

Searches for new physics in rare decays of heavy flavors at CMS

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on behalf of the CMS collaboration

Outline

- Introduction, $b \rightarrow sll$
- Angular observables in $b \rightarrow sll$ transitions
- Searches for LFUV in heavy-flavor decays

CMS experiment

CMS DETECTOR

Total weight : 14,000 tonnes
Overall diameter : 15.0 m
Overall length : 28.7 m
Magnetic field : 3.8 T

STEEL RETURN YOKE
12,500 tonnes

SILICON TRACKERS
Pixel (100x150 μm) $\sim 16\text{m}^2 \sim 66\text{M}$ channels
Microstrips (80x180 μm) $\sim 200\text{m}^2 \sim 9.6\text{M}$ channels

SUPERCONDUCTING SOLENOID
Niobium titanium coil carrying $\sim 18,000\text{A}$

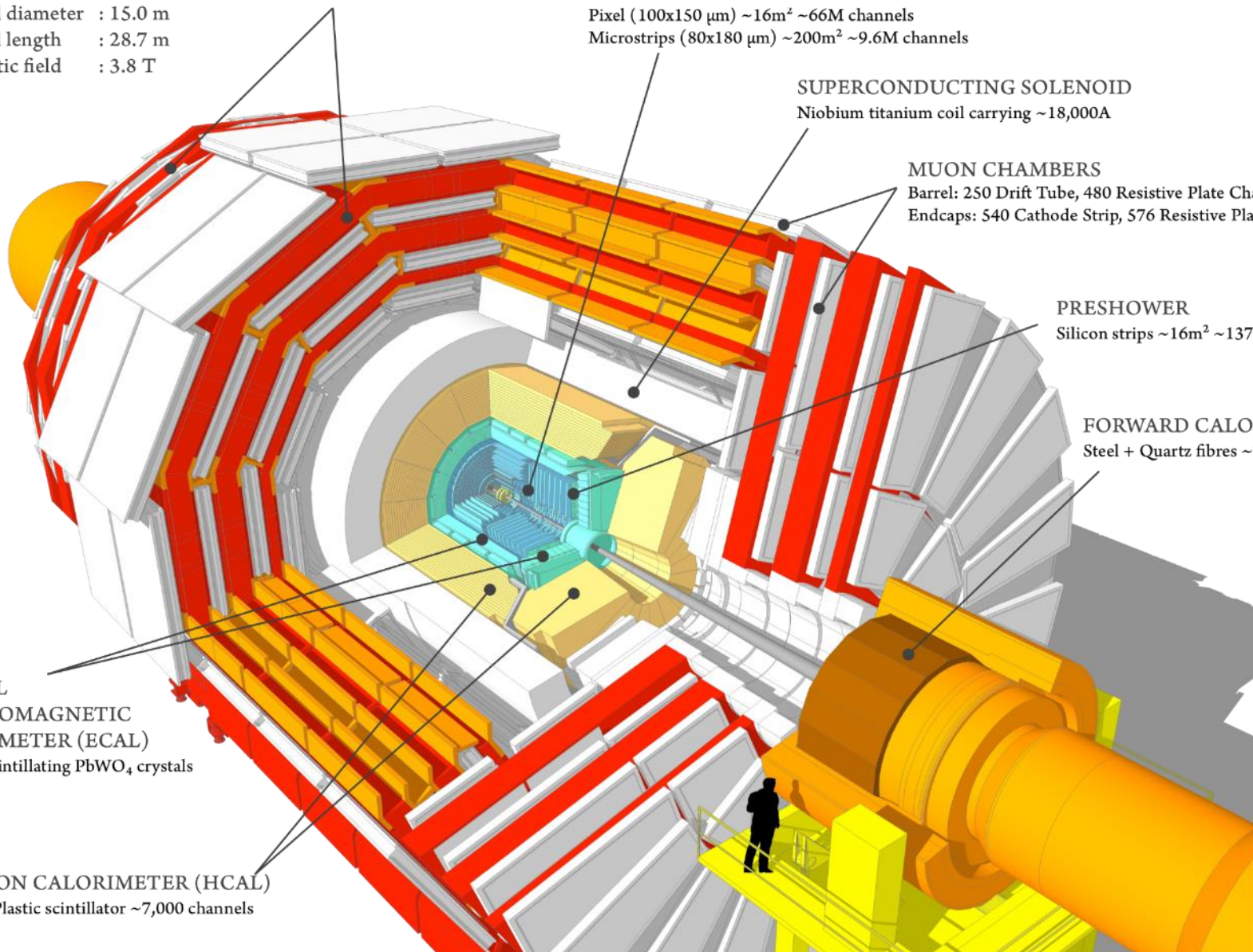
MUON CHAMBERS
Barrel: 250 Drift Tube, 480 Resistive Plate Chambers
Endcaps: 540 Cathode Strip, 576 Resistive Plate Chambers

PRESHOWER
Silicon strips $\sim 16\text{m}^2 \sim 137,000$ channels

FORWARD CALORIMETER
Steel + Quartz fibres $\sim 2,000$ Channels

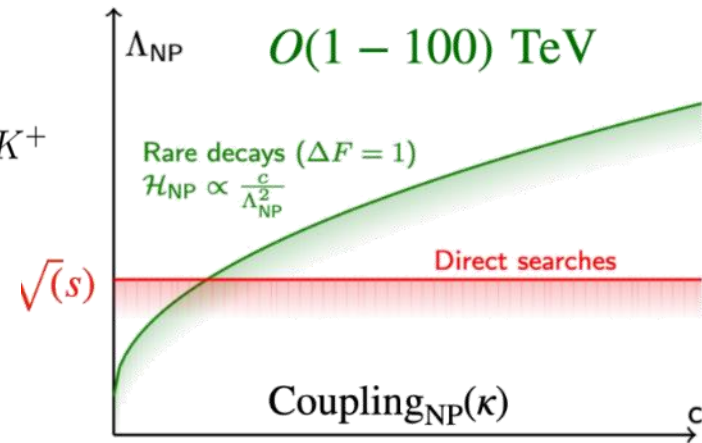
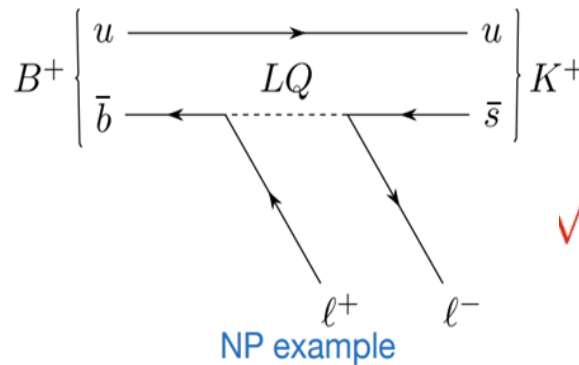
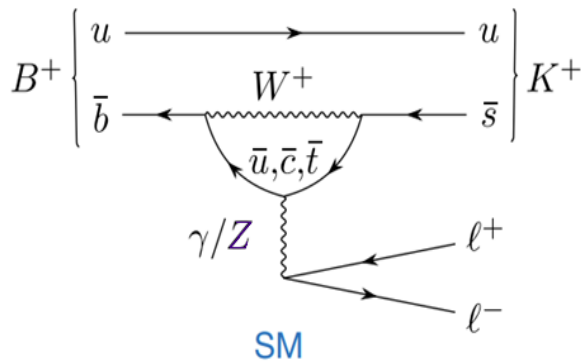
CRYSTAL
ELECTROMAGNETIC
CALORIMETER (ECAL)
 $\sim 76,000$ scintillating PbWO_4 crystals

HADRON CALORIMETER (HCAL)
Brass + Plastic scintillator $\sim 7,000$ channels



$b \rightarrow sll$ as New Physics probes

- $b \rightarrow sll$ transitions are precisely predicted by Standard Model
- Processes are rare (loop level, CKM-suppressed)
 - **new interactions can be major contribution**
- New interactions can have different symmetries from the SM
- NP can modify parameters of angular distributions observed in FCNC decays $B \rightarrow h l^+ l^-$



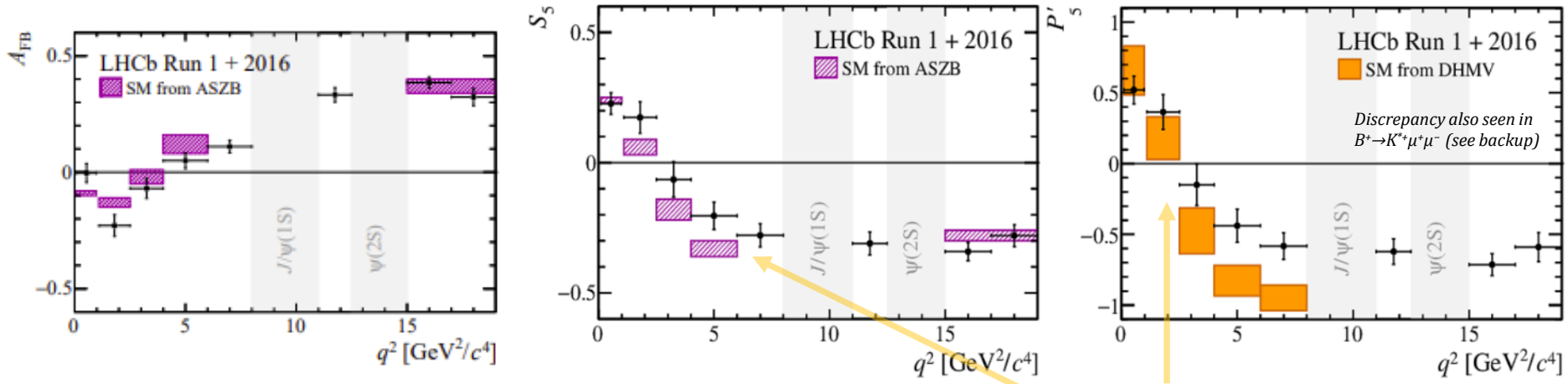
Angular analyses of $b \rightarrow s ll$ transitions

- Many recent results measuring angular parameters and differential branching fractions

$$q = m(ll)$$

$B^0 \rightarrow K^{*0} \mu^+ \mu^-$, Run-1+2016, ~ 4500 signal

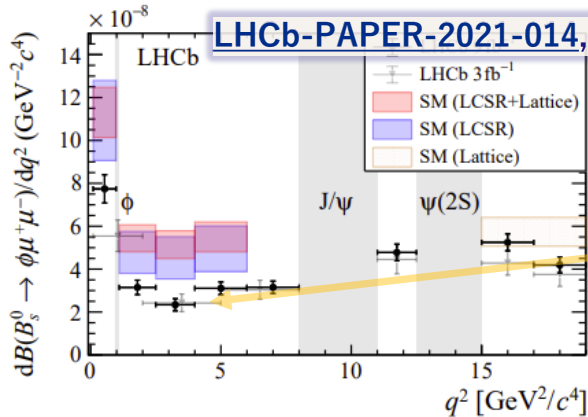
[LHCb-PAPER-2020-002, Phys.Rev.Lett.125\(2020\)011802](#)



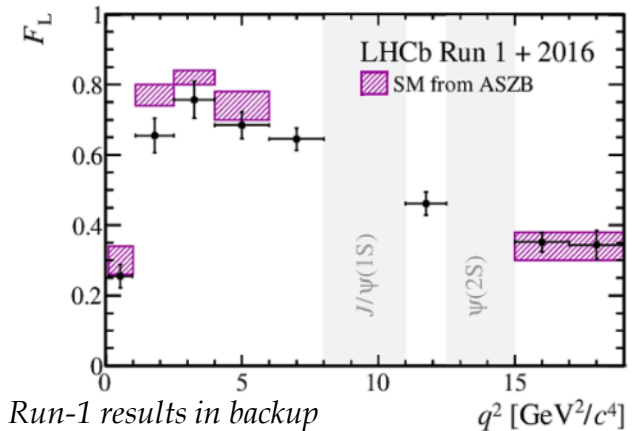
$\sim 3\sigma$ tension at low q^2 !

BF measurement of $B_s^0 \rightarrow \phi \mu^+ \mu^-$

[LHCb-PAPER-2021-014, Phys.Rev.Lett.127\(2021\)151801](#)

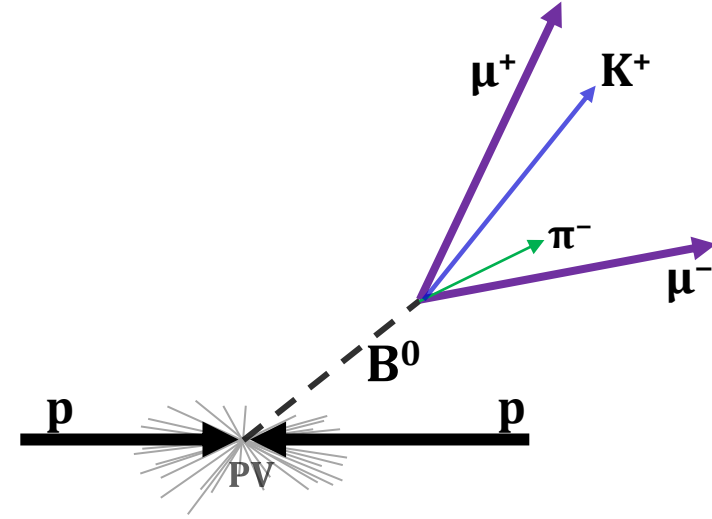
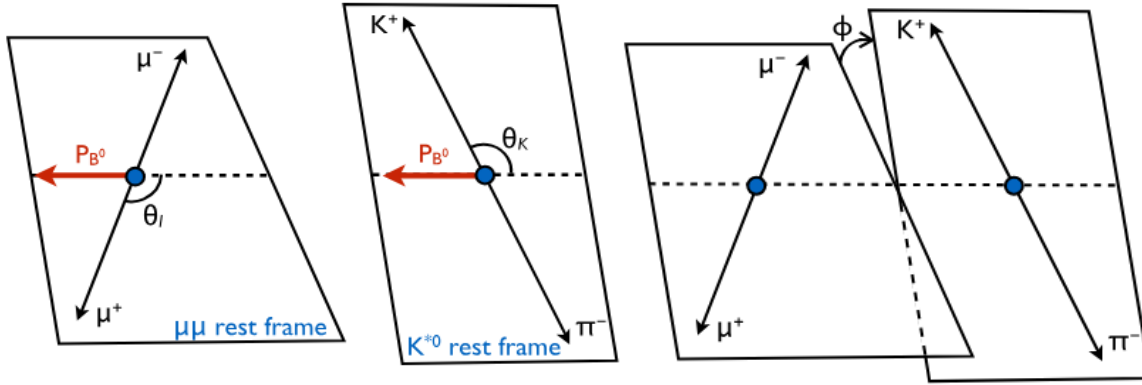


$\sim 3.6\sigma$ tension w.r.t. SM at low q^2



ATLAS and CMS Run-1 results in backup

$B^0 \rightarrow K^{*0} \mu^+ \mu^-$ Angular decay rate



$$\frac{d^4\Gamma_{P\text{-wave}}}{dq^2 d^3\Omega} = \frac{1}{d\Gamma/dq^2} \frac{d^4\Gamma}{dq^2 d\cos\theta_l d\cos\theta_K d\phi} = \frac{9}{32\pi} \left[\frac{3}{4} F_T \sin^2\theta_K + F_L \cos^2\theta_K \right. \\ \left. + \left(\frac{1}{4} F_T \sin^2\theta_K - F_L \cos^2\theta_K \right) \cos 2\theta_l + \frac{1}{2} P_1 F_T \sin^2\theta_K \sin^2\theta_l \cos 2\phi \right. \\ \left. + \sqrt{F_T F_L} \left(\frac{1}{2} P_4' \sin 2\theta_K \sin 2\theta_l \cos \phi + P_5' \sin 2\theta_K \sin \theta_l \cos \phi \right) \right. \\ \left. - \sqrt{F_T F_L} \left(P_6' \sin 2\theta_K \sin \theta_l \sin \phi - \frac{1}{2} P_8' \sin 2\theta_K \sin 2\theta_l \sin \phi \right) \right. \\ \left. + 2P_2 F_T \sin^2\theta_K \cos \theta_l - P_3 F_T \sin^2\theta_K \sin^2\theta_l \sin 2\phi \right]$$

F_L is the fraction of longitudinally polarised K^* (and $F_T = 1 - F_L$)

P_2 is related to A_{FB} , the muon forward backward asymmetry

$P_i^{(0)}$ is the base of optimised observables

The $K\pi$ system can also be in S-wave \rightarrow

$$\frac{d^4\Gamma_{total}}{dq^2 d^3\Omega} = (1 - F_S) \frac{d^4\Gamma_{P\text{-wave}}}{dq^2 d^3\Omega} + \frac{d^4\Gamma_{S/SP\text{-wave}}}{dq^2 d^3\Omega}$$

Data set and $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ selection

Run 2 dataset (2016-2018), corresponding to 140 fb^{-1}

Trigger requires 2 muons + a track

$B^0 \rightarrow J/\psi K^{*0}$ and $B^0 \rightarrow \psi(2S) K^{*0}$ decays with the same final-state particles $K^\pm \pi^\mp \mu^+ \mu^-$ are used for the control and validation

BDT to distinguish signal from background

- trained on signal MC and background from data sidebands
- different training per year of data taking, k-folding to avoid overtraining
- input features: decay-vertex quality and displacement, isolation, mass of $K\pi$ system
- chosen working point optimises signal significance

$B^0 \rightarrow J/\psi K^*$ and $B^0 \rightarrow \psi(2S) K^*$ leak in the nearby signal region, mainly due to unreconstructed photon

- combined cut on $m(K\pi\mu\mu)$ and q^2

Specific backgrounds

$B^+ \rightarrow K^+ \mu\mu$ (plus combinatorial track)

- ▶ additional veto on mass of two $h^+ \mu\mu$ systems

$B^+ \rightarrow K^+ \psi(2S)$, with $\psi(2S) \rightarrow J/\psi \pi\pi$ (a π track is lost)

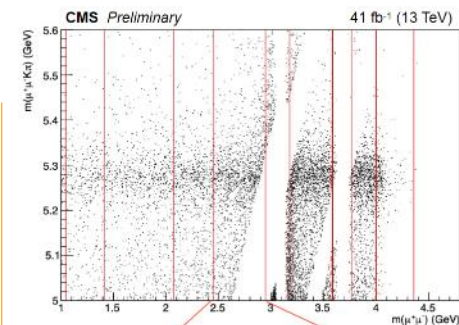
- ▶ only affects J/ψ control region
- ▶ combination of cuts on intermediate masses

$B_s \rightarrow \phi(\rightarrow KK) \mu\mu$

- ▶ veto at preselection level on KK mass hypothesis
- ▶ residual contribution negligible wrt signal ($< 1\%$)

$B_s \rightarrow KK\mu\mu$ contribution (4%) treated as combinatorial bkg

Negligible contribution from $B_s \rightarrow K^* \mu\mu$ ($< 1\%$), no evidence of $\Lambda_b \rightarrow p K^+ \mu^-$



4D UML simultaneous fit in 6 q^2 bins

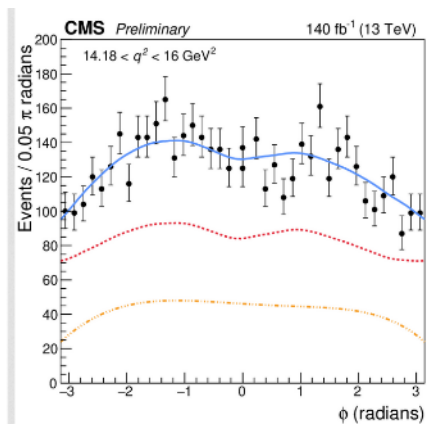
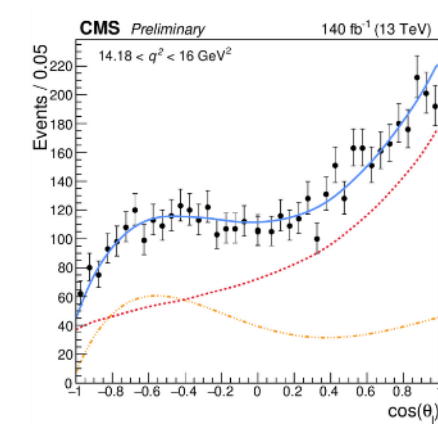
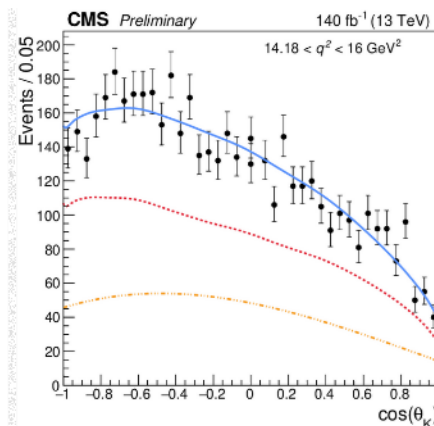
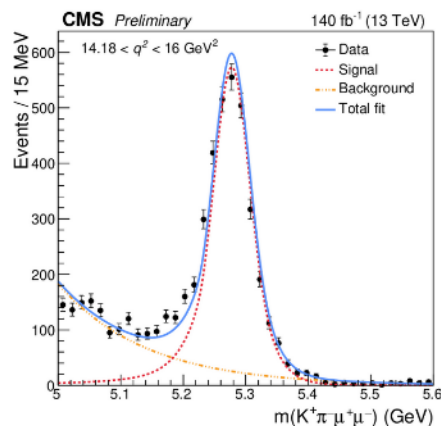
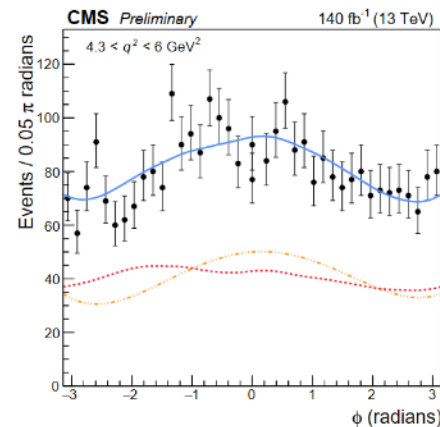
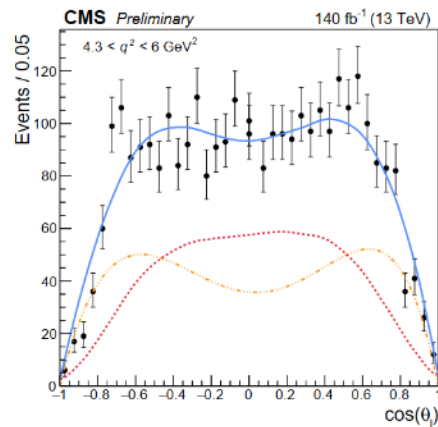
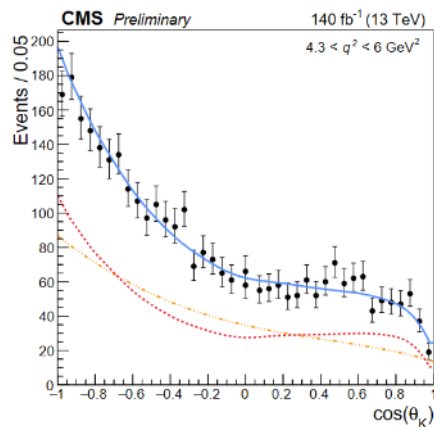
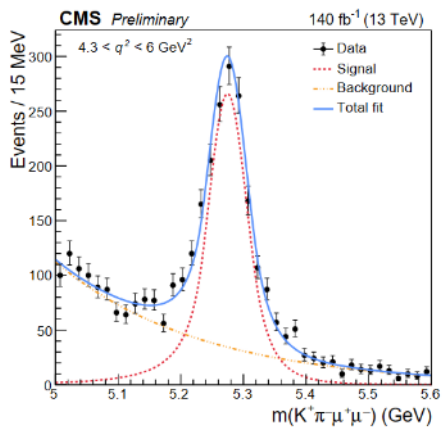
3 angles + mass

3 data-taking years

Fit to $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ sample additionally accounts for:

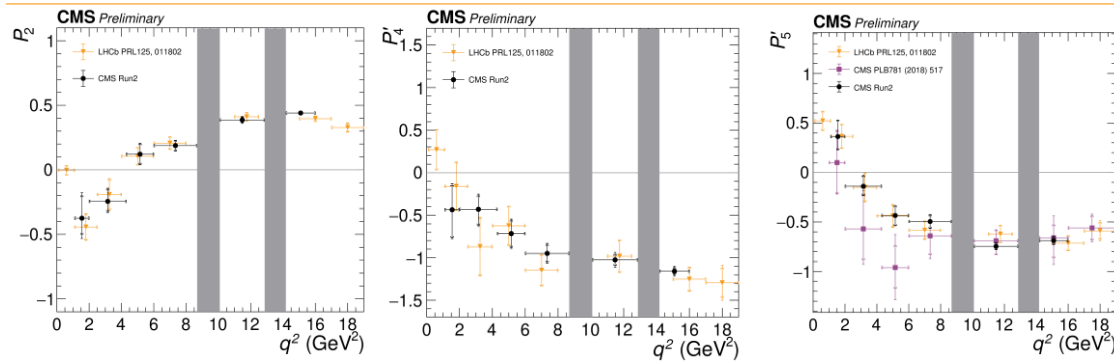
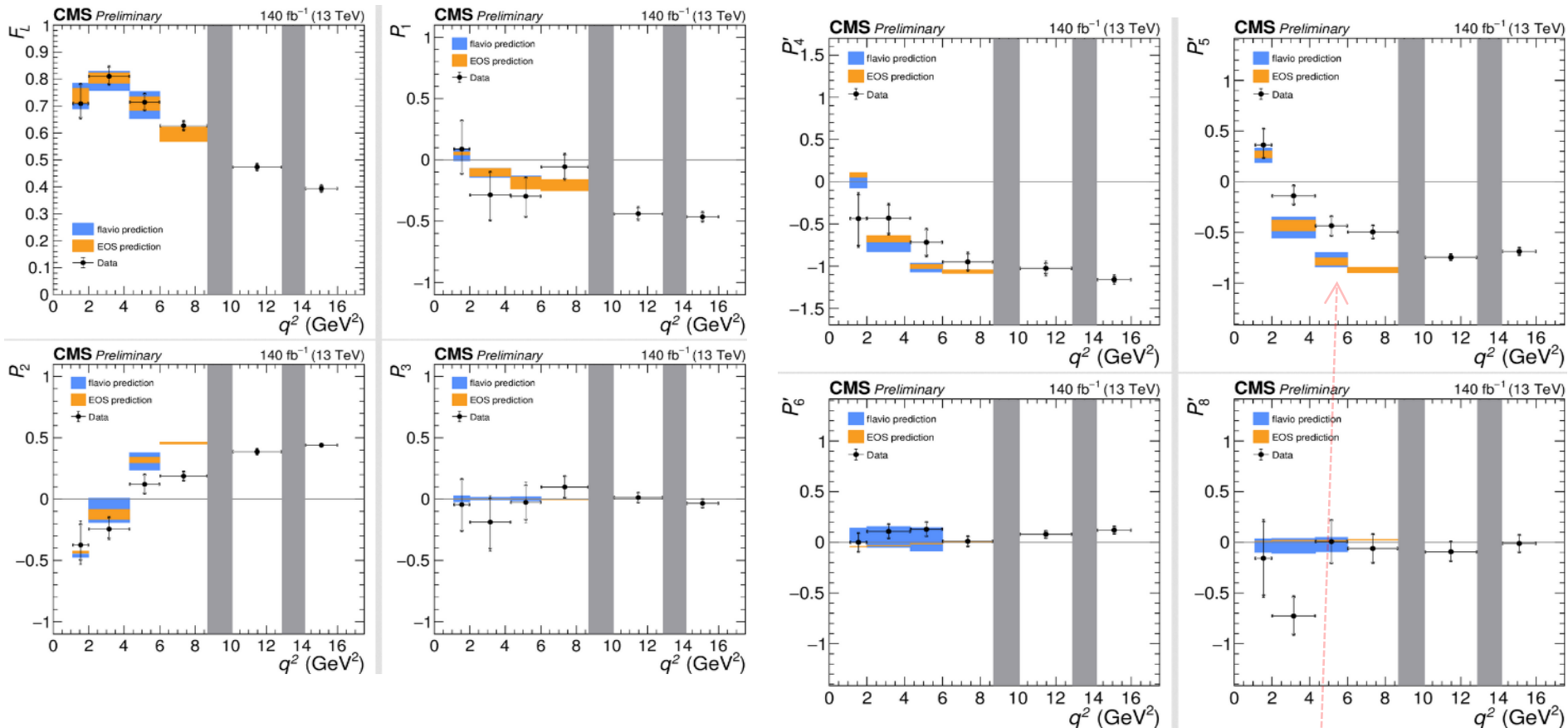
- the possibility that K and π are swapped
- Efficiency as a function of angular variables
- Background angular shape (from sidebands)
- Physical boundary in the angular observables space

[CMS-PAS-BPH-21-002](#), arXiv soon!



$B^0 \rightarrow K^{*0} \mu^+ \mu^-$ Results

$q = m(ll)$

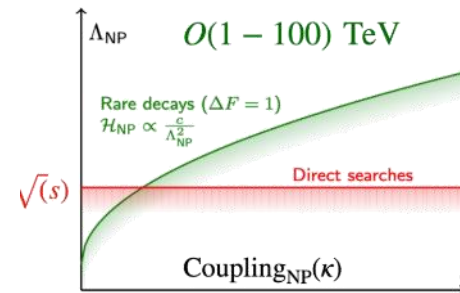
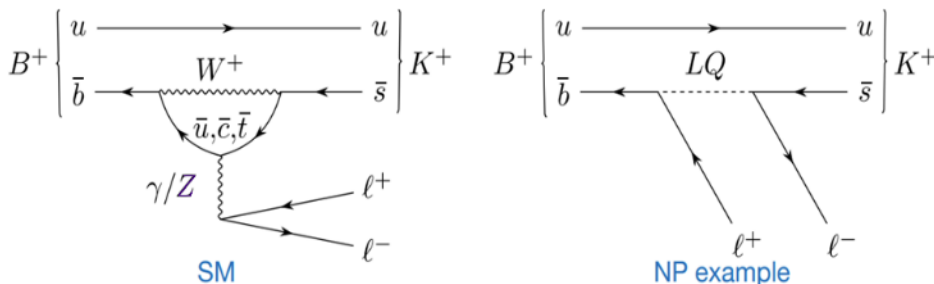


Good agreement with, and similar precision to the previous LHCb results
 P_5 discrepancy remains...

Lepton Flavour Universality Violation

- SM has identical couplings of charged leptons to W and Z bosons
 - ➔ **similar behavior of e, μ, τ** (some differences due to different mass)
- Observation of a significant **LFUV** would immediately point to **New Physics** contribution
- **LFU** tests in $b \rightarrow s ll$ are theoretically-clean observables

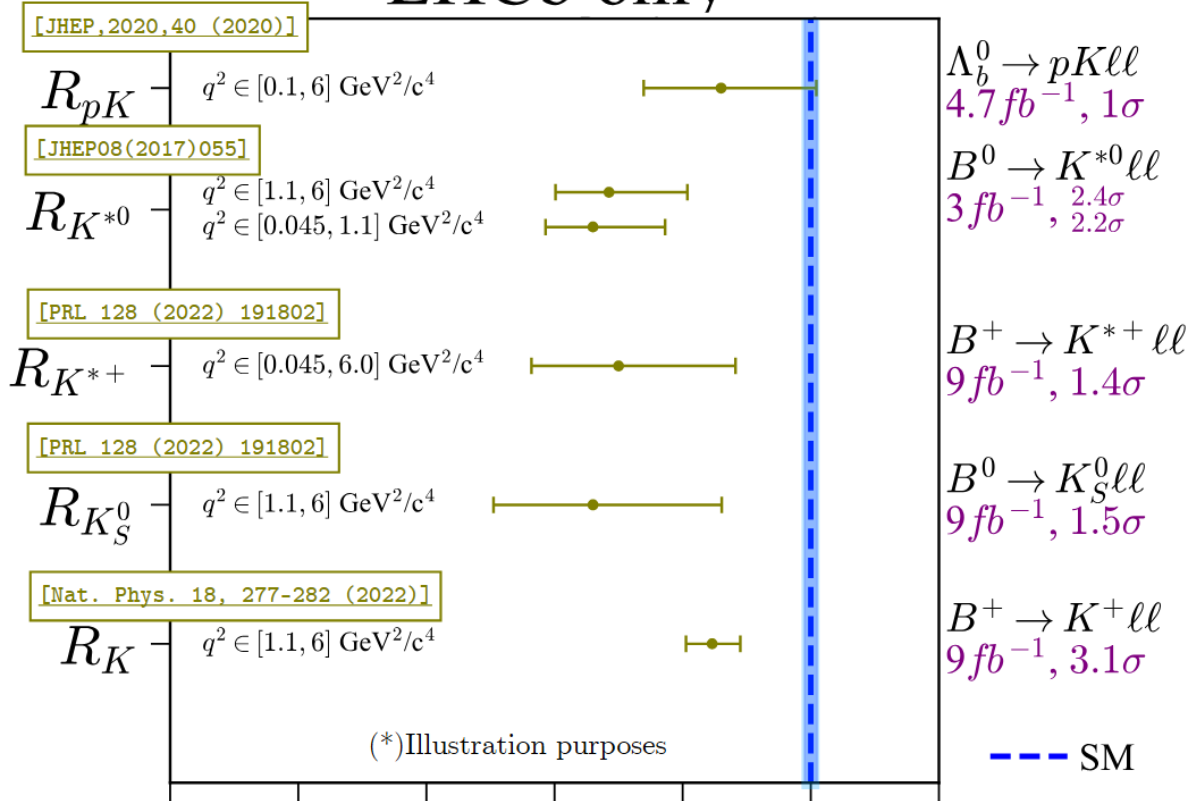
$$R(K)(q^2)[q_{\min}^2, q_{\max}^2] = \frac{\mathcal{B}(B^+ \rightarrow K^+ \mu^+ \mu^-)[q_{\min}^2, q_{\max}^2]}{\mathcal{B}(B^+ \rightarrow K^+ e^+ e^-)[q_{\min}^2, q_{\max}^2]} \quad q=m(ll)$$



LFUV searches: previous results

$$q=m(ll)$$

LHCb only



A set of previous results (with precision dominated by LHCb measurements) have indicated a discrepancy from SM prediction in R_x ratios (until 2023)

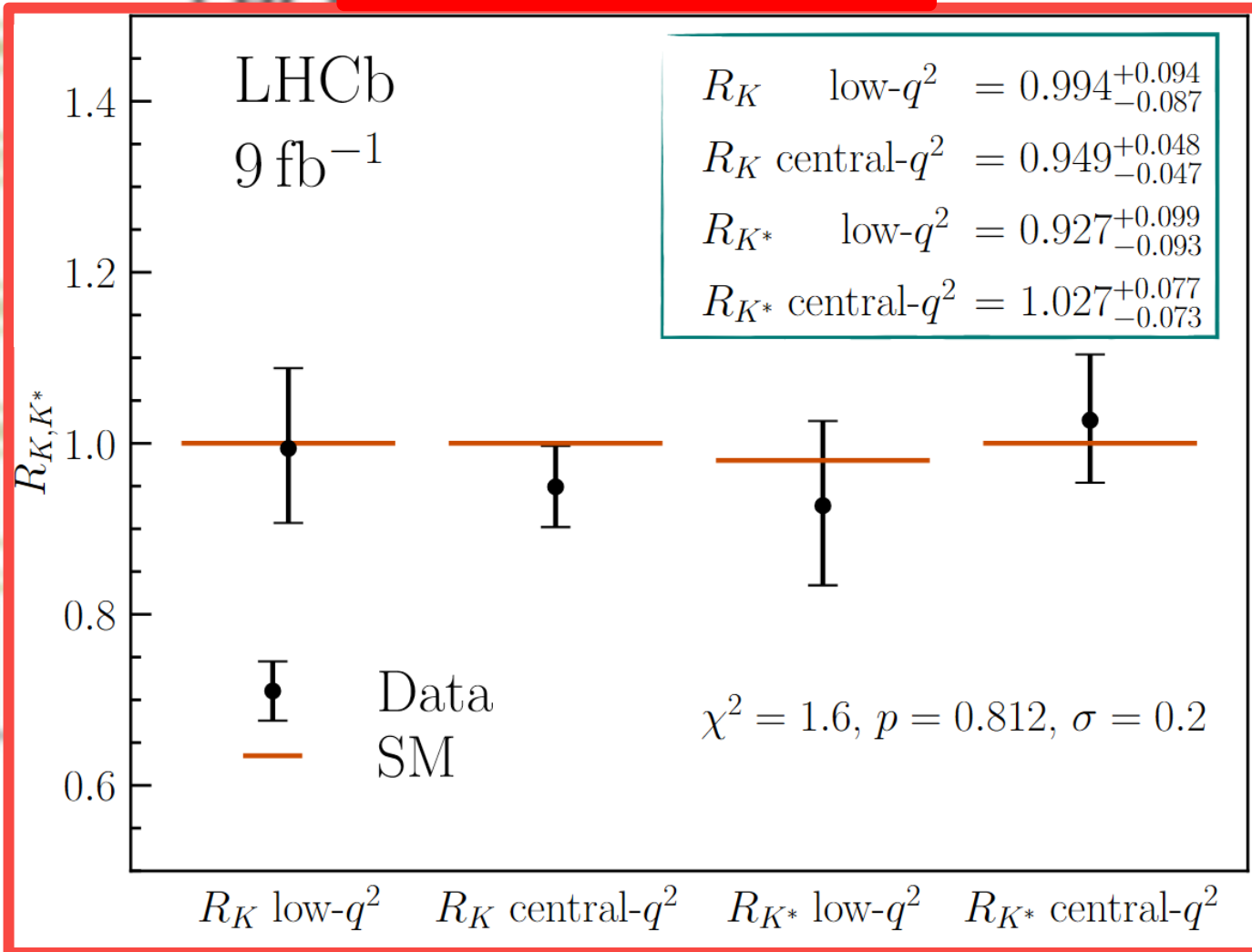
All the ratios are lower than prediction?

$$R_X = \frac{\mathcal{B}(b \rightarrow s\mu^+\mu^-)}{\mathcal{B}(b \rightarrow se^+e^-)}$$

(τ) Measurements from Belle not shown (larger statistical uncertainties)

LFUV searches: previous results

However, in 2023:

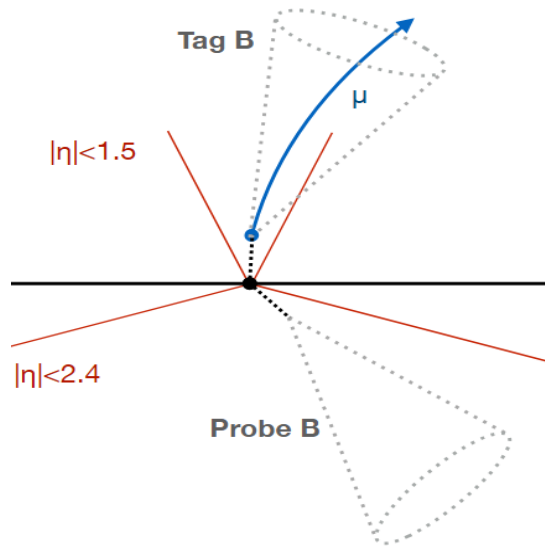


(*) Measurements from Belle and others (large statistical uncertainties)

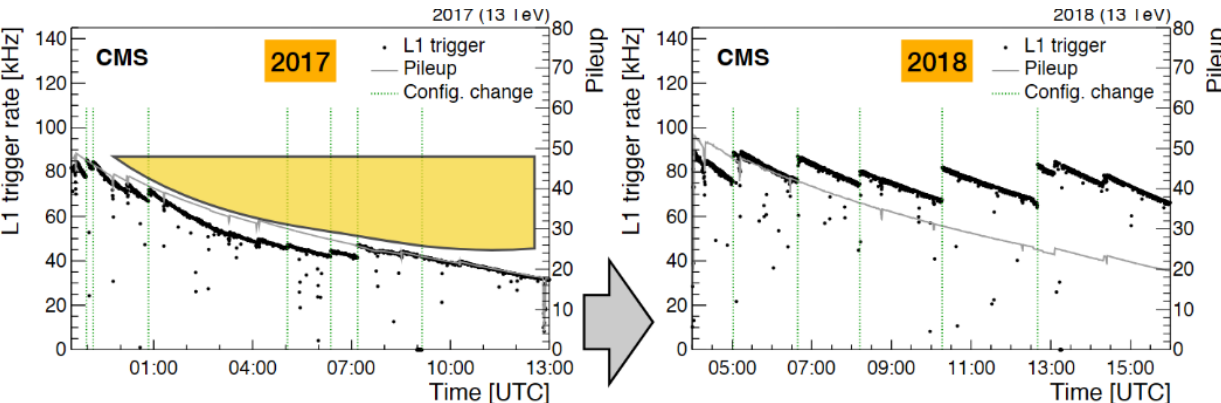
Data sample

➤ Data Parking stands for “data that will be reconstructed later”

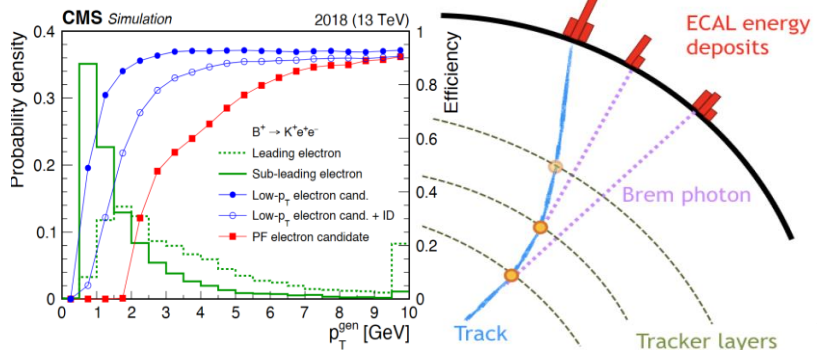
- In 2018, CMS implemented “B-parking”
- triggering on low momentum electrons is very hard
→ leverage the $b\bar{b}$ pair production, trigger on one B and investigate the other
- Triggering on displaced muon from SL b decays
- Constantly utilize full L1 bandwidth



Using a family of L1 seeds with decreasing thresholds as lumi decreases, to “fill the gap to 100 kHz”



- Require muon displacement at HLT
- Resulting purity ~75-80% (events with b hadron)
- Dedicated low- p_T electron reconstruction



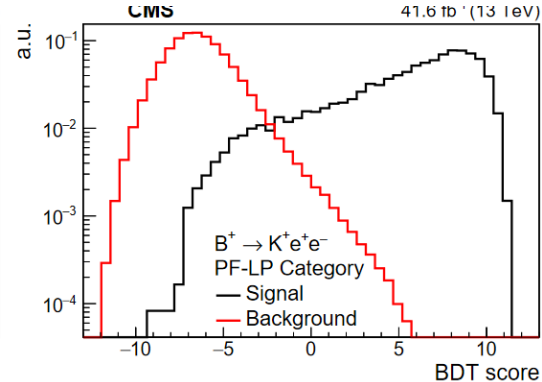
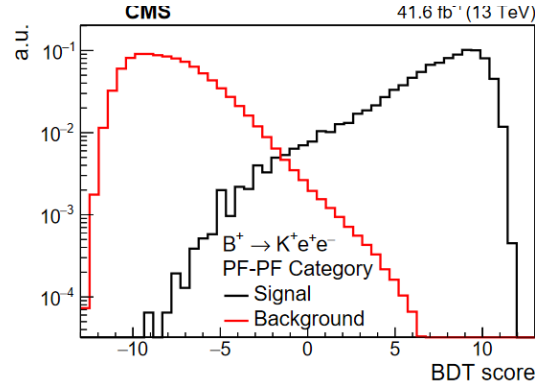
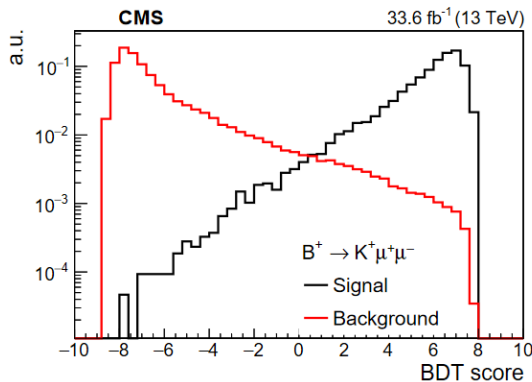
Measurement strategy

CMS-BPH-22-005,
Rep. Prog. Phys. 87 (2024) 077802

- Measuring double ratio to reduce systematics

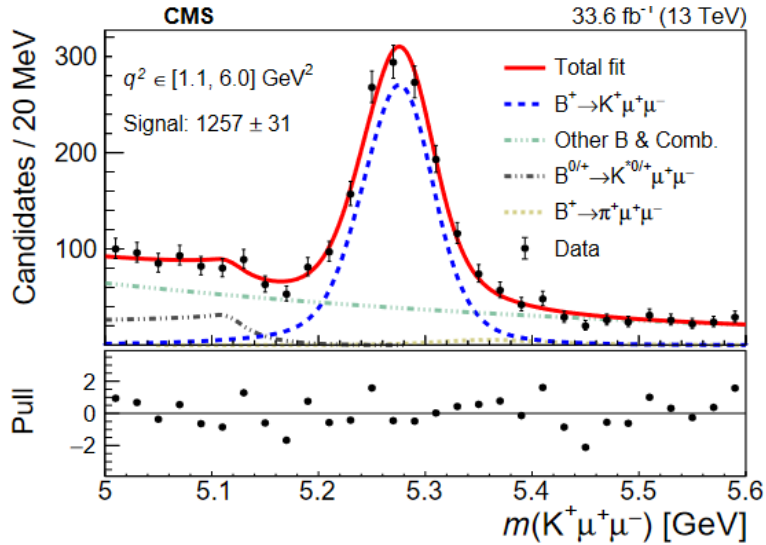
$$R_{K,K^*}(q_a^2, q_b^2) = \frac{\int_{q_a^2}^{q_b^2} \frac{d\Gamma(B^{(+,0)} \rightarrow K^{(+,*0)} \mu^+ \mu^-)}{dq^2} dq^2}{\int_{q_a^2}^{q_b^2} \frac{d\Gamma(B^{(+,0)} \rightarrow K^{(+,*0)} e^+ e^-)}{dq^2} dq^2} = \frac{\mathcal{B}(B^+ \rightarrow K^+ \mu^+ \mu^-)[q_{\min}^2, q_{\max}^2]}{\mathcal{B}(B^+ \rightarrow J/\psi(\mu^+ \mu^-) K^+)} \bigg/ \frac{\mathcal{B}(B^+ \rightarrow K^+ e^+ e^-)[q_{\min}^2, q_{\max}^2]}{\mathcal{B}(B^+ \rightarrow J/\psi(e^+ e^-) K^+)}$$

- Muon channel on tag side, electron channel on probe side
- Electrons from two algos: PF and LP → PF-PF and PFLP categories
- BDT discriminator(s) for background rejection
 - Features: vertex displacement, fit probability, pointing angle
muon isolation, kaon p_T , B p_T
electron p_T , isolation, K IP w.r.t ee vertex, ID

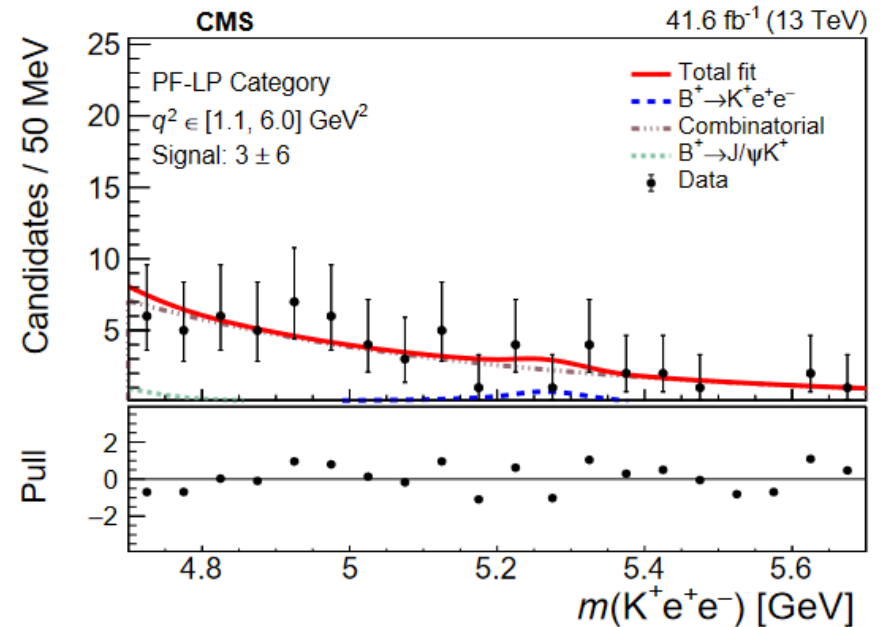
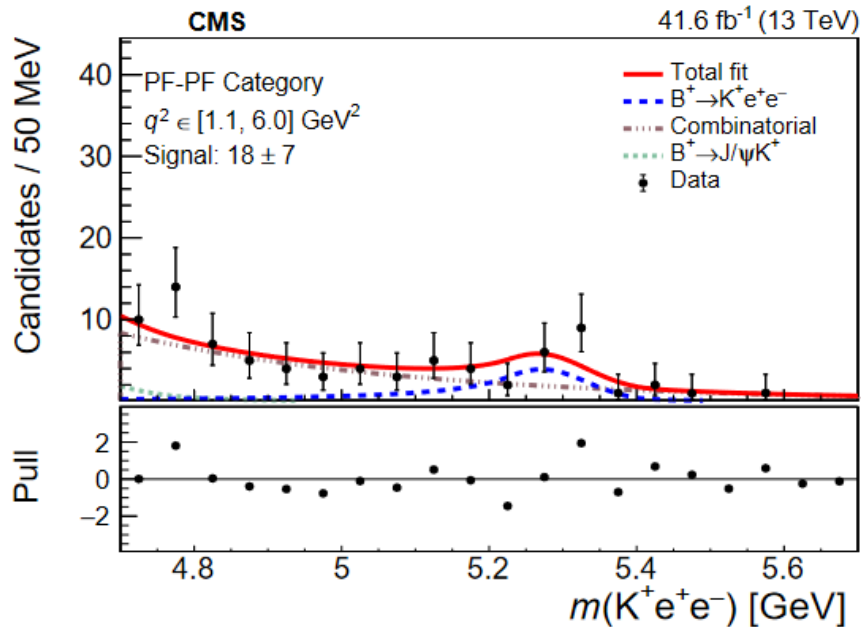


- Dedicated vetoes to kill backgrounds from charm and charmonia

Observed signals



At just 20 events in the signal channel, the precision of $R(K)$ measurement is quite low



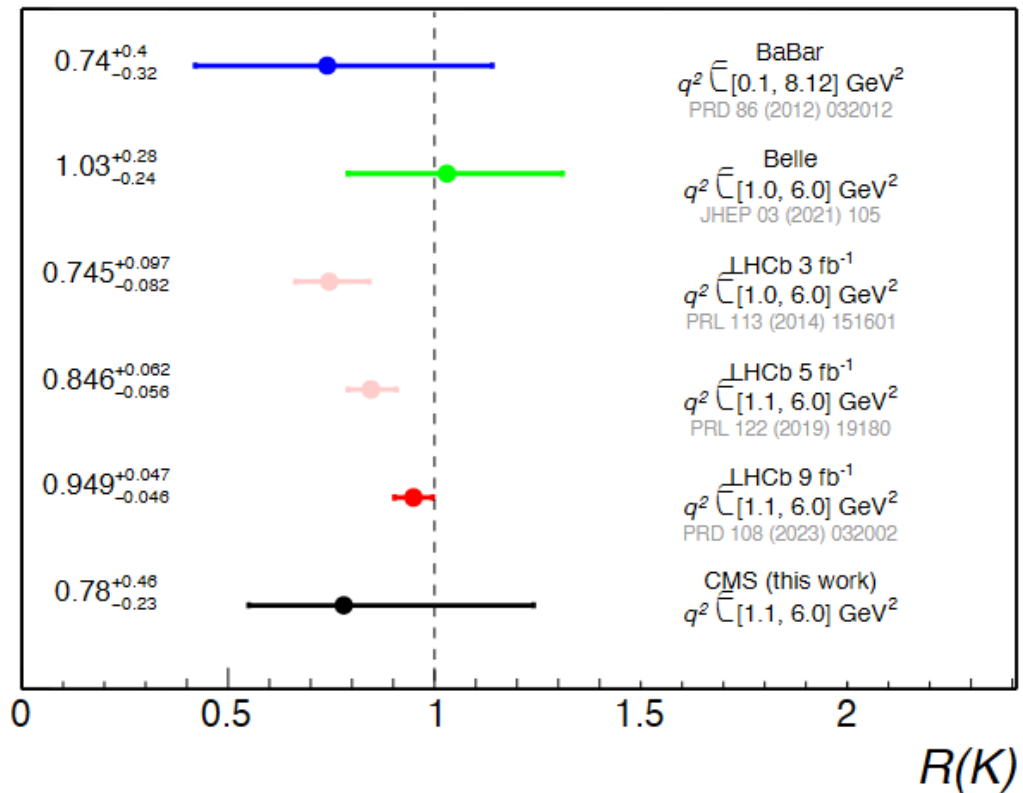
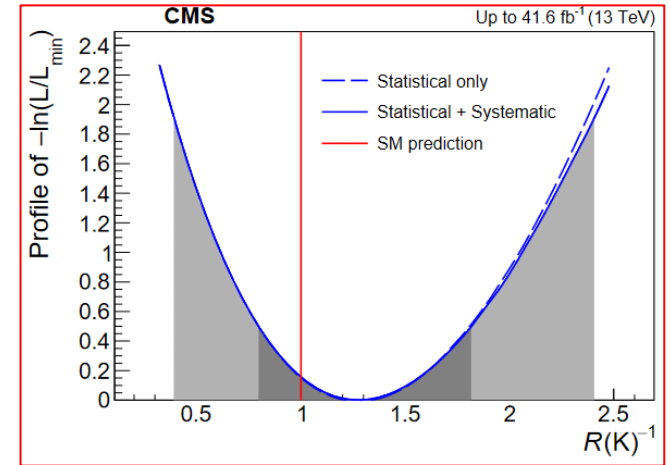
R(K) result

$$R(K) = 0.78^{+0.46}_{-0.23} (\text{stat})^{+0.09}_{-0.05} (\text{syst}) = 0.78^{+0.47}_{-0.23}$$

The resulting value is consistent both with SM and LHCb, but the precision is not competitive

The main reason is the difficulty with low- p_T electrons, especially at the trigger level, which is why the trigger for this analysis searches for *the other b* in an event

CMS has learned a lot and huge improvement is expected in Run-3, stay tuned!



Summary

- ❖ Many new results on rare $b \rightarrow sll$ decays, sensitive to NP, are obtained in the last few years
- ❖ Some **$\sim 3\text{-}4\sigma$ tensions** w.r.t. SM have been observed in B meson decays
- ❖ CMS made an angular analysis of the $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ decay with full Run-2 data
 - ❖ The results agree with and are as precise as earlier LHCb results
 - ❖ Confirming the tension between experiment and SM in P'_5 parameter
- ❖ $R(K)$ measurement is very challenging at CMS due to low- p_T electrons
 - ❖ The result agrees (with large uncertainties) with LHCb and LFU=SM
 - ❖ In the meanwhile, refreshed LHCb results from last year also agree with LFU=SM...
- ❖ Run-3 triggers will allow to make significantly more precise measurements

<http://cms-results.web.cern.ch/cms-results/public-results/publications/BPH/>

<https://cms-results.web.cern.ch/cms-results/public-results/preliminary-results/BPH/index.html>



Thank you !

BACKUP

Search for $D^0 \rightarrow \mu\mu$

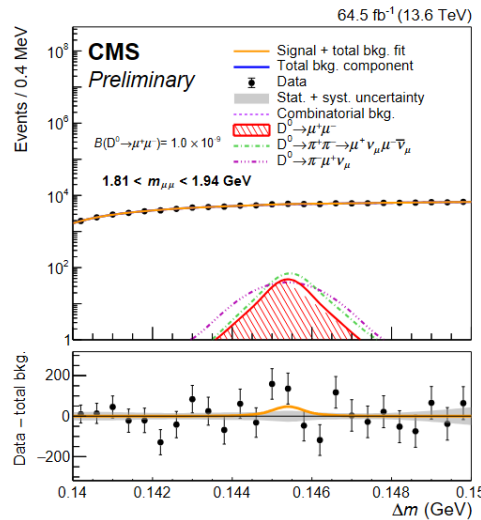
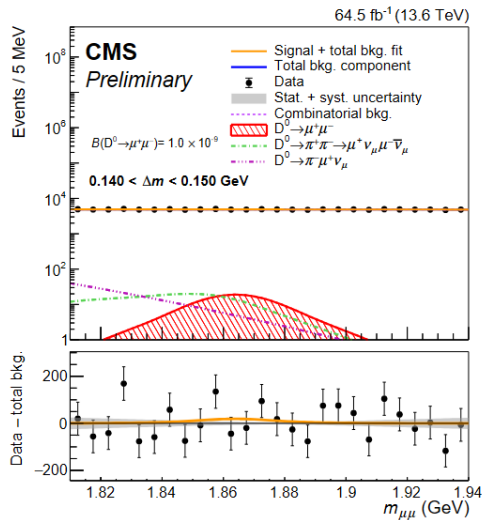
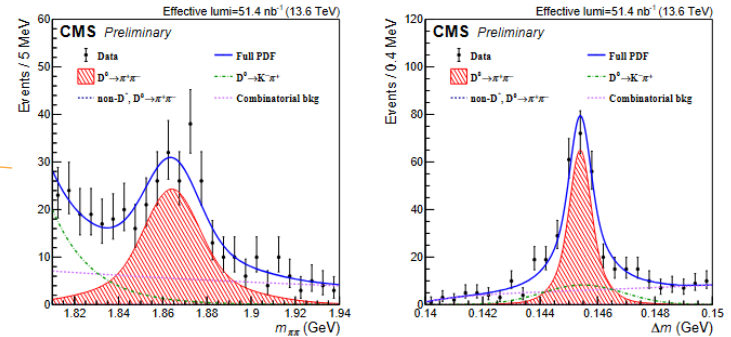
With many studies of $b \rightarrow sll$ transitions, even more rare $c \rightarrow ull$ ones are often forgotten/neglected

CMS searches for the $D^0 \rightarrow \mu\mu$ decay via $D^{*+} \rightarrow D^0 \pi^+$ to suppress bkg

One of the first analyses to:

- use Run-3 13.6 TeV data! (2022+2023 in this case)
- use new inclusive dimuon triggers

Normalization relative to $D^0 \rightarrow \pi^+ \pi^-$

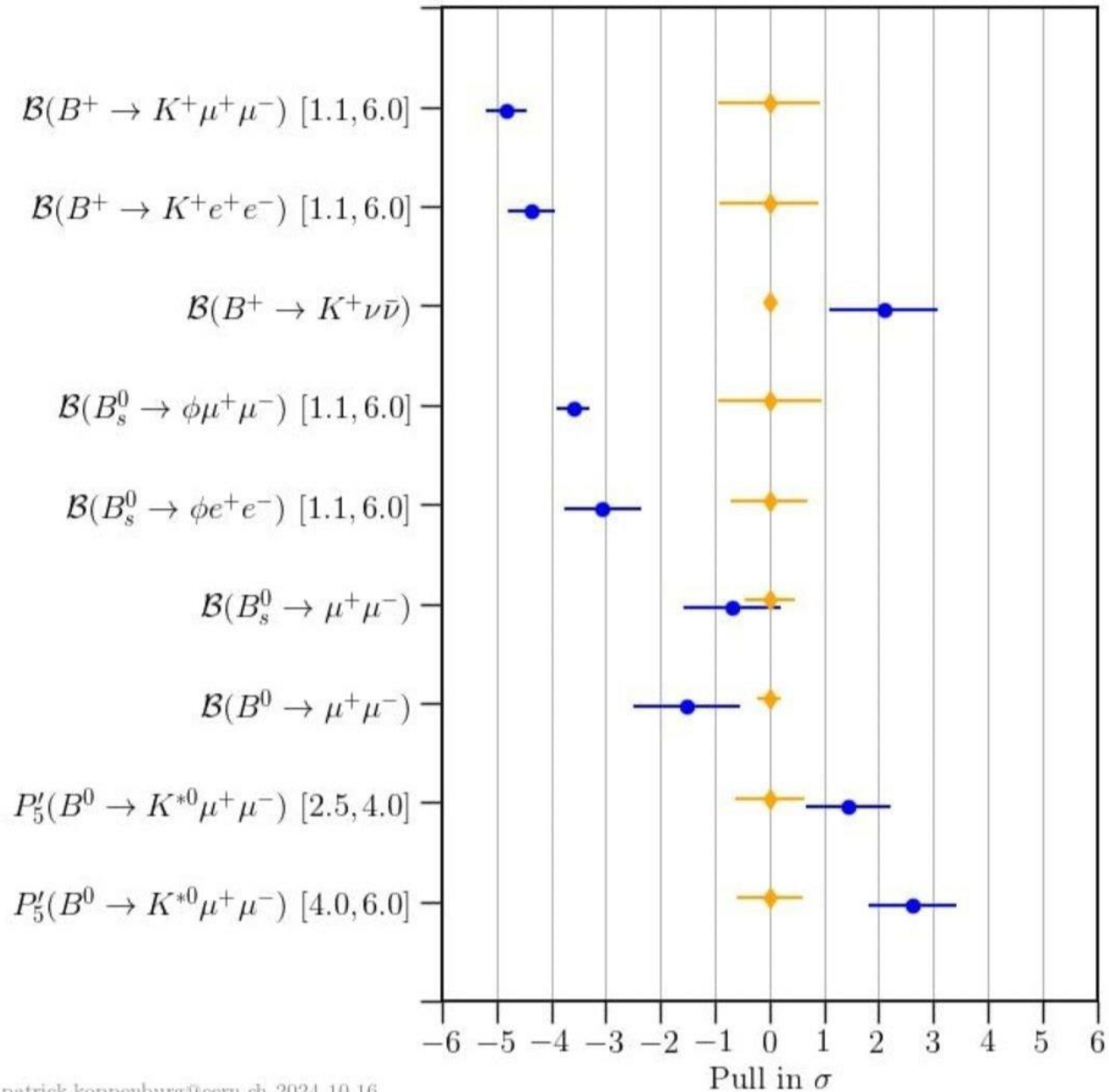


No signal observed, UL set at

$$B(D^0 \rightarrow \mu^+ \mu^-) < 2.6 \times 10^{-9} \text{ at 95\% CL}$$

Best limit to date, 35% improvement over [LHCb](#)
SM prediction $\sim 3 \cdot 10^{-13}$

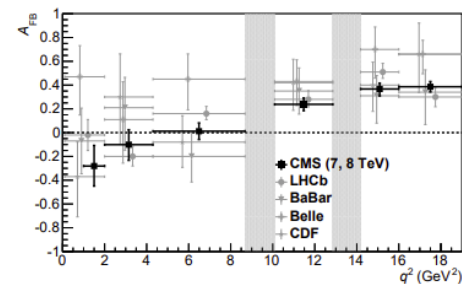
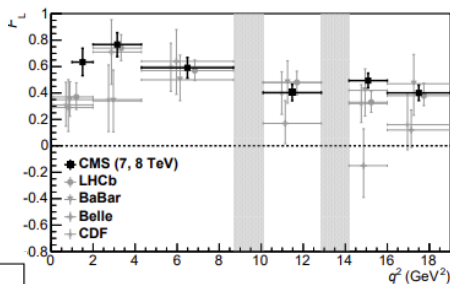
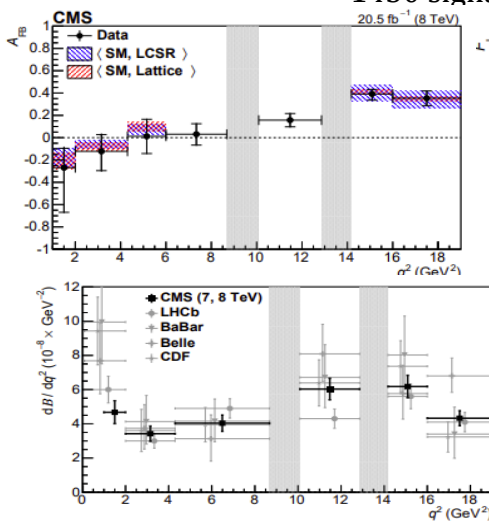
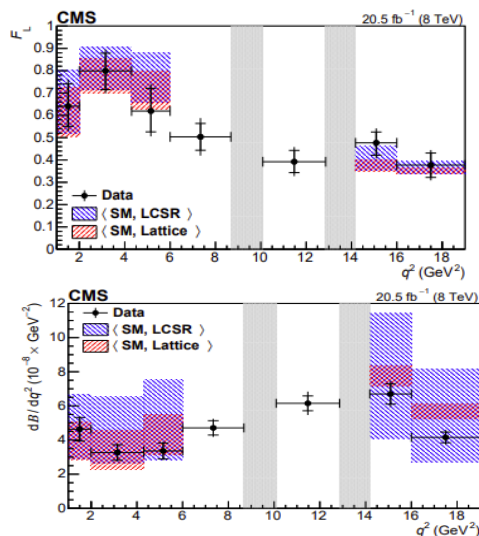
“anomalies” in FCNC transitions



$B^0 \rightarrow K^{*0} \mu^+ \mu^-$, CMS and ATLAS Run-1

~1450 signal, 7 q^2 bins

CMS-BPH-13-010, Phys. Lett. B 753 (2016) 424

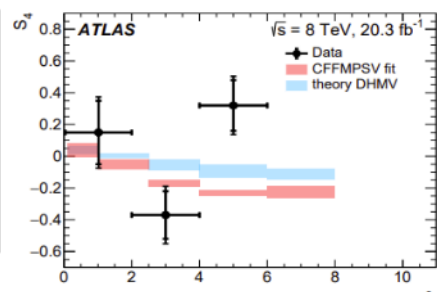
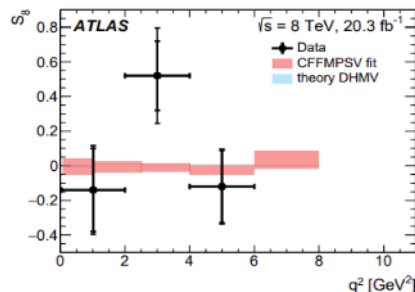
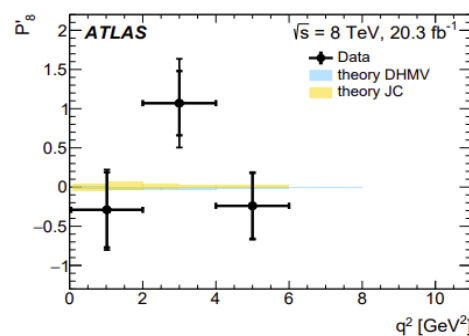
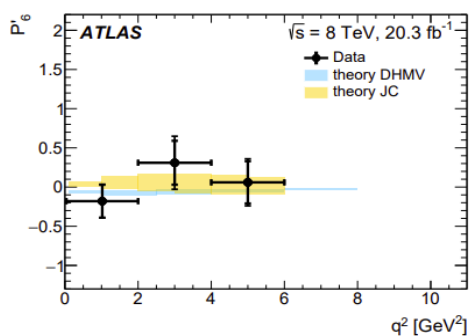
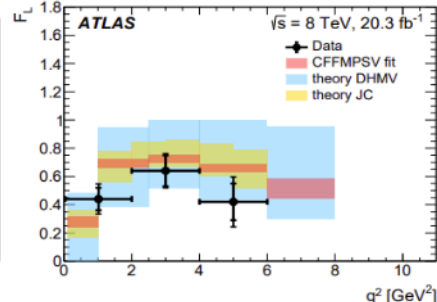
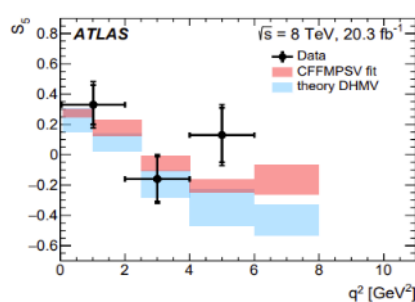
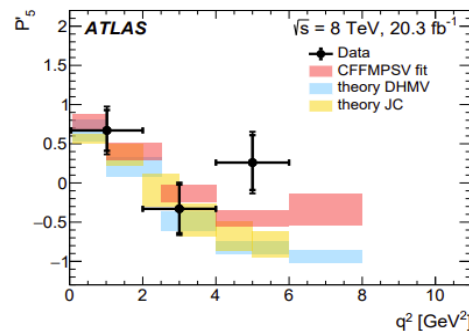
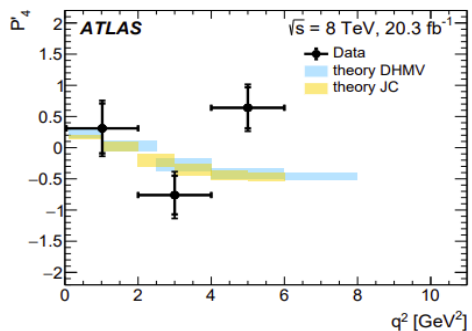


Consistent with SM

More details in the parallel talk by Samar N. yesterday

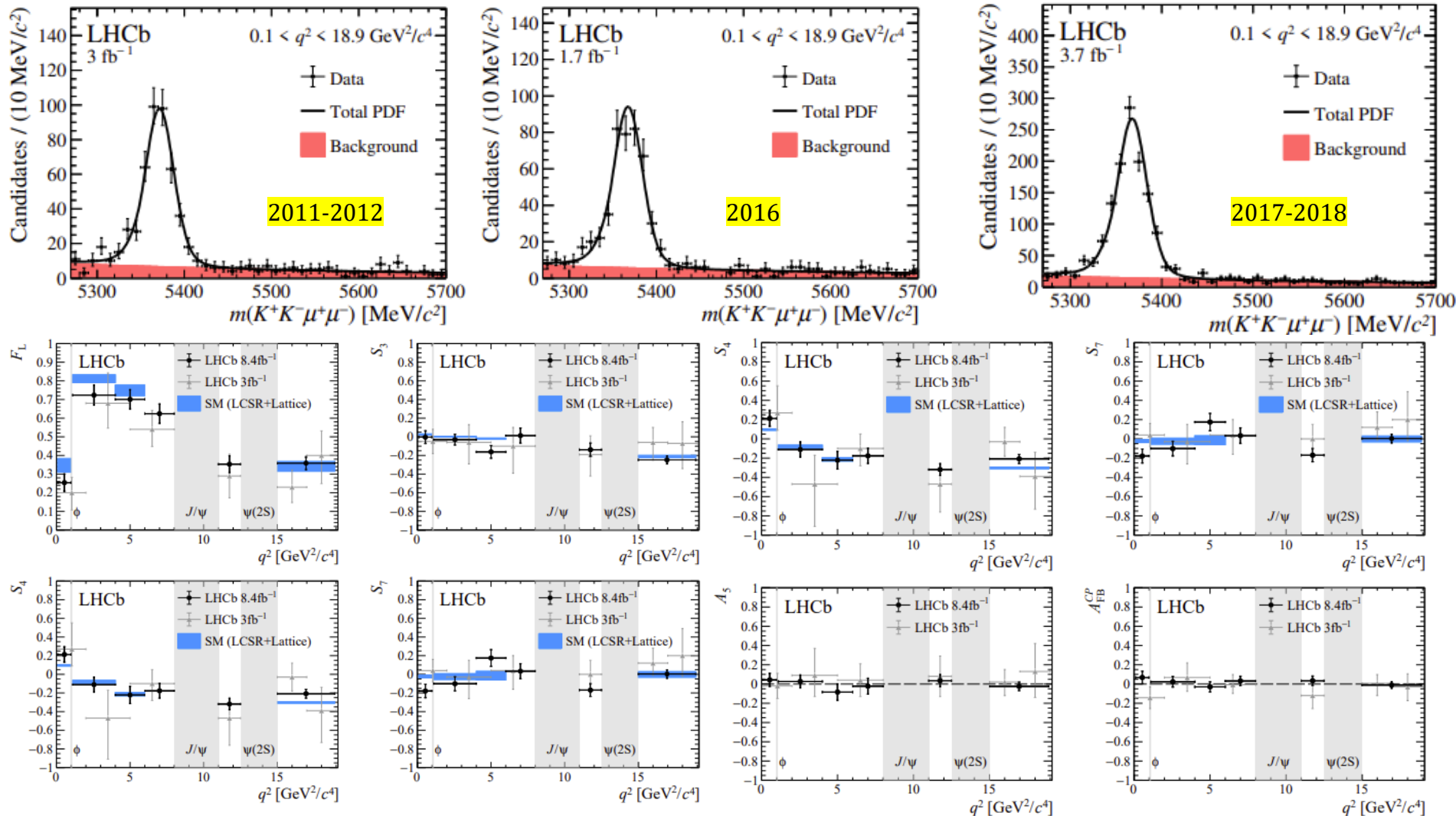
ATLAS-BPHY-13-02, JHEP10(2018)047

~350 signal, 3 q^2 bins



Angular analyses of $b \rightarrow s l l$ transitions

$B_s^0 \rightarrow \phi \mu^+ \mu^-$ angular, Run-1+Run-2 (no 2015), ~ 2000 signal



Compatible with SM

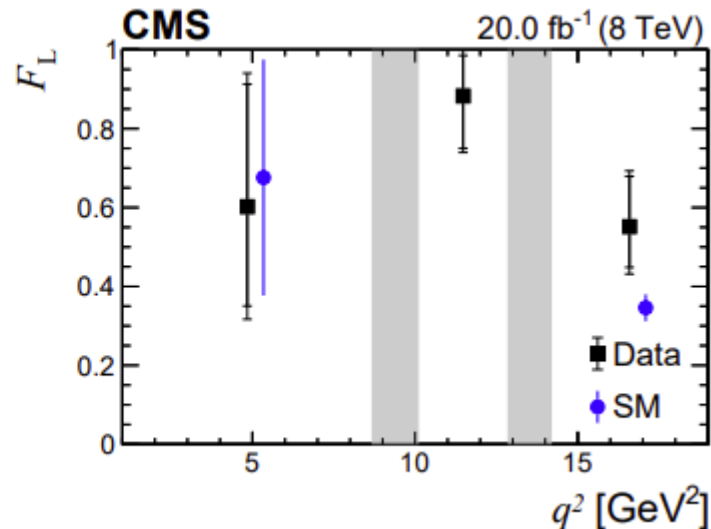
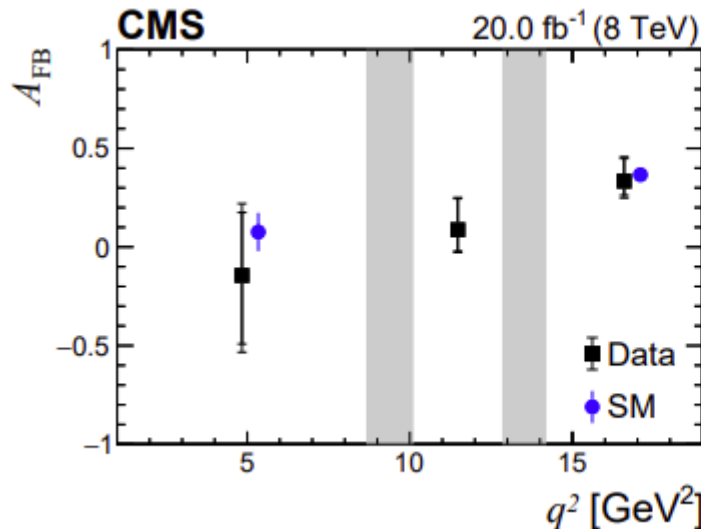
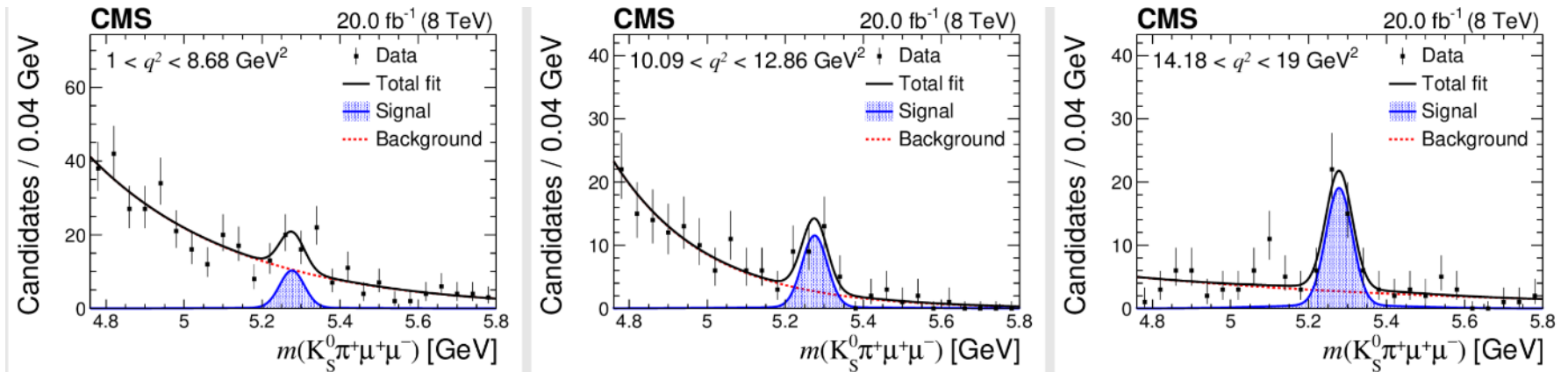
*CMS Run-2 analysis
in development*

Angular analyses of $b \rightarrow sll$ transitions

$B^+ \rightarrow K^{*+} \mu^+ \mu^-$, ($K^{*+} \rightarrow K_S^0 \pi^+$) CMS Run-1, ~ 90 signal

CMS-BPH-15-009, JHEP04(2021)124

Lower statistics compared to K^{*0} channel because of K_S^0



$$\frac{1}{d\Gamma/dq^2 d\cos\theta_l d\cos\theta_K d\phi} \frac{d^4\Gamma}{dq^2 d\cos\theta_l d\cos\theta_K d\phi} = (1 - F_S)\Gamma_V + \frac{3}{16\pi} (F_S \sin^2\theta_l + A_S \sin^2\theta_l \cos\theta_K + A_S^4 \sin\theta_K \sin 2\theta_l \cos\phi + A_S^5 \sin\theta_K \sin\theta_l \cos\phi + A_S^7 \sin\theta_K \sin\theta_l \sin\phi + A_S^8 \sin\theta_K \sin 2\theta_l \sin\phi)$$

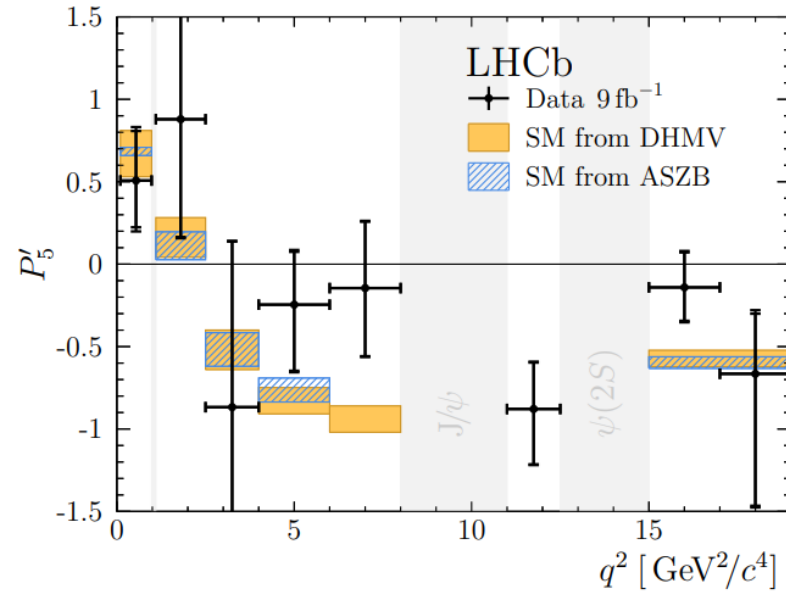
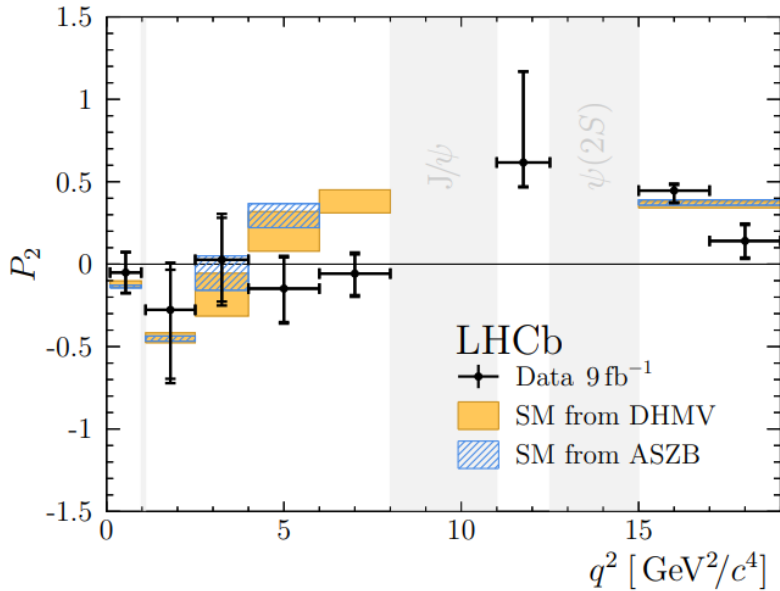
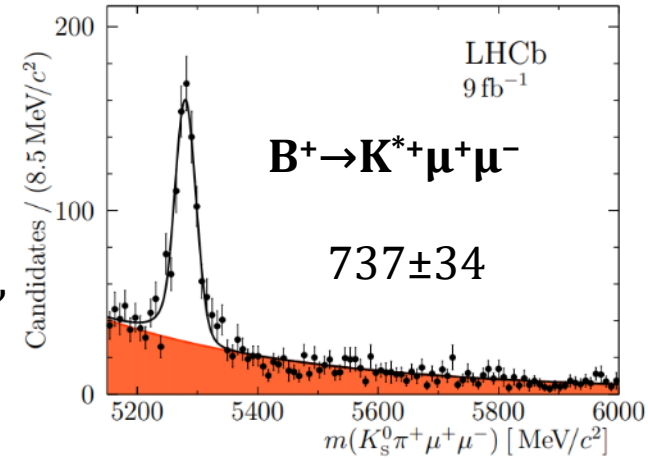
S-wave

Compatible with SM

Angular analyses of $b \rightarrow sll$ transitions

$B^+ \rightarrow K^{*+} \mu^+ \mu^-$, ($K^{*+} \rightarrow K_S^0 \pi^+$) LHCb Run-1 + Run-2, ~ 90 signal
 Lower statistics compared to K^{*0} channel because of K_S^0
 Two categories based on K_S^0 decay vertex position

Angular analysis, measuring full set of optimized variables,
 $F_L, S_3, S_4, S_5, A_{FB}, S_7, S_8, S_9, P_1, P_2, P_3, P'_4, P'_5, P'_6, P'_8$
 in 5 folds of the data, due to limited stat.



3.1 σ tension w.r.t SM at low q^2 !

Angular analyses of $b \rightarrow sll$ transitions

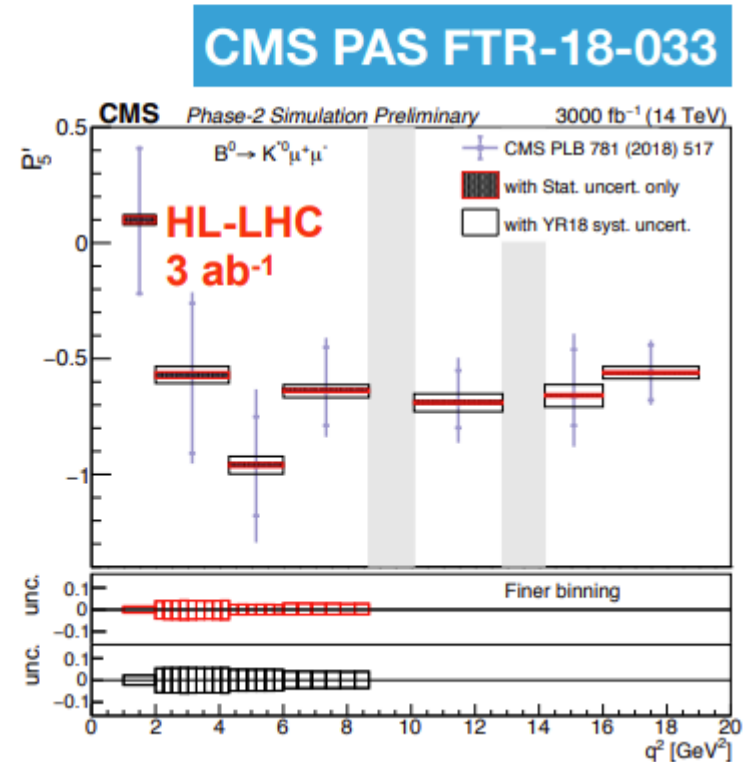
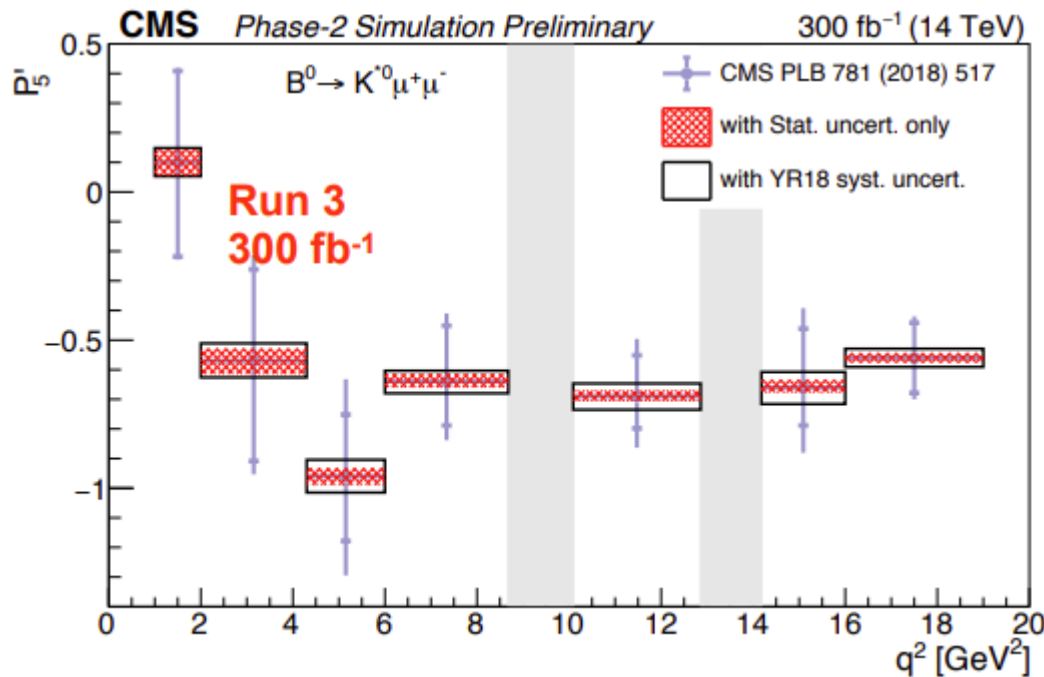
Table 2: Measured CP-averaged angular observables, in the corresponding q^2 bins. The first uncertainties are statistical and the second systematic.

	$1.1 < q^2 < 2 \text{ GeV}^2$	$2 < q^2 < 4.3 \text{ GeV}^2$	$4.3 < q^2 < 6 \text{ GeV}^2$
F_L	$0.709^{+0.073}_{-0.054} \pm 0.021$	$0.810^{+0.036}_{-0.030} \pm 0.016$	$0.714^{+0.032}_{-0.030} \pm 0.012$
P_1	$0.089^{+0.234}_{-0.204} \pm 0.040$	$-0.285^{+0.187}_{-0.208} \pm 0.051$	$-0.297^{+0.153}_{-0.168} \pm 0.038$
P_2	$-0.374^{+0.173}_{-0.125} \pm 0.095$	$-0.244^{+0.094}_{-0.077} \pm 0.039$	$0.121^{+0.080}_{-0.076} \pm 0.030$
P_3	$-0.045^{+0.209}_{-0.216} \pm 0.044$	$-0.187^{+0.196}_{-0.218} \pm 0.089$	$-0.027^{+0.143}_{-0.143} \pm 0.081$
P'_4	$-0.436^{+0.289}_{-0.323} \pm 0.111$	$-0.431^{+0.160}_{-0.185} \pm 0.075$	$-0.717^{+0.154}_{-0.158} \pm 0.074$
P'_5	$0.363^{+0.165}_{-0.132} \pm 0.028$	$-0.139^{+0.103}_{-0.087} \pm 0.039$	$-0.435^{+0.096}_{-0.101} \pm 0.027$
P'_6	$0.000^{+0.094}_{-0.097} \pm 0.021$	$0.108^{+0.075}_{-0.071} \pm 0.018$	$0.129^{+0.074}_{-0.071} \pm 0.011$
P'_8	$-0.157^{+0.368}_{-0.369} \pm 0.113$	$-0.727^{+0.193}_{-0.184} \pm 0.056$	$0.007^{+0.215}_{-0.216} \pm 0.036$
	$6 < q^2 < 8.68 \text{ GeV}^2$	$10.09 < q^2 < 12.86 \text{ GeV}^2$	$14.18 < q^2 < 16 \text{ GeV}^2$
F_L	$0.627^{+0.016}_{-0.016} \pm 0.011$	$0.474^{+0.011}_{-0.013} \pm 0.009$	$0.394^{+0.012}_{-0.012} \pm 0.009$
P_1	$-0.056^{+0.101}_{-0.102} \pm 0.046$	$-0.439^{+0.051}_{-0.047} \pm 0.030$	$-0.465^{+0.037}_{-0.037} \pm 0.025$
P_2	$0.188^{+0.039}_{-0.040} \pm 0.014$	$0.386^{+0.021}_{-0.019} \pm 0.018$	$0.440^{+0.008}_{-0.010} \pm 0.008$
P_3	$0.099^{+0.092}_{-0.090} \pm 0.014$	$0.013^{+0.041}_{-0.043} \pm 0.007$	$-0.034^{+0.037}_{-0.038} \pm 0.010$
P'_4	$-0.949^{+0.102}_{-0.101} \pm 0.058$	$-1.025^{+0.064}_{-0.066} \pm 0.059$	$-1.159^{+0.042}_{-0.038} \pm 0.041$
P'_5	$-0.495^{+0.067}_{-0.067} \pm 0.023$	$-0.746^{+0.033}_{-0.032} \pm 0.014$	$-0.688^{+0.038}_{-0.036} \pm 0.021$
P'_6	$0.010^{+0.052}_{-0.052} \pm 0.016$	$0.080^{+0.037}_{-0.041} \pm 0.011$	$0.121^{+0.040}_{-0.039} \pm 0.011$
P'_8	$-0.061^{+0.143}_{-0.143} \pm 0.042$	$-0.093^{+0.104}_{-0.094} \pm 0.029$	$-0.011^{+0.086}_{-0.089} \pm 0.022$

P5' HL-LHC

Run 3 and HL-LHC projections

- Up to x15 improvement w/ 3 ab⁻¹ compared to the 8 TeV CMS result [PLB 781 (2018) 517]
- Should be possible to resolve the situation experimentally already in Run 3



CMS experiment

architecture of the CMS Trigger

Trigger System

- reduce the number of events from the **LHC collision rate (40 MHz)** to the data rate that can be stored, reconstructed and analysed **Offline (~~~1 kHz~~) O(few kHz)**
- maximising the physics reach of the experiment

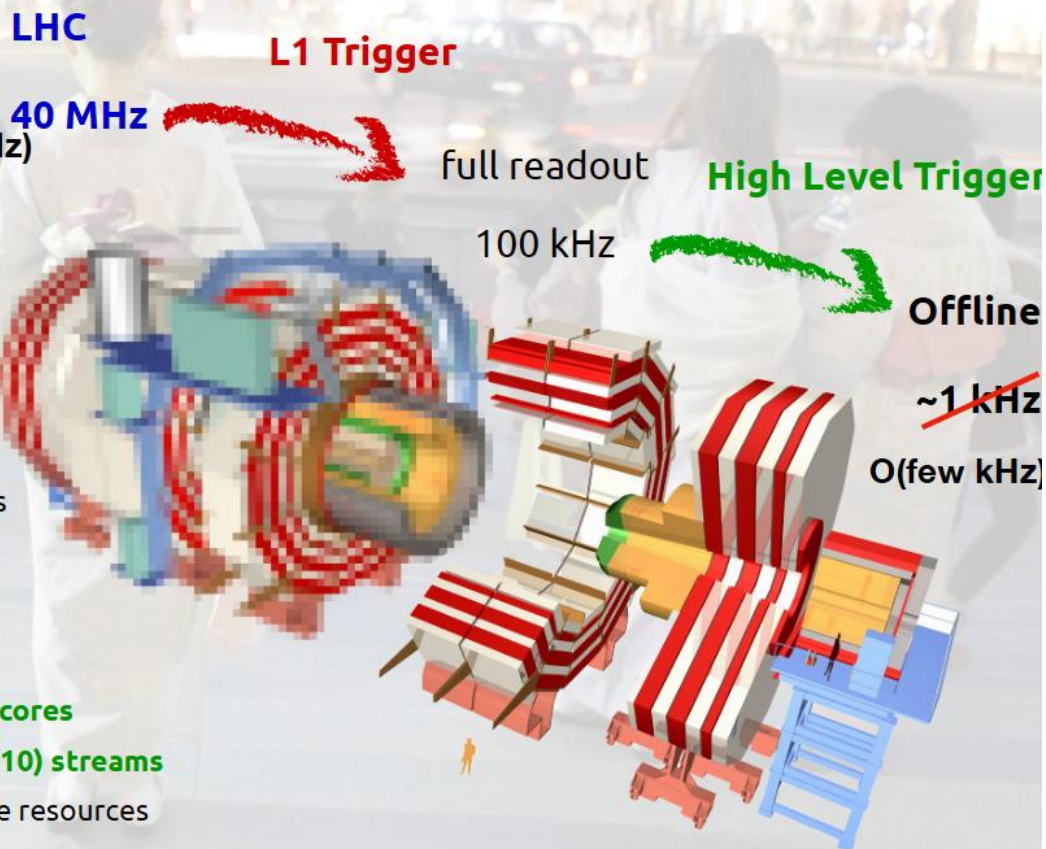
L1 Trigger

- coarse readout of the Calorimeters and Muon detectors
- implemented in custom electronics, ASICs and FPGAs
- **3 event types, 128 "physics" bits + 64 "technical" bits**
- output rate limited to **100 kHz** by the readout electronics
a bit more than 100 kHz in Run3

High Level Trigger

- readout of the whole detector with full granularity
- based on the CMSSW software, running on **O(15k) Xeon cores**
- organised in **O(2500) modules, O(400) trigger paths, O(10) streams**
- output rate limited to an average of ~~~1 kHz~~ by the Offline resources

O(few kHz)



$B^0 \rightarrow \phi \mu \mu$

LHCb-PAPER-2021-042, JHEP05(2022)067

It proceeds mainly via the color-suppressed penguin annihilation diagrams $\mathcal{O}(10^{-12})$

$\omega - \phi$ mixing can have a sizeable contribution $\mathcal{O}(10^{-10})$

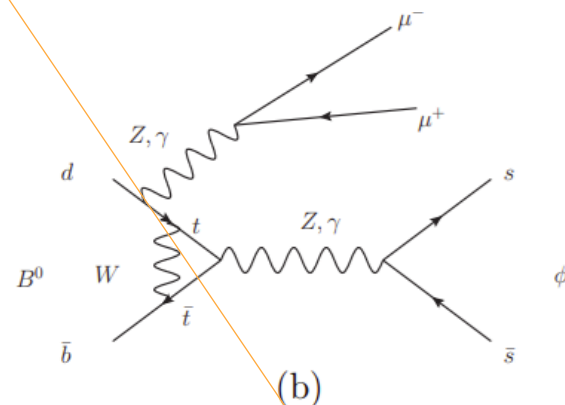
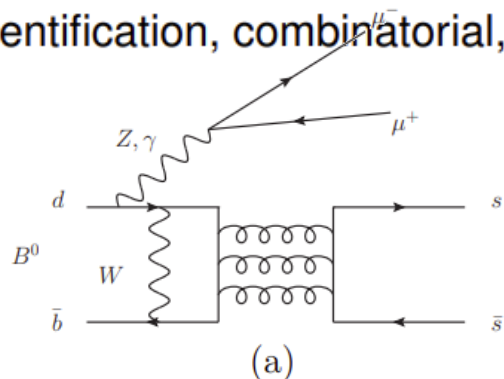
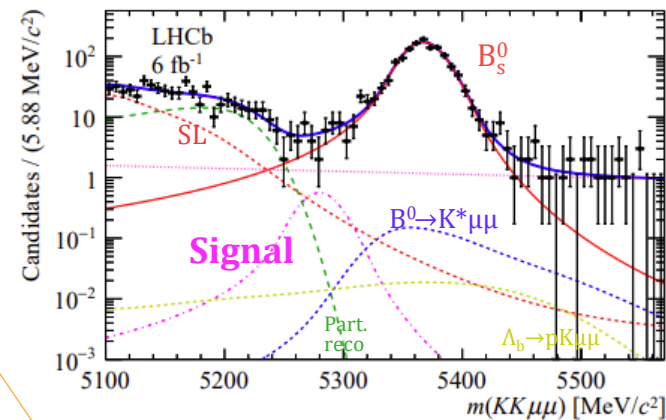
Run 1 + 2 data

Exclude regions in q^2 corresponding to ϕ , J/ψ and $\psi(2S)$

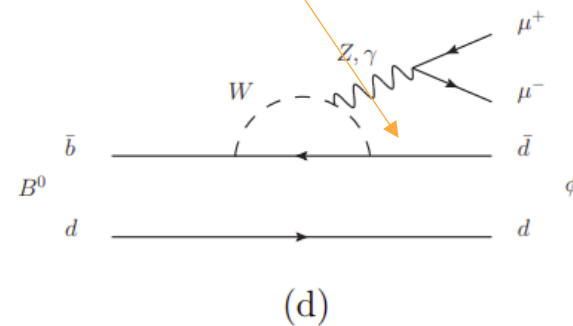
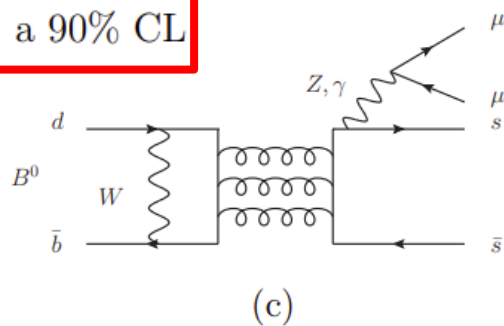
$B_s^0 \rightarrow \phi \mu^+ \mu^-$ used as normalisation

$B_s^0 \rightarrow J/\psi \phi$ used to develop a MVA discriminator

Dominant backgrounds: misidentification, combinatorial, semileptonic



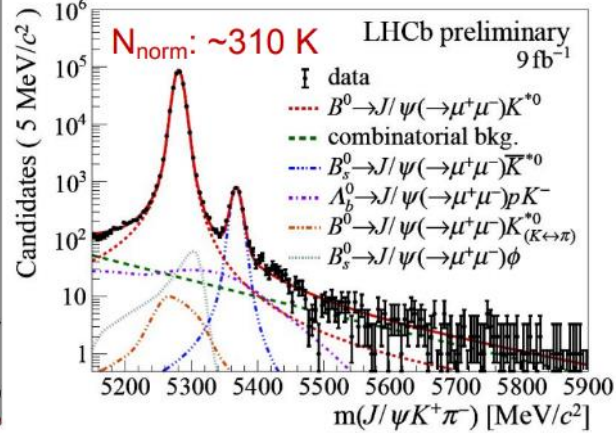
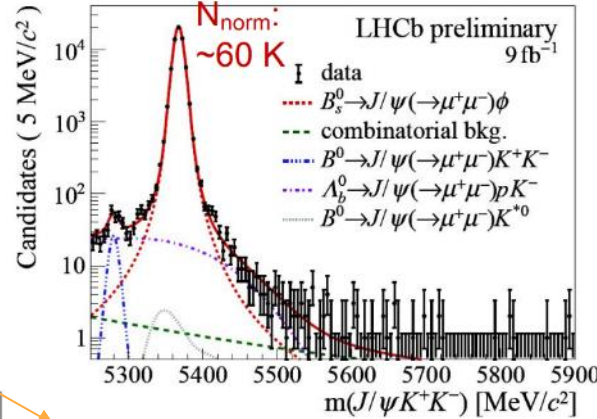
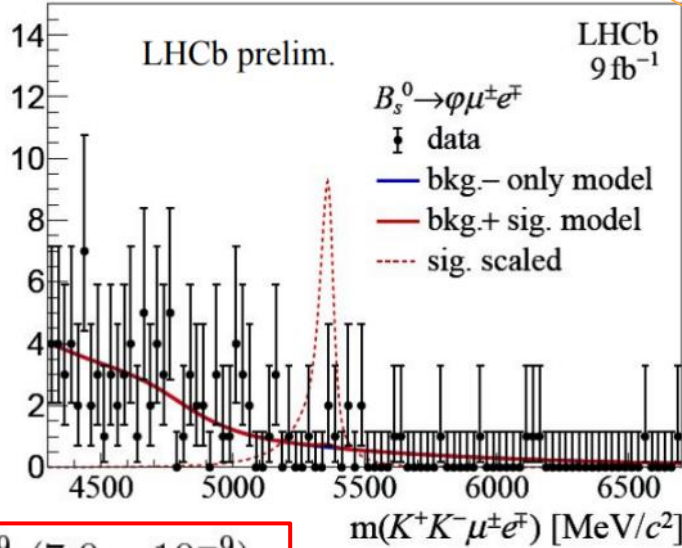
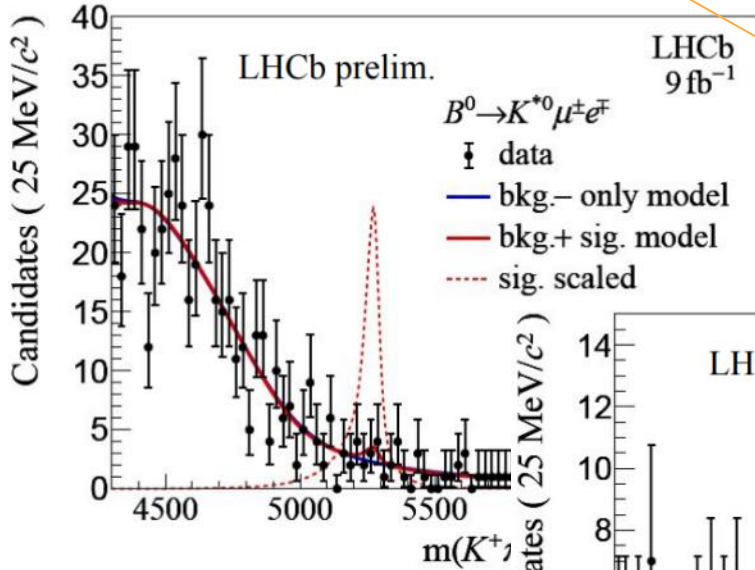
$$\mathcal{B}(B^0 \rightarrow \phi \mu^+ \mu^-) < 3.2 \times 10^{-9} \text{ at a 90\% CL}$$



$B^0 \rightarrow K^{*0} \mu^\mp e^\pm$ and $B_s^0 \rightarrow \phi \mu^\mp e^\pm$

LHCP-PAPER-2022-008. Full Run-1 + Run-2 analysis

Normalization using $B^0 \rightarrow J/\psi K^{*0}$ and $B_s^0 \rightarrow J/\psi \phi$



$$\mathcal{B}(B^0 \rightarrow K^{*0} \mu^+ e^-) < 5.7 \times 10^{-9} \quad (7.0 \times 10^{-9}),$$

$$\mathcal{B}(B^0 \rightarrow K^{*0} \mu^- e^+) < 6.7 \times 10^{-9} \quad (7.9 \times 10^{-9}),$$

$$\mathcal{B}(B^0 \rightarrow K^{*0} \mu^\pm e^\mp) < 9.9 \times 10^{-9} \quad (11.6 \times 10^{-9})$$

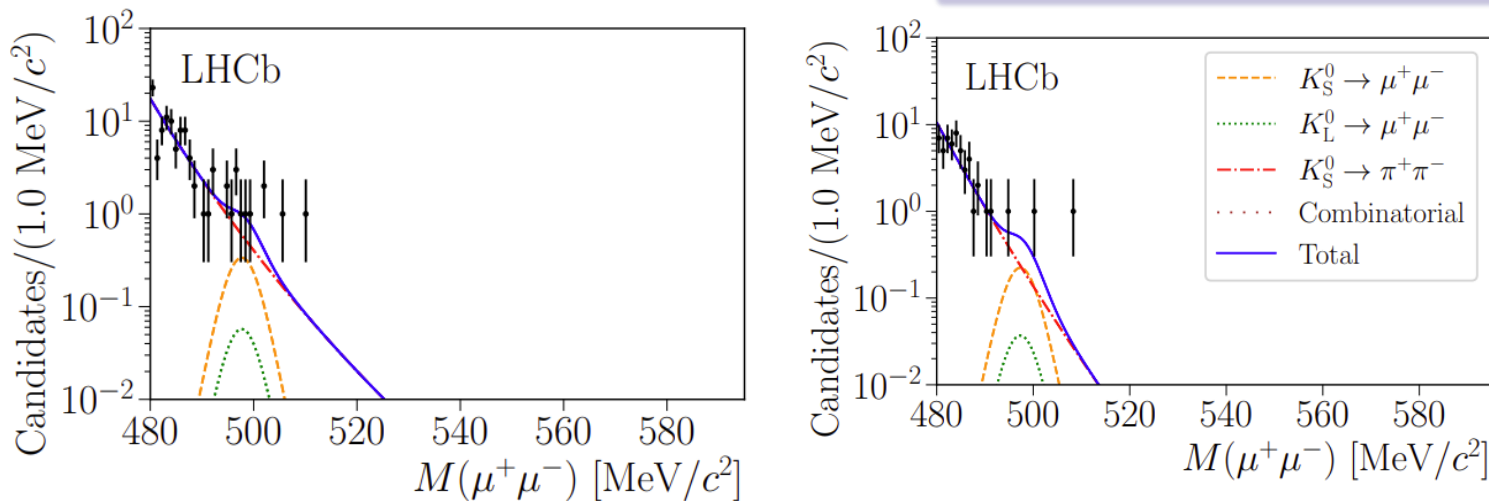
$$\mathcal{B}(B_s^0 \rightarrow \phi \mu^\pm e^\mp) < 15.9 \times 10^{-9} \quad (19.4 \times 10^{-9})$$

World-best limits

$K_S^0 \rightarrow \mu\mu$

- FCNC process, in SM $\mathcal{B}(K_S^0 \rightarrow \mu^+\mu^-)_{\text{SM}} = (5.18 \pm 1.50_{\text{LD}} \pm 0.02_{\text{SD}}) \times 10^{-12}$
- Some NP (SUSY/LQ) models modify the **B**
- LHCb performed a search using Run-2 data
- Normalization using decay to $\pi^+\pi^-$
 - This decay is also the main background

LHCb-PAPER-2019-038, Phys.Rev.Lett.125(2020)231801



Statistically combined with Run-1 result upper limit is **most stringent to date:**

$$\mathcal{B}(K_S^0 \rightarrow \mu^+\mu^-) < 2.1 \times 10^{-10} \text{ at 90\% CL}$$

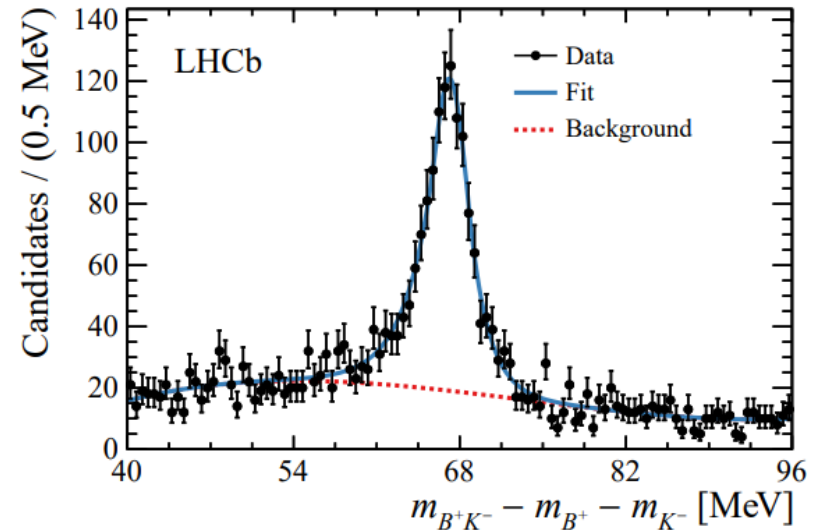
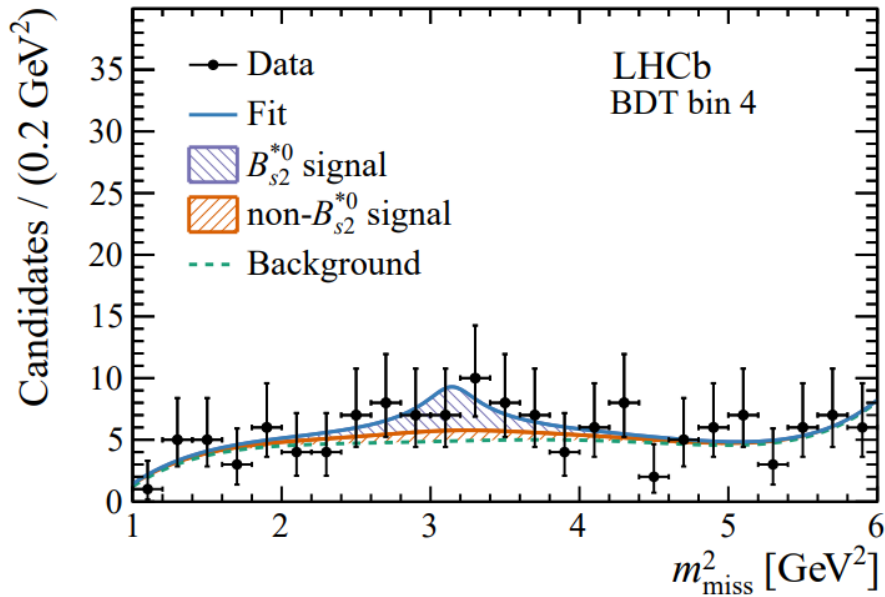
$B^+ \rightarrow K^+ \mu^- \tau^+$

LHCb-PAPER-2019-043, JHEP06(2020)129

LHCb Full Run 1 + Run 2 analysis

Using $B_{s2}^* \rightarrow B^+ K^-$ decays to tag partially-reconstructed B^+ mesons

Normalization using $B^+ \rightarrow K^+ \mu^- \mu^+$ (with J/ψ)



$$\mathcal{B}(B^+ \rightarrow K^+ \mu^- \tau^+) < 3.9 \times 10^{-5} \text{ at } 90\% \text{ CL}$$

Weaker than 2012 BaBar upper limit

$$\mathcal{B}(B^+ \rightarrow K^+ \mu^- \tau^+) < 2.8 \times 10^{-5} \text{ at } 90\%$$

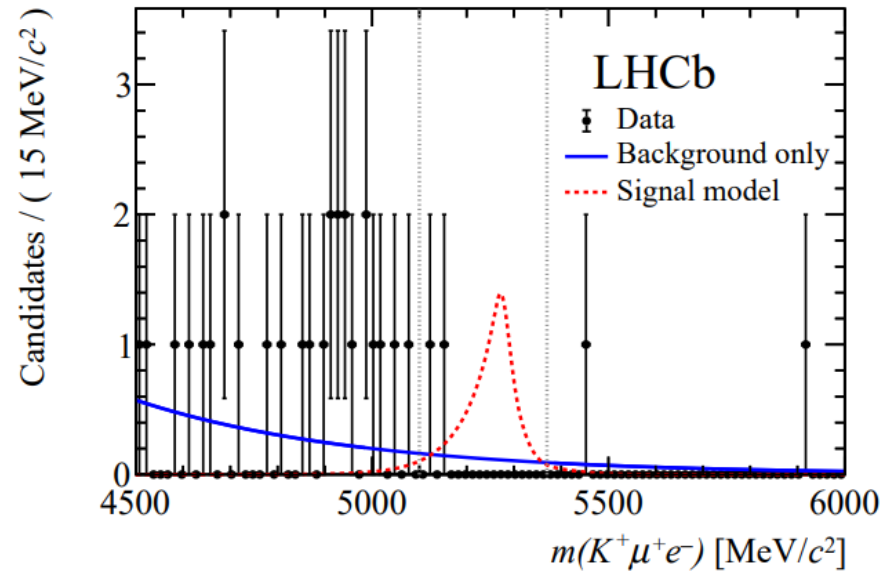
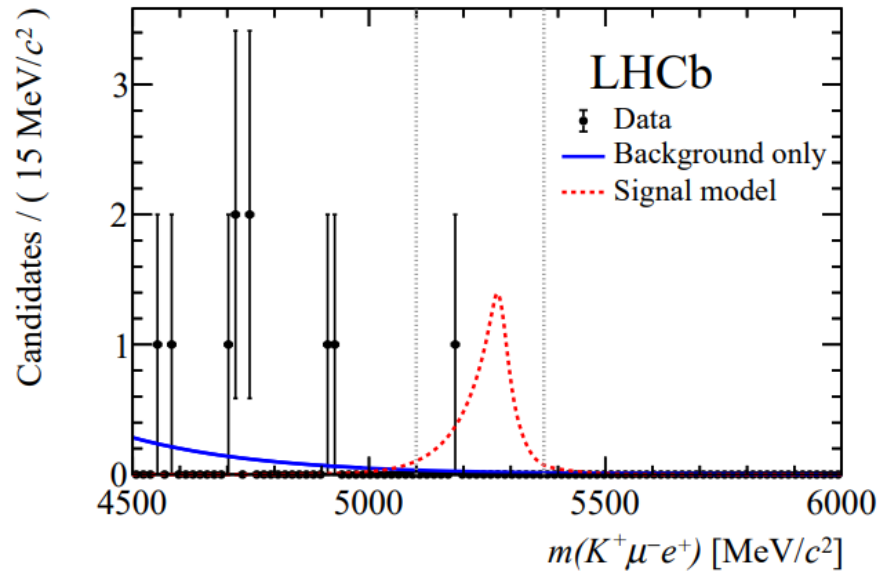
[More details in the parallel talk by Liang S. yesterday](#)

$B^+ \rightarrow K^+ \mu^\mp e^\pm$

LHCb Run-1 analysis

LHCb-PAPER-2019-022, [Phys.Rev.Lett.123\(2019\)241802](#)

Normalization using $B^+ \rightarrow K^+ \mu^- \mu^+$ (with J/ψ)



$$\mathcal{B}(B^+ \rightarrow K^+ \mu^- e^+) < 7.0 \text{ (9.5)} \times 10^{-9}$$

$$\mathcal{B}(B^+ \rightarrow K^+ \mu^+ e^-) < 6.4 \text{ (8.8)} \times 10^{-9}$$

World-best limits

[More details in the parallel talk by Liang S. yesterday](#)

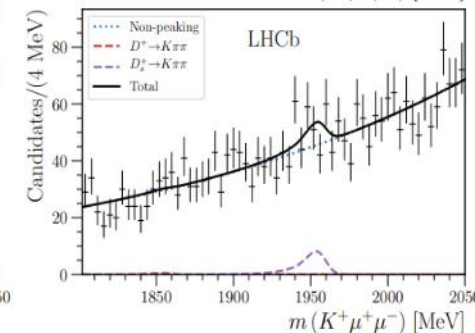
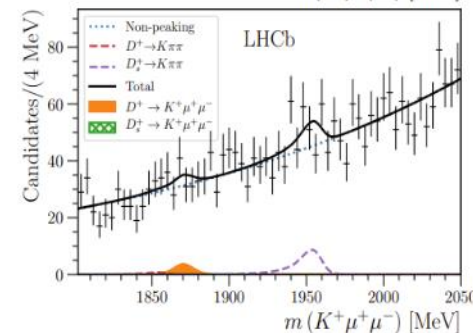
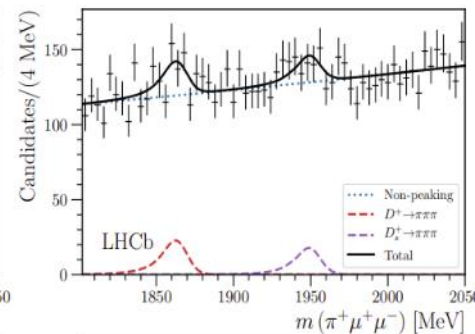
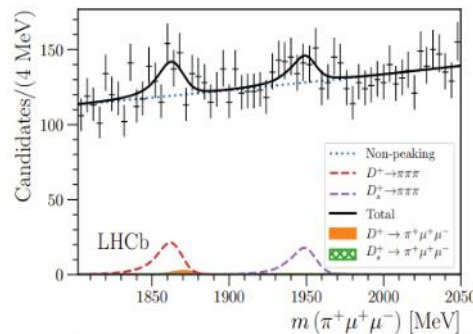
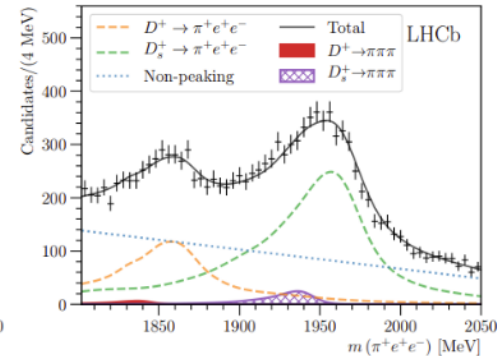
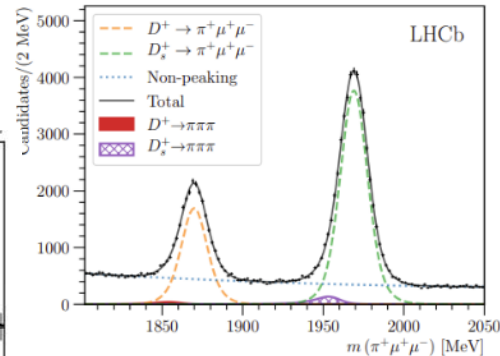
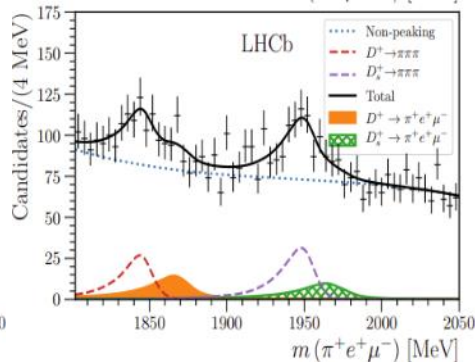
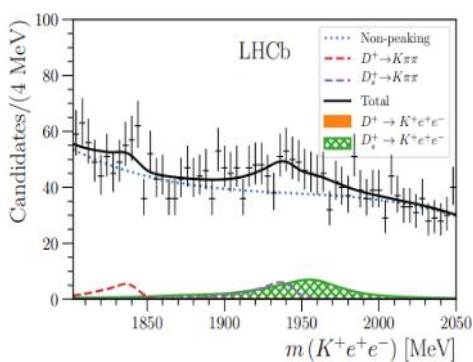
Run-2 analysis ongoing

Rare charm decays

LHCb [JHEP 06 \(2021\) 044](#)

Searches for 25 rare and forbidden decays of D^+ and D_s^+ mesons

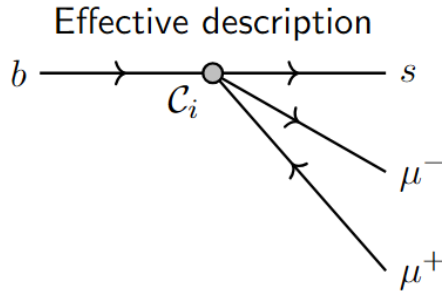
$$\mathcal{B}_{D_{(s)}^+ \rightarrow h^\pm \ell^+ \ell(\ell)^\mp} = \frac{N_{D_{(s)}^+ \rightarrow h^\pm \ell^+ \ell(\ell)^\mp}}{N_{D_{(s)}^+ \rightarrow (\phi \rightarrow \mu^+ \mu^-) \pi^+}} \cdot \frac{\epsilon_{D_{(s)}^+ \rightarrow (\phi \rightarrow \mu^+ \mu^-) \pi^+}}{\epsilon_{D_{(s)}^+ \rightarrow h^\pm \ell^+ \ell(\ell)^\mp}} \cdot \mathcal{B}_{D_{(s)}^+ \rightarrow (\phi \rightarrow \mu^+ \mu^-) \pi^+}$$



Decay	Branching fraction upper limit [10^{-9}]			D_s^+		
	SES	90% CL	95% CL	SES	90% CL	95% CL
$D_{(s)}^+ \rightarrow \pi^+ \mu^+ \mu^-$	0.6	67	74	2.4	180	210
$D_{(s)}^+ \rightarrow \pi^- \mu^+ \mu^+$	0.3	14	16	1.8	86	96
$D_{(s)}^+ \rightarrow K^+ \mu^+ \mu^-$	1.2	54	61	3.8	140	160
$D_{(s)}^+ \rightarrow K^- \mu^+ \mu^+$	-	-	-	1.2	26	30
$D_{(s)}^+ \rightarrow \pi^+ e^+ \mu^-$	0.6	210	230	3.1	1100	1200
$D_{(s)}^+ \rightarrow \pi^+ \mu^+ e^-$	0.4	220	220	2.2	940	1100
$D_{(s)}^+ \rightarrow \pi^- \mu^+ e^+$	0.4	130	150	2.0	630	710
$D_{(s)}^+ \rightarrow K^+ e^+ \mu^-$	0.7	75	83	3.7	790	880
$D_{(s)}^+ \rightarrow K^+ \mu^+ e^-$	0.5	100	110	2.5	560	640
$D_{(s)}^+ \rightarrow K^- \mu^+ e^+$	-	-	-	2.4	260	320
$D_{(s)}^+ \rightarrow \pi^+ e^+ e^-$	1.9	1600	1800	8.1	5500	6400
$D_{(s)}^+ \rightarrow \pi^- e^+ e^+$	0.9	530	600	4.1	1400	1600
$D_{(s)}^+ \rightarrow K^+ e^+ e^-$	4.4	850	1000	14.8	4900	5500
$D_{(s)}^+ \rightarrow K^- e^+ e^+$	-	-	-	4.1	770	840

EFT for $b \rightarrow sll$ decays

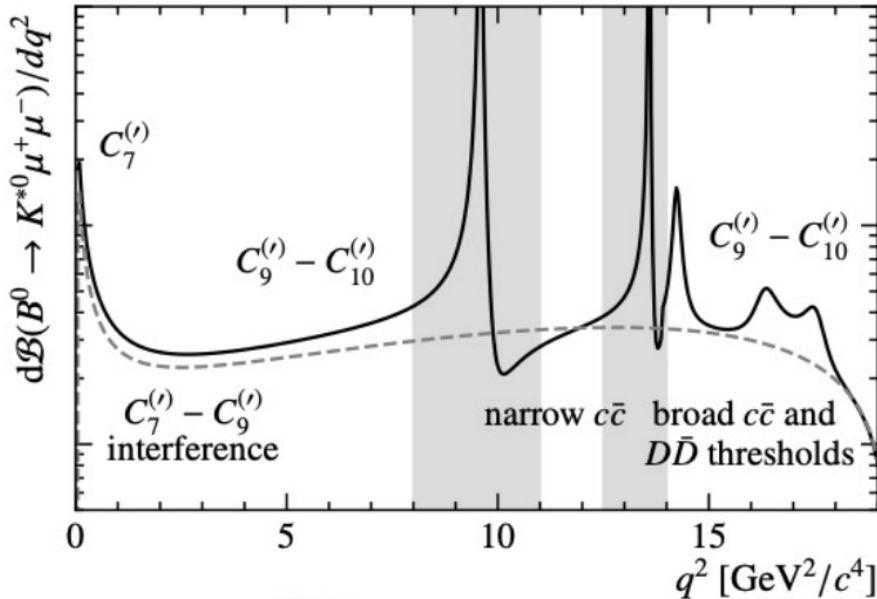
In general, $b \rightarrow sll$ transitions can be described using an EFT approach



$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \frac{e^2}{16\pi^2} \sum_i C_i O_i$$

Effective coupling
"Wilson coefficient"
short distance physics

Local operator
long distance physics



NP could modify the Wilson Coefficients,
or enable new operators

different dimuon mass (q) ranges
sensitive to different operators

NP contribution to C_9 expected to
be constant vs q^2 (M. B. talk)