

The comparison of calculated atmospheric neutrino spectra with measurement data of IceCube and ANTARES experiments

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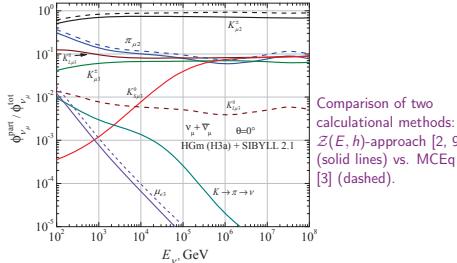
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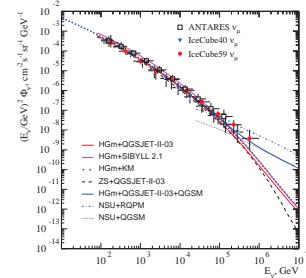
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Contributions to the $(\nu_\mu + \bar{\nu}_\mu)$ flux: $\mathcal{Z}(E, h)$ compared to MCEq



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Energy spectra of the atmospheric $\nu_\mu + \bar{\nu}_\mu$ (averaged over the range $90 - 180^\circ$): IceCube, ANTARES



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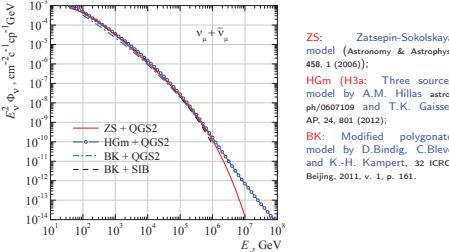
Abstract

The problem of the atmospheric neutrino background became really important after comparison with the IceCube data of neutrinos induced by very high-energy cosmic rays of extraterrestrial origin [1]. Comprehensive study of the high-energy neutrino spectrum and zenith angle distribution of atmospheric neutrinos is required since it is not improbable that the atmospheric neutrinos arising from decays of π , K mesons (conventional neutrinos) or from decays of charmed particles (prompt neutrinos), produced in collisions of cosmic ray particles with the Earth's atmosphere, and astrophysical neutrinos overlap.

We calculate the atmospheric neutrino spectra in the energy range of 10^2 GeV - 10^7 PeV using the set of the hadronic models and several parametrizations of cosmic ray spectra supported by experimental data. Above 10^3 TeV calculated spectra of muon neutrinos show the apparent dependence on the spectrum and composition of primary cosmic rays around to the knee. Also in this region uncertainties appear due to production cross sections and decays of the charmed particles which imprint on the prompt neutrino flux. It is also shown also the contribution of rare decays of short-lived neutral mesons (comprise about a third of the conventional flux) at the energies above 100 TeV. Taking into account K -meson production in π -A collisions gives addition 5 - 7% to the ν_μ flux in the energy range of 100 GeV - 100 TeV. The detailed comparison of our calculations, based on the $Z(E, h)$ -functions approach [2] with those of the MCEq method by A.Fedynitch et al. [3], shows the consistency on the whole at least in the energy range 100 GeV - 1 PeV. Calculated neutrino spectra agree rather well with the measurement data in IceCube [4, 5, 6, 7] and ANTARES [8]. Uncertainties of the experimental data above 500 TeV leave a window for the prompt neutrino component predicted with use of the quark-gluon string model (QGSM) [2].

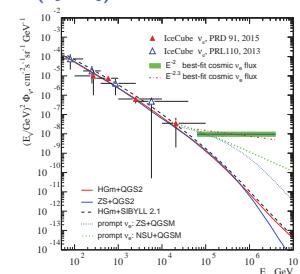
Influence of CR spectrum on the $(\nu_\mu + \bar{\nu}_\mu)$

Comparison of the calculations (zenith-angle averaged) for the CR spectra parametrizations involving the "knee"



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Atmospheric and astrophysical fluxes ($\nu_e + \bar{\nu}_e$)



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The band and dash-dotted line = IceCube best fits for astrophysical neutrino flux $(\nu_e + \bar{\nu}_e)$ [1]; $E_\nu^2 \phi_{\nu_e} = (0.95 \pm 0.3) \cdot 10^{-8}$, $E_\nu^2 \phi_{\bar{\nu}_e} = 1.5 \cdot 10^{-8} (E/100\text{TeV})^{-0.3}$ $(\text{cm}^2 \text{sr}^{-1} \text{GeV}^{-1})$
 $(\nu_e : \nu_\mu : \mu\tau = 1 : 1 : 1)$

Prompt neutrinos: calculations with QGSM for CR spectra by Nikolsky-Stamenov-Ushev (NSU) and Zatsepin-Sokolskaya (ZS) (dotted)

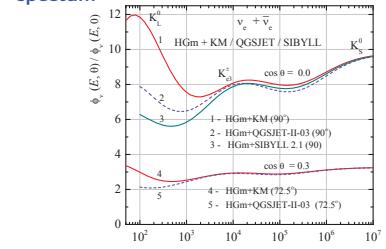
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Sources of atmospheric neutrinos

Conventional (π / K) neutrinos		Prompt (D/Λ_c) neutrinos	
Decay modes	Fraction, Γ/Γ_0	Decay modes	Fraction
μ^\pm	$\pi^\pm + \nu_\mu (\bar{\nu}_e) + \bar{\nu}_\mu (\nu_e)$	D^\pm	$\mu^\pm + \nu_\mu (\bar{\nu}_\mu) + X$
π^\pm	$99, 99\%$	D^0, \bar{D}^0	$\mu^\pm + \nu_\mu (\bar{\nu}_\mu) + X$
K^\pm	$\mu^\pm + \nu_K (\bar{\nu}_K) + \bar{\nu}_K (\nu_K)$ $\pi^0 + \mu^\pm + \nu_\mu (\bar{\nu}_\mu)$ $\pi^0 + \pi^0 + \nu_e (\bar{\nu}_e)$	D_s^\pm	$e^\pm + X$
K_L^0	$\pi^\pm + \mu^\mp + \bar{\nu}_\mu (\nu_\mu)$ $\pi^\pm + e^\mp + \bar{\nu}_e (\nu_e)$	Λ_c^\pm	$\Lambda + \mu^\pm + \nu_\mu$ $\Lambda + e^\pm + \nu_e$
K_S^0	$\pi^+ + \pi^-$ $\pi^\pm + \mu^\mp + \bar{\nu}_\mu (\nu_\mu)$		$2.0 \pm 0.7\%$ $2.1 \pm 0.6\%$

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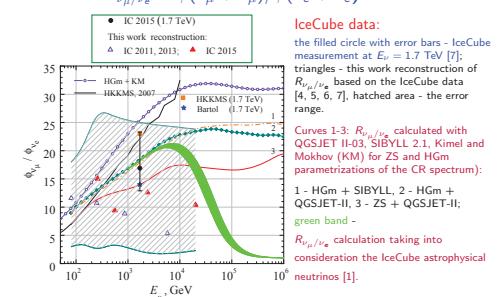
Zenith-angle enhancement of the $(\nu_e + \bar{\nu}_e)$ spectrum



Zenith-angle enhancement of the $(\nu_e + \bar{\nu}_e)$ flux reflects successive "switching-on" of the kaon sources

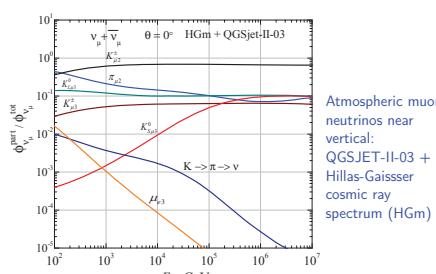
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Ratio $R_{\nu_\mu/\nu_e} = \phi(\nu_\mu + \bar{\nu}_\mu)/\phi(\nu_e + \bar{\nu}_e)$



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Partial contributions to the flux of $(\nu_\mu + \bar{\nu}_\mu)$



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Neutrino fluxes depending on the hadronic model

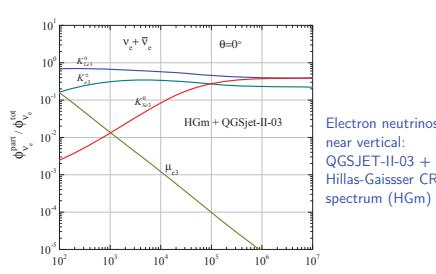
Atmospheric neutrino flux ratios calculated with SIBYLL 2.1, QGSJET II-03 and Kimel-Mokhov model (KM) for CR spectra ZS and HGM: 1 (4) - sib/qgs2; 2 (5) - km/qgs2; 3 (6) - sib/km.

E_ν, TeV	1	2	3	4	5	6
ZS: $(\nu_\mu + \bar{\nu}_\mu)$				$(\nu_e + \bar{\nu}_e)$		
1	1.70	1.05	1.62	1.41	0.51	2.76
10	1.53	1.04	1.47	1.32	0.48	2.75
10^2	1.53	1.10	1.39	1.29	0.54	2.39
10^3	1.79	1.64	1.09	1.41	0.84	1.68
10^4	1.85	2.08	0.89	1.38	1.06	1.30
HGM: $(\nu_\mu + \bar{\nu}_\mu)$				$(\nu_e + \bar{\nu}_e)$		
1	1.59	0.85	1.87	1.39	0.49	2.84
10	1.57	1.12	1.40	1.33	0.50	2.66
10^2	1.57	1.27	1.24	1.31	0.57	2.30
10^3	1.63	1.63	1.00	1.31	0.70	1.87
10^4	1.47	1.53	0.96	1.23	0.59	2.08

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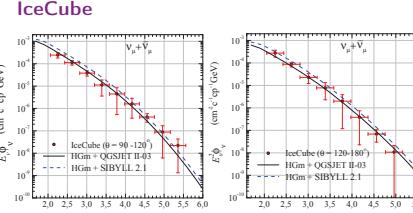
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Partial contributions to the $(\nu_e + \bar{\nu}_e)$ flux



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Spectra of the atmospheric $(\nu_\mu + \bar{\nu}_\mu)$ fluxes: IceCube



$(\nu_\mu + \bar{\nu}_\mu)$ fluxes averaged over zenith angle ranges $90 - 120^\circ$ and $120 - 180^\circ$. Curves - the calculations with QGSJET II-03 and SIBYLL 2.1, symbols - IceCube measurements.

References

- [1] M.G. Aartsen et al. (IceCube Collaboration), Phys. Rev. Lett. **111**, 021103 (2013); Science **342**, 124286 (2013); Phys. Rev. Lett. **113**, 101101 (2014); PoS (ICRC2015) 1081; arXiv:1510.05223
- [2] T.S. Sinegovskaya, A.D. Morozova, S.I. Sinegovsky, Phys. Rev. D **91**, 063011 (2015)
- [3] A. Fedynitch et al. EPJ Web Conf. **99**, 08001 (2015); arXiv:1503.00544; PoS (ICRC2015) 1129
- [4] R. Abbasi et al. (IceCube Collaboration), Phys. Rev. D **83**, 012001 (2011)
- [5] M.G. Aartsen et al. (IceCube Collaboration), Eur. Phys. J. C **75**, 116 (2015)
- [6] M.G. Aartsen et al. (IceCube Collaboration), Phys. Rev. Lett. **110**, 151105 (2013)
- [7] M.G. Aartsen et al. (IceCube Collaboration), Phys. Rev. D **91**, 122004 (2015)
- [8] S. Adrian-Martinez et al. Eur. Phys. J. C **73**, 2606 (2013)
- [9] A. A. Kochanov, T. S. Sinegovskaya, S. I. Sinegovsky, Astropart. Phys. **30**, 219 (2008); arXiv:0803.2943v2 [astro-ph]

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