

# Dynamical $O(4)$ -symmetry in the light meson spectrum within the framework of the Regge approach

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24.10.2024

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# Meson clustering

If one's up to analyze PDG tables, an interesting pattern is seen. Light mesons such as  $\pi$ ,  $\eta$ ,  $\rho$ ,  $\omega$ ,  $a_J$ ,  $f_J$ ,  $b_J$ ,  $h_J$  all have close masses:

$\approx 1300$ MeV	11 states
$\approx 1600$ MeV	8 states
$\approx 2000$ MeV	22 states
$\approx 2250$ MeV	24 states

## O(4)-symmetry

Inside the excited unflavoured mesons, spin-orbital and spin-spin correlations are suppressed. The discrete mesonic spectrum depends only on a quantum number  $N = n_r + L$ .

$$\vec{S} = \vec{s}_q + \vec{s}_{\bar{q}}$$

Singlets and triplets:  $\vec{J} = \vec{L} + \vec{S}$ , with  $S = 0$  or  $S = 1$ , respectively.

The spectrum can be fitted by

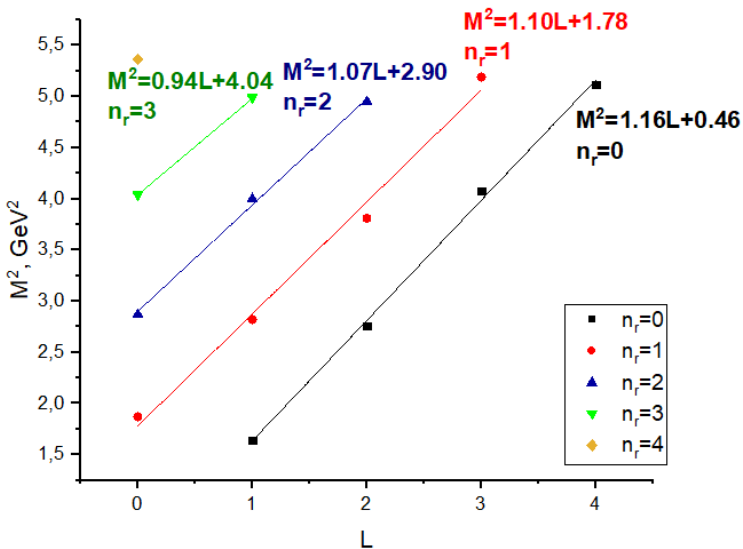
$$M^2 = a(n_r + L) + c,$$

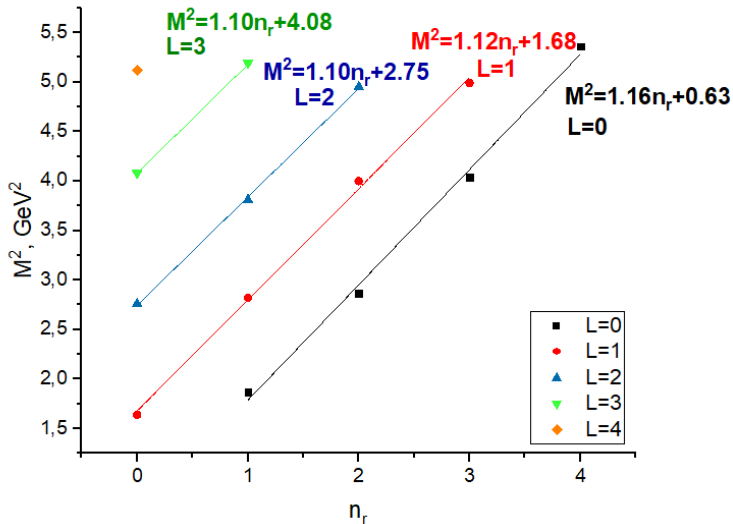
where  $n_r$  is the radial excitation and  $L$  is the meson angular momentum.

$L \downarrow$		$S \downarrow$	$J \downarrow$	$n = 0$	$n = 1$	$n = 2$	$n = 3$	$n = 4$
0	0	0		$\pi$	$\pi(1300)$	$\pi(1800)$	$\pi(2070)$	$\pi(2360)$
	1	1		$\rho$	$\eta(1295)$	$\eta(1760)$	$\eta(2010)$	$\eta(2320)$
				$\omega$	$\rho(1450)$	$\rho(1700)$	$\rho(2000)$	$\rho(2270)$
					$\omega(1420)$	$\omega(1650)$	$\omega(1960)$	$\omega(2290)$
1	0	1		$b_1(1235)$	$b_1?$	$b_1(1960)$	$b_1(2240)$	$n = 4$
	1	0		$h_1(1170)$	$h_1(1595)$	$h_1(1965)$	$h_1(2215)$	
		1		$a_0(1450)$	$a_0(1710)$	$a_0(2020)$	$a_0?$	
		2		$f_0(1370)$	$f_0(1770)$	$f_0(2020)$	$f_0(2200)$	
				$a_1(1260)$	$a_1(1640)$	$a_1(1930)$	$a_1(2270)$	
				$f_1(1285)$	$f_1?$	$f_1(1970)$	$f_1(2310)$	
				$a_2(1320)$	$a_2(1700)$	$a_2(2030)$	$a_2(2255)$	
				$f_2(1270)$	$f_2(1750)$	$f_2(1950)$	$f_2(2300)$	
2	0	2		$\pi_2(1670)$	$\pi_2(2005)$	$\pi_2(2285)$	$n = 3$	
	1	1		$\eta_2(1645)$	$\eta_2(2030)$	$\eta_2(2250)$		
		2		$p?$	$p?$	$p?$		
		3		$\omega?$	$\omega?$	$\omega(2220)$		
				$p_2?$	$p_2(1940)$	$p_2(2225)$		
				$\omega_2?$	$\omega_2(1975)$	$\omega_2(2195)$		
				$p_3(1690)$	$p_3(1990)$	$p_3(2250)$		
				$\omega_3(1670)$	$\omega_3(1945)$	$\omega_3(2255)$		
3	0	3		$b_3(2030)$	$b_3(2245)$	$n = 2$	$N = 0$ $N = 1, M^2 \approx 1.72 \text{ GeV}$ $N = 2, M^2 \approx 2.63 \text{ GeV}$ $N = 3, M^2 \approx 3.98 \text{ GeV}$ $N = 4, M^2 \approx 5.11 \text{ GeV}$	
	1	2		$h_3(2025)$	$h_3(2275)$			
		3		$a_2(1990)$	$a_2(2175)$			
		4		$f_2(2010)$	$f_2(2240)$			
				$a_3(2030)$	$a_3(2275)$			
				$f_3(2050)$	$f_3(2300)$			
				$a_4(1970)$	$a_4(2255)$			
				$f_4(2050)$	$f_4(2300)$			
4	0	4		$b_4(2245)$	$n = 1$	$M^2 = 1.14 N + 0.58$	$N = n + L$	
	1	3		$h_4(2330)$				
		4		$\rho_3?$				
		5		$\omega_3(2285)$				
				$p_4(2230)$				
				$\omega_4(2250)$				
				$p_5(2350)$				
				$\omega_5(2250)$				
$L \uparrow$	$S \uparrow$	$J \uparrow$	$n = 0$					

 $I = 1$   
 $I = 0$ 

Резонансы с более бледным фоном не участвовали в анализе, и есть сомнения, что это связанное состояние кварк-антикварковой пары.

$n_r = \text{constant}$ 

$L = \text{constant}$ 

$L \downarrow$		$S \downarrow$	$J \downarrow$	$n = 0$	$n = 1$	$n = 2$	$n = 3$	$n = 4$
0	0	0		$\pi$	$\pi(1300)$	$\pi(1800)$	$\pi(2070)$	$\pi(2360)$
	1	1		$\rho$	$\rho(1450)$	$\rho(1700)$	$\rho(2000)$	$\rho(2270)$
				$\omega$	$\omega(1420)$	$\omega(1650)$	$\omega(1960)$	$\omega(2290)$
1	0	1		$b_1(1235)$	$b_1?$	$b_1(1960)$	$b_1(2240)$	$n = 4$
	1	0		$h_1(1170)$	$h_1(1595)$	$h_1(1965)$	$h_1(2215)$	
		1		$a_0(1450)$	$a_0(1710)$	$a_0(2020)$	$a_0?$	
		2		$f_0(1370)$	$f_0(1770)$	$f_0(2020)$	$f_0(2200)$	
		1		$a_1(1260)$	$a_1(1640)$	$a_1(1930)$	$a_1(2270)$	
		2		$f_1(1285)$	$f_1?$	$f_1(1970)$	$f_1(2310)$	
		2		$a_2(1320)$	$a_2(1700)$	$a_2(2030)$	$a_2(2255)$	
		2		$f_2(1270)$	$f_2(1750)$	$f_2(1950)$	$f_2(2300)$	
2	0	2		$\pi_2(1670)$	$\pi_2(2005)$	$\pi_2(2285)$	$n = 3$	
	1	1		$\eta_2(1645)$	$\eta_2(2030)$	$\eta_2(2250)$		
		2		$\rho?$	$\rho?$	$\rho?$		
		3		$\omega?$	$\omega?$	$\omega(2220)$		
		2		$p_2?$	$p_2(1940)$	$p_2(2225)$		
		3		$\omega_2?$	$\omega_2(1975)$	$\omega_2(2195)$		
		3		$p_3(1690)$	$p_3(1990)$	$p_3(2250)$		
		3		$\omega_3(1670)$	$\omega_3(1945)$	$\omega_3(2255)$		
3	0	3		$b_3(2030)$	$b_3(2245)$	$n = 2$		
	1	2		$h_3(2025)$	$h_3(2275)$			
		3		$a_2(1990)$	$a_2(2175)$			
		4		$f_2(2010)$	$f_2(2240)$			
		3		$a_3(2030)$	$a_3(2275)$			
		4		$f_3(2050)$	$f_3(2300)$			
		4		$a_4(1970)$	$a_4(2255)$			
		4		$f_4(2050)$	$f_4(2300)$			
4	0	4		$\pi_4(2250)$	$n = 1$			
	1	3		$\eta_4(2330)$				
		4		$\rho_3?$				
		5		$\omega_3(2285)$				
		4		$p_4(2230)$				
		5		$\omega_4(2250)$				
		5		$p_5(2350)$				
		5		$\omega_5(2250)$				
$L \uparrow$	$S \uparrow$	$J \uparrow$	$n = 0$					

$N = 0$
$N = 1, M^2 \approx 1.72 \text{ GeV}$
$N = 2, M^2 \approx 2.63 \text{ GeV}$
$N = 3, M^2 \approx 3.98 \text{ GeV}$
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$M^2 = 1.14 N + 0.58$
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$I = 1$
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Резонансы с более бледным фоном не участвовали в анализе, и есть сомнения, что это связанное состояние кварк-антикварковой пары.

# Slopes

Fitting this data by multiple linear regression, we acquire that in the relation  $M^2 = a \cdot n_r + b \cdot L + c$  slopes for  $n_r$  and  $L$  are the same,

$$a = 1.14 \pm 0.05, \quad b = 1.14 \pm 0.07, \quad c \approx 0.5$$



# The mass operator

Meson: a gluon string with quark and antiquark at the ends. If  $L$  and  $n_r$  are big enough, we can use semiclassical methods. The mass operator:

$$M = 2p + \sigma r$$

Here we assume two UR quarks with  $m_q = m_{\bar{q}} \ll p$  and linear confining potential  $\sigma r$ ,  $\sigma$  is equal to effective string tension. The maximal string length is  $\ell = M/\sigma$ .

## Semiclassical quantization for radial excitations

The mass operator is  $M = 2p + \sigma r$ .

The wave-function for fermions is **antisymmetric**. The Bohr-Sommerfeld quantization condition should be

$$\oint p(r)dr = \pi(n_r + \gamma) \quad \Rightarrow \quad M^2 = 2\pi\sigma(n_r + \gamma)$$

Thus in the  $M^2 = an_r + bl + c$  the slope for  $n_r$  is  $a = 2\pi\sigma$ . If one forgets about fermionic nature of quarks, the different slope  $4\pi\sigma$  appears.

## A collective gluon excitation

The second option: placing quarks at the 0 and  $\ell$ , assume interaction by exchanging massive scalar particle. The mass operator is now  $M = p + \sigma r$  and Bohr-Zommerfeld quantization condition reads

$$\oint p(r)dr = 2\pi(n_r + \gamma) \quad \Rightarrow \quad M^2 = 2\pi\sigma(n_r + \gamma)$$

Again,  $a = 2\pi\sigma$ .

## Chew-Frautchi formula

Suppose two massless quark rotating at the speed of light at radius  $\ell/2$ . The flux tube at the distance  $r$  from the center of mass rotates with the speed  $v(r) = 2r/\ell$ . The mass and angular momentum of such a "solid" gluon flux tube are

$$M = 2 \int_0^{\ell/2} \frac{\sigma dr}{\sqrt{1 - v^2(r)}} = \frac{\pi\sigma\ell}{2}, \quad L = 2 \int_0^{\ell/2} \frac{\sigma r v(r) dr}{\sqrt{1 - v^2(r)}} = \frac{\pi\sigma\ell^2}{8}$$
$$\Rightarrow M^2 = 2\pi\sigma L$$

This is the Chew-Frautschi formula and we've got the same slope for  $L$  which is  $2\pi\sigma$ .

## Semiclassical look at angular excitations

Quarks stay in circular classical orbits with large  $r$  and  $p$ . There is also a centripetal acceleration acting on them. Using orbit quantization conditions and applying Newton's law, we obtain

$$\oint p(r)dr = L\hbar, \quad p = \frac{\sigma r}{2} \quad \Rightarrow \quad L = \frac{\sigma r^2}{2}$$

After some calculations,  $M^2 = 8\sigma L$  and the slope for  $L$  differs crucially from  $2\pi\sigma$  for radial excitation, although  $M^2(L)$  is still linear.

## Conclusions

- There is a broad mass degeneracy in light unflavoured mesons. The spectrum depends linearly on meson radial quantum number  $n_r$  and its angular momentum  $L$ ,

$$M^2 = a n_r + b L + c.$$

- The analysis of experimental data shows that  $a \approx b \approx 1.14$ . The  $O(4)$  degeneracy predicts such a dependency too, discrete meson spectrum must depend on a single quantum number  $N = n_r + L + 1$ .
- Semiclassical approach is also capable to give us the needed relation between slopes,  $a = b = 2\pi\sigma$ , under certain conditions.