

Stopping of protons in pA collisions at SPS and NICA energies in analytical hydrodynamic model and in SMASH event generator



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Motivation of this study

STOPPING

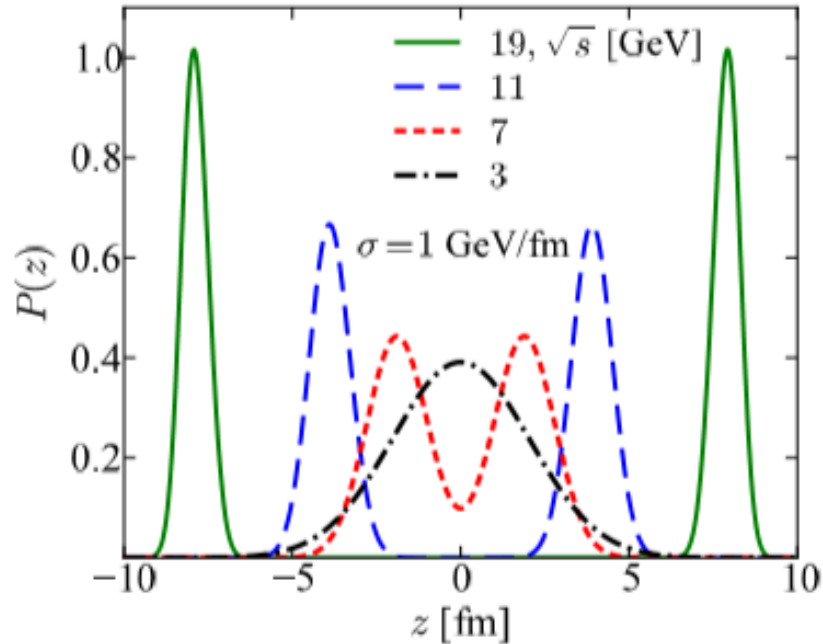


Fig.1. Stopped nucleons in configuration space [1]

[1] Andrzej Bialas, Adam Bzdak, Volker Koch (2016)

Non-Linear relationship[2]

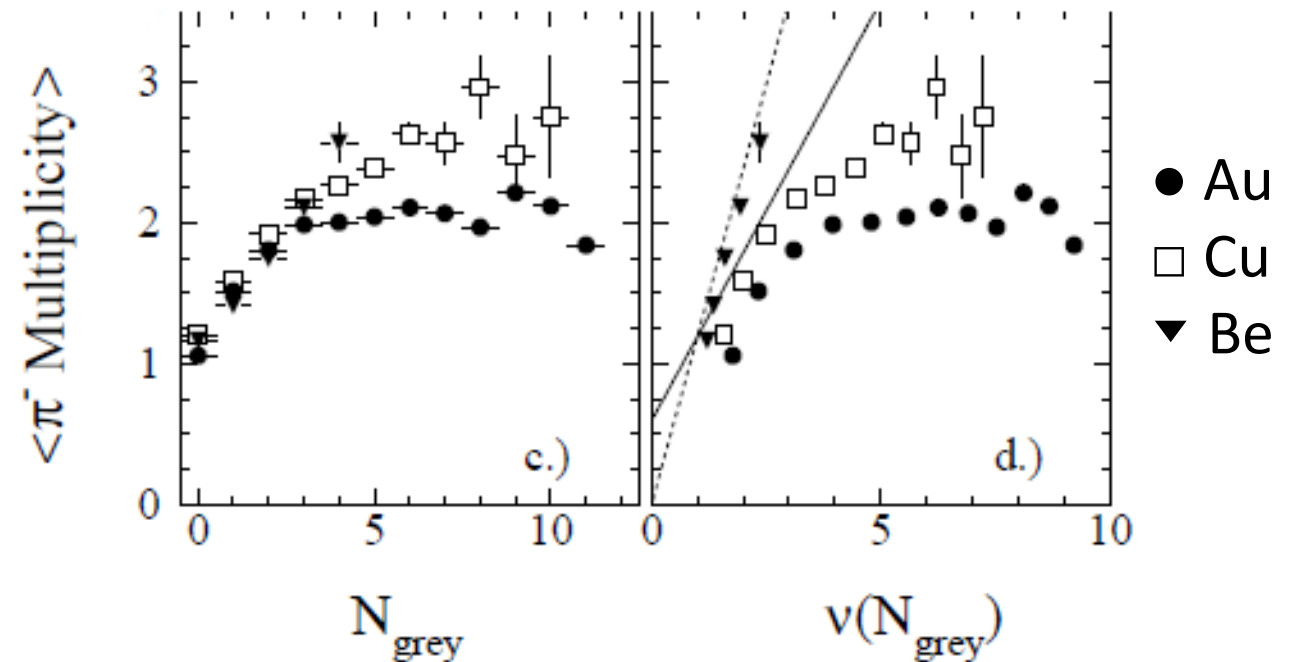


Fig.2. Multiplicity vs. number of "grey nucleons" (N_{grey}) and number of binary collisions ($\nu(N_{\text{grey}})$) in p-A collisions at 18 GeV/c[2]

[2] I. Chemakin et al. (1999)

Our hydrodynamic model of proton stopping

[3] L. Landau, E. Lifshits (1986)

$$dF^i = T^{ik} df_k; \quad T^{\alpha\beta} = \begin{pmatrix} \varepsilon & 0 & 0 & 0 \\ 0 & p & 0 & 0 \\ 0 & 0 & p & 0 \\ 0 & 0 & 0 & p \end{pmatrix}; \quad \varepsilon = \frac{M}{V} \\ p = 0 \quad T^{ik} = \frac{\varepsilon v^i v^k}{1 - v^2} \Rightarrow F = -\frac{1}{1 - v^2} \varepsilon v^2 S$$

Known quantities: $v_0, \rho(r), \sigma_{inel}^{NN}$

$$m\vec{a}\gamma + m\vec{v}(\vec{v}\vec{a})\gamma^3 = \vec{F} \Rightarrow v(t) \quad \left. \begin{array}{l} \\ \text{Speed of the «stop» } V_0 \text{ (when become gray)} \end{array} \right\} \Rightarrow l$$

$$\frac{dv}{dt} = (-1) \frac{\varepsilon S}{m_{0p}} * v^2 \gamma (1 + v^2 \gamma^2)^{-1} \\ \text{Const} = (-1) \frac{\varepsilon S}{m_{0p}}$$

Quantities used for calculation (p-Au - collision):

R, fm	a, fm	A	V_0	σ_{inel}^{NN}, fm^2	r_p, fm	m_{0p}, GeV	M, GeV	$p_{lab}, GeV/c$
7.64	0.538	197	0.5	2.85	1.2	1	197.09	18

Const ≈ 0.1 GeV

Glauber-like approach:

$$◆ T(b) = \int_{\Omega} \int \rho(z, \mathbf{b}_A) \delta(\mathbf{b} - \mathbf{b}_A) dz d\mathbf{b}_A$$

-thickness function in case of pA-collision (ρ – Woods-Saxon distribution; normalized to unity)

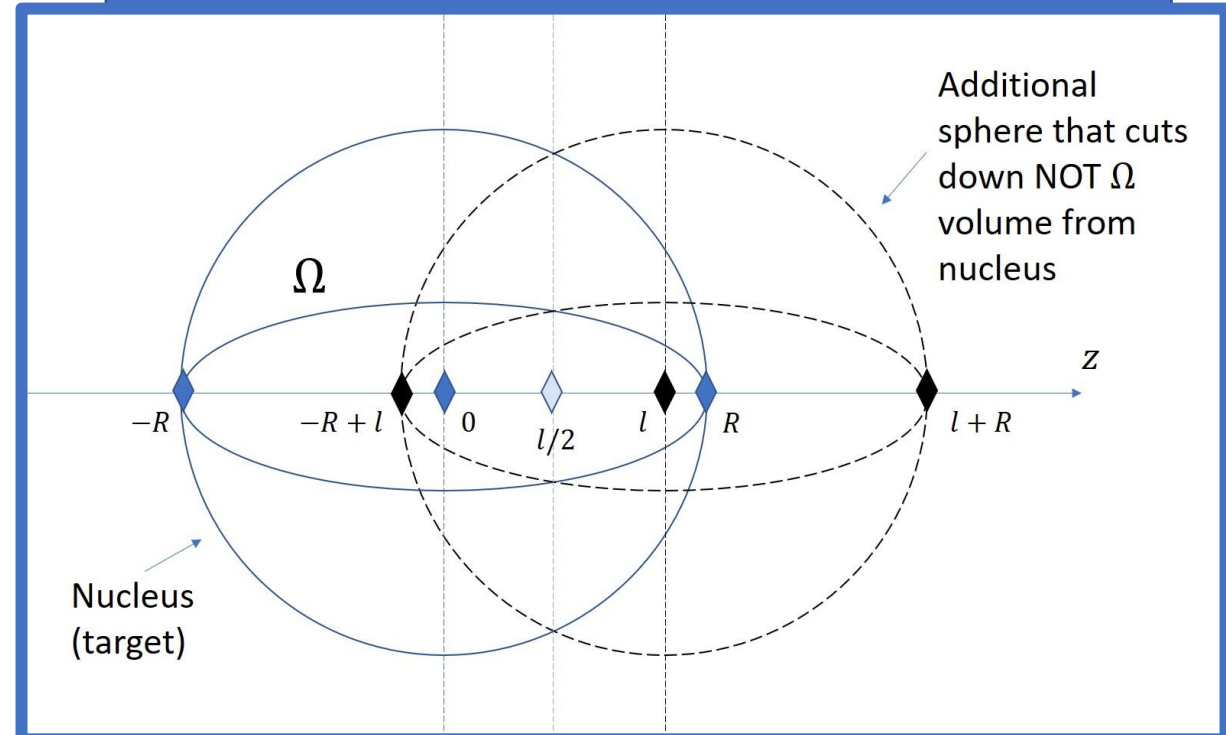
$$◆ \sigma_{inel}^{NN} T(b) = \int_{\Omega} \rho(z, \mathbf{b}_A) \delta(\mathbf{b} - \mathbf{b}_A) dz d\mathbf{b}_A \sigma_{inel}^{NN}$$

- probability to have ONE baryon-baryon inelastic collision when proton and target situated at an impact parameter \mathbf{b} relative to each other

$$◆ P(n, b) = C_A^n \left(\sigma_{inel}^{NN} T(b) \right)^n \left(1 - \sigma_{inel}^{NN} T(b) \right)^{A-n}$$

- probability to have n baryon-baryon inelastic collision when proton and target situated at an impact parameter \mathbf{b} relative to each other

[4] Cheuk-Yin Wong (1994)



Glauber-like approach:

[4] Cheuk-Yin Wong (1994)

- ◆ $N_{bin}(\mathbf{b}) = \sum_{i=1}^A nP(n, \mathbf{b})$ - average number on inelastic baryon-baryon collisions at an impact parameter \mathbf{b}
- ◆ $N_{ch}(\mathbf{b}) = \sum_{i=1}^N N_{ch}^{pp}(v_i) * \Delta N_{bin}(v_i | \mathbf{b})$ - average number of multiplicity at an impact parameter \mathbf{b} !


$$N_{ch}^{pp}(s_{NN}) = a + b * \log(s_{NN}) + c(\log(s_{NN}))^2$$

[5] W.Thomé, K.Eggert, K.Gibini [etc.], (1997)

$\Delta N_{bin}(v_i | \mathbf{b})$ - number of binary collisions that happened while proton was decelerating from v_i to v_{i+1}

!

In fact, this is an integral that was written in a discrete form for better understanding.

Compare with MC event-generators (full acceptance)

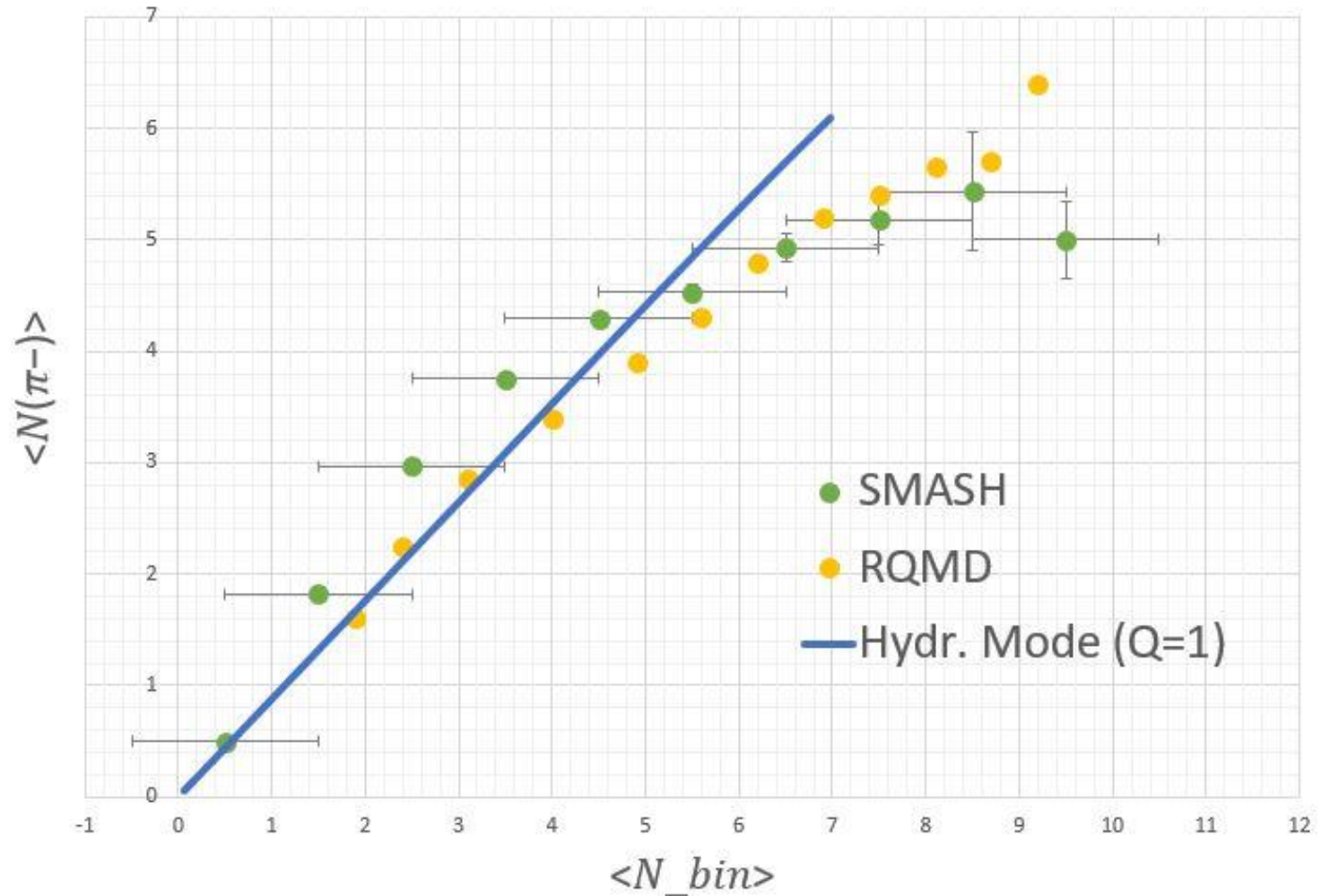
$$N_{\pi^-} = N_{ch}/2;$$



In the Hydrodynamic model of stopping

Fig.2 Mean charged particle yields vs. number of binary collisions in RQMD [2], SMASH and in Hydrodynamic model of stopping

p-Au collision, $p_{lab} = 18 \text{ GeV}/c$



[2] I. Chemakin et al. (1999)

Account of the limited acceptance:

◆
$$N_{ch} = \sum_{i=1}^N N_{ch}^{pp}(v_i) * \Delta N_{bin}(v_i|\mathbf{b}) * Q^{\sum_{k=1}^i \Delta N_{bin}(v_k|\mathbf{b})}$$
 - average multiplicity considering acceptance !

Q – probability to detect multiplicity that we got from ONE inelastic binary collision. Q^y - probability to detect multiplicity that we got from y inelastic binary collisions.

$\Delta N_{bin}(v_i|\mathbf{b})$ - number of binary collisions that happened while proton was decelerating from v_k to v_{k+1}

! In fact, this is an integral that was written in a discrete form for better understanding.

Account of the limited acceptance:

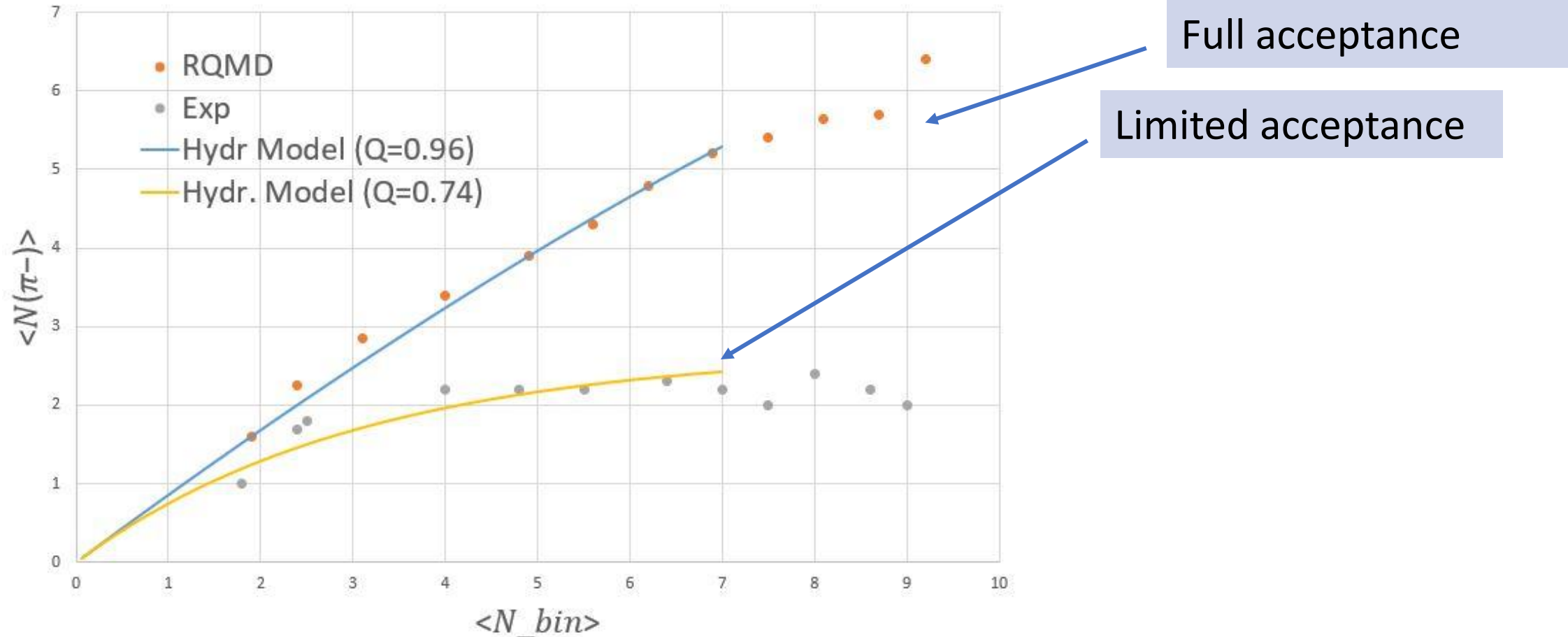


Fig.2 Mean charged particle yields vs. number of binary collisions in RQMD [2], experiment at BNL and in Hydrodynamic model of stopping considering acceptance p-Au – collision, $p_{lab} = 18 \text{ GeV}/c$

[2] I. Chemakin et al. (1999)

Account of the limited acceptance:

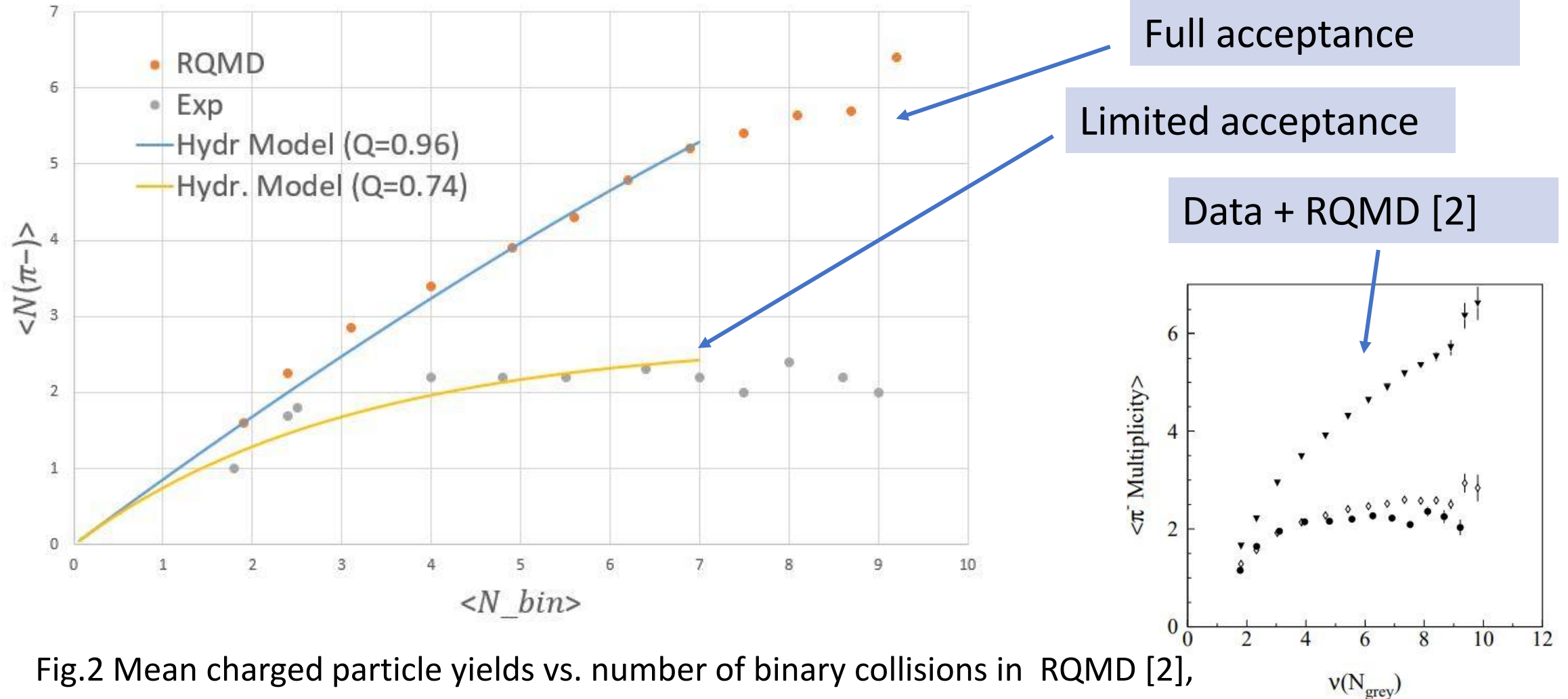


Fig.2 Mean charged particle yields vs. number of binary collisions in RQMD [2], experiment at BNL and in Hydrodynamic model of stopping considering acceptance p-Au – collision, $p_{lab} = 18 \text{ GeV}/c$

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Conclusion

- ◆ A new hydrodynamic model of nucleon stopping that describes the deceleration of a proton in a nucleus and based on hydrodynamics is proposed. No fitting coefficients are required.
- ◆ The linear dependences for the p-A collisions were obtained. This is in line with experiment on p-Au collisions at $p_{\text{lab}} = 18 \text{ GeV}/c$ in the first approximation. Similar dependence is demonstrated by the MC models RQMD and SMASH.
- ◆ The non-linear behavior of multiplicity vs. number of binary collisions -description is found to be a result of limited acceptance of the experimental data.
- ◆ Results of these studies of nucleon stopping are important for the future analysis of centrality selection in p-A and A-A collisions at NICA experiments.

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Back Up

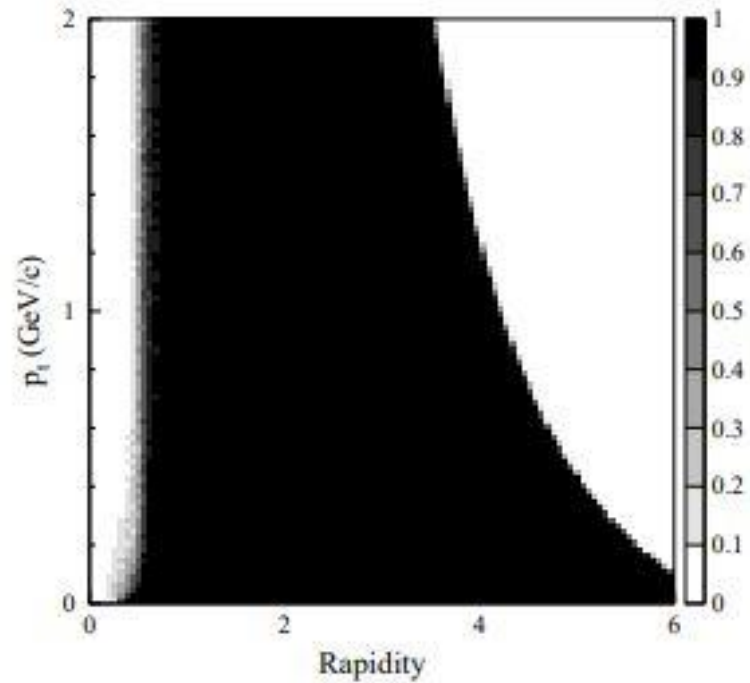
SMASH - Simulating **M**any **A**ccelerated **S**trongly-Interacting **H**adrons

Based on the relativistic Boltzmann equation :

$$p^\mu \partial_\mu f_i(x, p) + m_i F^\alpha \partial_\alpha^p f_i(x, p) = C_{coll}^i$$

At each step along the time axis, for each particle, its trajectory is calculated using the Boltzmann equation.

Back Up



Acceptance in BNL experiment

[2] I. Chemakin et al. (1999)