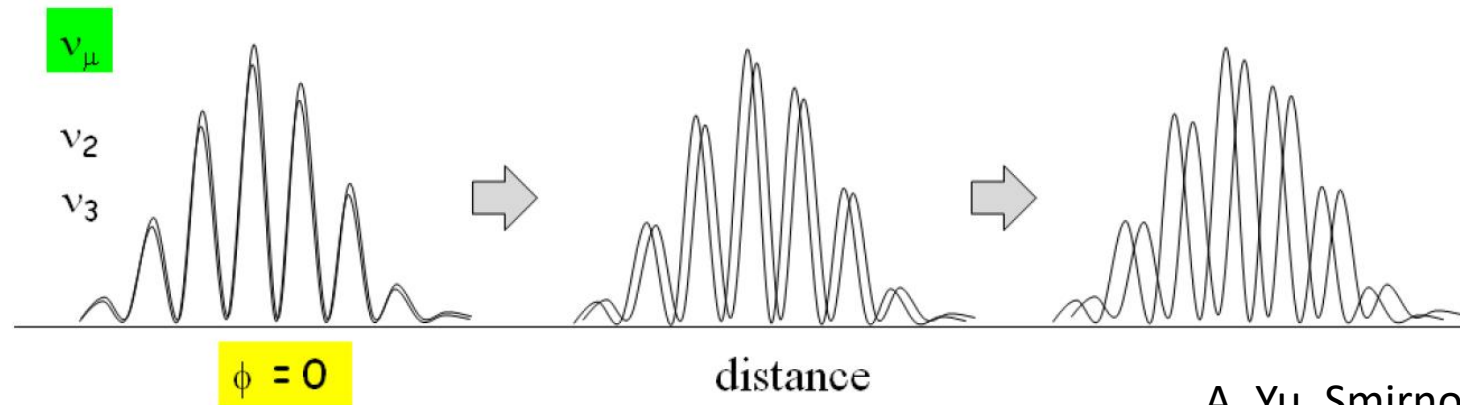
through the initial conditions consistent with the

Heisenberg principle

Decoherence by wave packet separation

$$i\partial_t\psi_i(t, \mathbf{x}) = H_{ij}(t, \mathbf{x})\psi_j(t, \mathbf{x})$$

$$\psi = \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix}$$



A. Yu. Smirnov
[1609.02386](https://arxiv.org/abs/1609.02386)

$$\psi_i(0, \mathbf{p}) \propto \exp\left(-\frac{(\mathbf{p} - \mathbf{p}_w)^2}{4\sigma_p^2}\right) e^{-i\mathbf{p}\mathbf{x}_w}$$

Kinetic approach to neutrino oscillations

Liouville term

$$\partial_t \varrho(t, \mathbf{x}, \mathbf{p}) + \frac{1}{2} \{ \partial_{\mathbf{p}} H(t, \mathbf{x}, \mathbf{p}), \partial_{\mathbf{x}} \varrho(t, \mathbf{x}, \mathbf{p}) \} - \frac{1}{2} \{ \partial_{\mathbf{x}} H(t, \mathbf{x}, \mathbf{p}), \partial_{\mathbf{p}} \varrho(t, \mathbf{x}, \mathbf{p}) \} \\ \approx -i \underbrace{[H(t, \mathbf{x}, \mathbf{p}), \varrho(t, \mathbf{x}, \mathbf{p})]}_{\text{Oscillation term}} + \underbrace{\mathcal{C}}_{\text{Collision term}} .$$



In quantum mechanics coordinate and momentum are not measurable simultaneously

In the kinetic approach the coordinate and momentum arguments have definite values

This is frequently interpreted as a sign, that the kinetic approach is inherently classical

Heisenberg principle

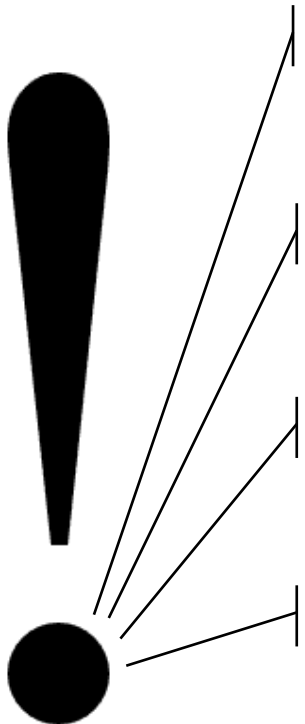
$$\hat{Q}_{ij}(t, \mathbf{x}, \mathbf{p}) = \int \frac{d^3 \Delta}{(2\pi)^3} e^{i\Delta \mathbf{x}} \hat{a}_j^\dagger(t, \mathbf{p} - \Delta/2) \hat{a}_i(t, \mathbf{p} + \Delta/2)$$

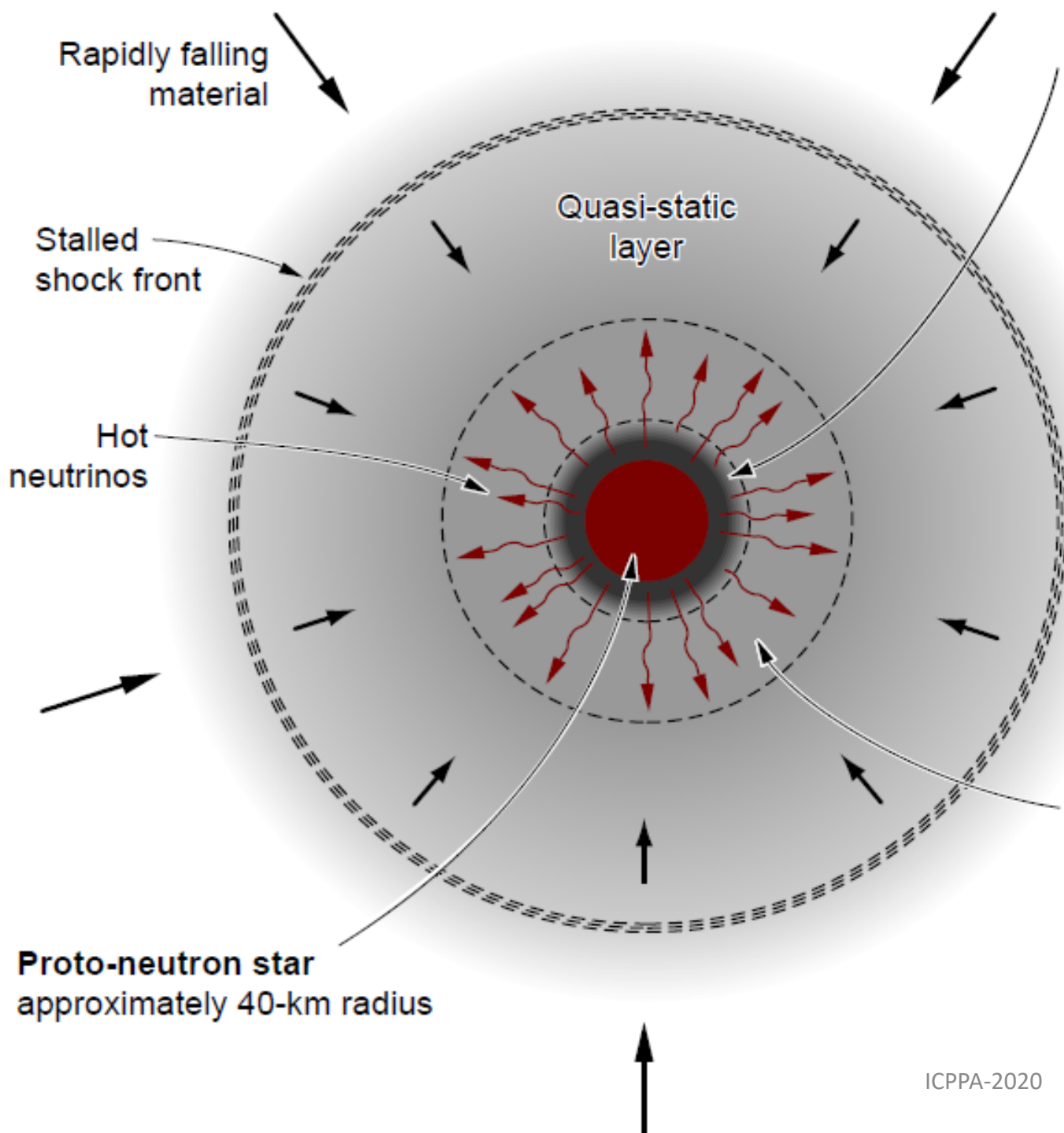
The coordinate and momentum arguments are *not* the neutrino coordinate and momentum

$$|\psi\rangle = \sum_i \int \frac{d^3 \mathbf{p}}{(2\pi)^3} \psi_i(0, \mathbf{p}) \hat{a}_i^\dagger(0, \mathbf{p}) |0\rangle$$

$$\varrho_{ij}(0, \mathbf{x}, \mathbf{p}) = \int \frac{d^3 \Delta}{(2\pi)^3} e^{i\Delta \mathbf{x}} \psi_i(0, \mathbf{p} + \Delta/2) \psi_j^*(0, \mathbf{p} - \Delta/2)$$

$$\varrho_{ij}(0, \mathbf{x}, \mathbf{p}) \propto \exp\left(-\frac{(\mathbf{p} - \mathbf{p}_w)^2}{2\sigma_p^2}\right) \exp\left(-\frac{(\mathbf{x} - \mathbf{x}_w)^2}{2\sigma_x^2}\right)$$





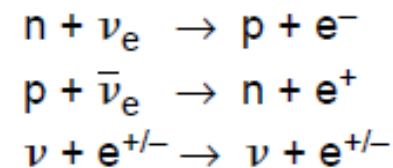
Region cooled by neutrinos
40- to 50-km radius;

$$|\psi\rangle = \frac{1}{\sqrt{N!}} \sum_{i_1 \dots i_N} \int \frac{d^3 \mathbf{p}_1}{(2\pi)^3} \cdots \frac{d^3 \mathbf{p}_N}{(2\pi)^3} \times \psi_{i_1 \dots i_N}(0, \mathbf{p}_1 \dots \mathbf{p}_N) \hat{a}_{i_1}^\dagger(0, \mathbf{p}_1) \cdots \hat{a}_{i_N}^\dagger(0, \mathbf{p}_N) |0\rangle$$

[Akhmedov & Mirrizi](#)
“decoherence by wave packet separation”

$$\rho_{ij}(0, \mathbf{x}, \mathbf{p}) \approx \sum_{l=1}^N \underbrace{\rho_{ij}^{(l)}(0, \mathbf{x}, \mathbf{p})}_{\text{decoherence by dephasing of many neutrinos}}$$

Region heated by neutrinos
50- to 100-km radius;
neutrinos are absorbed by the following reactions:



[Raffelt & Tamborra](#)
“decoherence by dephasing of many neutrinos”

Summary

The effect of

decoherence by wave packet separation

can be incorporated into the

kinetic approach to neutrino oscillations

through the initial conditions consistent with the

Heisenberg principle

Wigner approach

$$i\partial_t\psi_i(t, \mathbf{x}) = H_{ij}(t, \mathbf{x})\psi_j(t, \mathbf{x})$$

$$\varrho_{ij}(t, \mathbf{x}, \mathbf{p}) = \int \frac{d^3\Delta}{(2\pi)^3} e^{i\Delta\mathbf{x}} \psi_i(t, \mathbf{p} + \Delta/2) \psi_j^*(t, \mathbf{p} - \Delta/2)$$

$$\begin{aligned} \partial_t\varrho(t, \mathbf{x}, \mathbf{p}) + \frac{1}{2}\{\partial_{\mathbf{p}}H(t, \mathbf{x}, \mathbf{p}), \partial_{\mathbf{x}}\varrho(t, \mathbf{x}, \mathbf{p})\} - \frac{1}{2}\{\partial_{\mathbf{x}}H(t, \mathbf{x}, \mathbf{p}), \partial_{\mathbf{p}}\varrho(t, \mathbf{x}, \mathbf{p})\} \\ \approx -i[H(t, \mathbf{x}, \mathbf{p}), \varrho(t, \mathbf{x}, \mathbf{p})] . \end{aligned}$$