

Comparison of different centrality determination methods at the BM@N experiment

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- Evolution of matter produced in heavy-ion collisions depend on its initial geometry
- Centrality procedure maps initial geometry parameters with measurable quantities (multiplicity or transverse energy of the produced particles)
- **This allows comparison of the BM@N and future MPD results with the data from other experiments (STAR BES, NA49/NA61 scans) and theoretical models**

The Bayesian inversion method (Γ-fit): 2D fit 2D Gamma distribution

y

mean value and variance in the new coordinate system

 $D(x) = D(E)\cos(\alpha)^2 + R(E,M)\sqrt{D(E)D(M)}\sin(2\alpha) + D(M)\sin(\alpha)^2$ y = $-\sin(\alpha)\langle E\rangle + \cos(\alpha)\langle M\rangle$ $D(y) = D(E)\sin(\alpha)^2 - R(E,M)\sqrt{D(E)D(M)}\sin(2\alpha) + D(M)\cos(\alpha)^2$ $\langle x \rangle = \cos(\alpha) \langle E \rangle + \sin(\alpha) \langle M \rangle$ $D(x) = D(E)$ co

• A number of produced protons is stronger correlated with the number of produced particles (track & RPC+TOF hits)

• The fluctuation kernel for energy and multiplicity at fixed impact parameter can be describe by 2D Gamma distr.:

- than with the total charge of spectator fragments (FW) • to suppress self-correlation biases, it is necessary to
	- use spectators fragments for centrality estimation

The centrality determination procedure provides an estimate of the initial geometry in heavy ion collisions through a connection between the observable and the impact parameter. The multiplicity of produced charged particles and the Monte Carlo Glauber model is usually used to determine the centrality. However, there are difficulties in using this approach at NICA collider energies due to the large uncertainty of the impact parameter at small multiplicity, autocorrelation effect, etc. New approach to determine the centrality based on Bayes' theorem and two-dimensional Gamma distribution is proposed. This method allows to obtain information about the impact parameter by using only the measured two-dimensional energy distribution of spectator nucleons and the multiplicity of produced charged particles. A new method to account for the efficiency of the detector system and the pileup is also proposed. The performance of the proposed approach has been tested on the data from the BM@N experiment and simulation data from the DCM-QGSM-SMM model for Xe+CsI collisions at a beam energy of 3.8A GeV.

Comparison with MC-Glauber fit

Summary and outlook

- A new approach for centrality determination with energy of spectators and multiplicity of charged particles was proposed and compared with MC-Glauber method.
- The proposed method was applied to the data from BM@N experiment
- It is planned to create a two-dimensional method based on a signal from a hodoscope and energy from the FHCal.

Fit results

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 $(E'(c_b))$ from the rec. mod
 $\langle E \rangle = \varepsilon_1 \langle E'(c_b) \rangle + \varepsilon_0$, $\langle b_{b} \rangle$ – average value and var. of energy/mult. *b*^{*b*}) from the rec. model data **EXECUTE:**
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 ethod (Γ **-fit): 2D fit**
 $\langle E'(c_b) \rangle$ - average value and v
 $D(E'(c_b))$ from the rec. mod E <sub> i = $\varepsilon_1 \langle E'(c_b) \rangle + \varepsilon_0$, $D(E) = \varepsilon_2 D(E'(c_b))$ $M \rangle = m_1 \langle M'(\mathcal{C}_b) \rangle,$

> 2 $D(11)$ $D(M) = m_1^2 \cdot D(M') + m_1 \cdot m_2 \langle M' \rangle$

 $E'(c_b)$, – can be approximated by $D(E'(c_b))$ polynomials

$$
\langle E'(c_b) \rangle = \sum_{j=1}^{12} a_j c_b^j, \quad D(E'(c_b)) = \sum_{j=1}^{19} b_j c_b^j
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\langle M'(c_b) \rangle = \sum_{j=1}^{12} a_j c_b^j, \quad D(M'(c_b)) = \sum_{j=1}^{6} b_j c_b^j
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 ε_0 , ε_1 , ε_2 , m_1 , m_2 - fit parameters

 E_{FH} ($x \rightarrow^{k_x(c_b)-1} e^{-x/\theta_x}$ ($x \rightarrow^{k_y(c_b)-1} e^{-y/\theta_y}$ X
 M_{ch} $G_{2D}(E_{FH}, M_{ch}, \langle E \rangle, \langle M \rangle, D(E), D(M), R) = G(x, \theta_x, k_x) \cdot G(y, \theta_y, k_y)$ $(x, \theta_x, k_x) \cdot G(y, \theta_y, k_y) = \frac{(x)^{\kappa_x(c_b)-1}e^{-x/\theta_x}}{\Gamma(1-(x))\theta^2} \cdot \frac{(y)^{\kappa_y(c_b)-1}e^{-y/\theta_y}}{\Gamma(1-(x))\theta^2}$ $(k_{y}(c_{h}))\theta_{y}^{2} = \Gamma(k_{y}(c_{h}))\theta_{y}^{2}$ $k_x(c_b) - 1$ $2^{-x/\theta_x}$ $(x, y) - k_y(c_b) - 1$ $2^{-y/\theta_y}$ $x^{y-x}x^{y} - (y^{y-x}y^{y-y})$ $\Gamma(1 - (x^{y}y^{y}y^{y-y}))$ $\int x^{(b)}/b^{y}$ **c** *c* $\int b^{y}y^{(b)}y^{(c)}$ *y* $G(x, \theta_x, k_x) \cdot G(y, \theta_y, k_y) = \frac{(x)^{x+y} - e^{-x}}{\sum_{x=0}^{x} (x - y)^2} \cdot \frac{(y)^{y-x} - e^{-y}}{\sum_{x=0}^{x} (x - y)^2}$ $k(c_i)$ θ^2 $\Gamma(k(c_i))\theta^2$ θ_{n} $\lambda_{n}(c_{h})-1$ $-y/\theta_{n}$ θ , k) $\cdot G(y, \theta, k) = \frac{(y, k)}{k}$ θ^2 $\Gamma(k \ (c_i))\theta^2$ $(M, K) = G(X, \mathcal{O}_X, K_X) \cdot G(Y, \mathcal{O}_Y, K_Y)$
-1 -x/ θ ($\searrow k_y(c_b)$ -1 -y/ θ _y \cdot Cr(\vee , θ , κ) = ———————————————————— $\Gamma(k(c_i))\theta^2 \qquad \Gamma(k(c_i))\theta^2$ (x) $\langle x \rangle^2$ ∂ $D(y)$ $\langle y \rangle^2$ | 2 , , , () () *^x ^x ^y ^y* $\frac{D(x)}{x}$, $k = \frac{\langle x \rangle}{x}$, $\theta = \frac{D(y)}{y}$, $k = \frac{\langle y \rangle}{y}$ $D(x)$ $\langle y \rangle$ $D(y)$ $\theta = \frac{\mathcal{L}(x)}{\mathcal{L}(x)}$, $k = \frac{\mathcal{L}(x)}{\mathcal{L}(x)}$, $\theta = \frac{\mathcal{L}(y)}{\mathcal{L}(x)}$, $k = \frac{\mathcal{L}(y)}{\mathcal{L}(x)}$ $\arctan\left(\frac{2\sqrt{D(E)D(M)R(E,M)}}{2\sqrt{D(E)D(M)R(E,M)}} \right)$ $(E) - D(M)$) $D(E)D(M)R(E,M)$ $D(E) - D(M)$ α = arctan $\frac{1}{1}$ $\left(2\sqrt{D(E)D(M)}R(E,M)\right)$ $= \arctan\left(\frac{2\sqrt{D(L)D(M)N(L)M}}{D(E)-D(M)}\right)$ It is possible to find such a rotation angle of the system that $cov(x, y) = 0$ Then the two-dimensional distribution in the new coordinate system will be $= \sum_{i=1} a_i c_b^j, \quad D(E'(c_b)) = \sum_{i=1} b_i c_b^j$

Good agreement between fit and data. Difference in energy distribution due to efficiency