**Analysis of experimental data on neutron decay for the possibility of the existence of the right vector boson** 

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**7th International Conference on Particle Physics and Astrophysics** 22-25 October 2024

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#### **1. Precision studies of neutron decay and the search for deviations from the Standard Model**

**2. Inconsistencies between experimental results and the − theory of weak interaction**

**3. The decay of a neutron within the left-right manifest model of mixing left and right vector bosons can be successfully described by**  $W_R$  **with mass**  $(M_{W_R} = 304^{+28}_{-22}$  GeV,  $\zeta = -0.038$ )

**4. Why wasn't**  $W_R(300 \text{ GeV})$ ,  $\zeta = -0.038$ ) detected at FermiLab and CERN?

#### **5. Prospects for neutron decay experiments**

#### **Precision studies of neutron decay and the search for deviations from the Standard Model Neutron Spin J**  $\frac{d^3\Gamma}{dE_e d\Omega_e d\Omega_v} = \frac{1}{2(2\pi)^5} G_F^2 |V_{ud}|^2 (1+3|\lambda|^2) p_e E_e (E_0 - E_e)^2$  Jackson, Treiman, Wyld,<br> $\frac{dE_e d\Omega_e d\Omega_v}{dE_e d\Omega_v d\Omega_v} = \frac{1}{2(2\pi)^5} G_F^2 |V_{ud}|^2 (1+3|\lambda|^2) p_e E_e (E_0 - E_e)^2$  Jackson, Treiman, Wyld, **Electron** в Neutrino p<sub>v</sub>  $\times \begin{bmatrix} 1 \\ -2 \frac{\lambda^2 + \lambda}{1 + 3\lambda^2} \\ 2 \frac{\lambda^2 - \lambda}{1 + 3\lambda^2} \\ = g_{\text{A}}/g_{\text{V}} \\ \frac{(1 - \lambda^2)}{(1 + 3\lambda^2)} \end{bmatrix}$  $\begin{array}{c|c}\n-2 & \frac{\lambda^2 + \lambda}{1 + 3\lambda^2} & \frac{\lambda^2 - \lambda}{1 + 3\lambda^2} \\
2 & \frac{\lambda^2 - \lambda}{1 + 3\lambda^2} & \frac{\lambda^2}{1 + 3\lambda^2} \\
= g_{\text{A}}/g_{\text{V}} & & \\
\frac{(1 - \lambda^2)}{(1 + 3\lambda^2)} & & \\
2 & \frac{\text{Im}(\lambda)}{1 + 3\lambda^2} & & \\
\end{array}$ Proton p<sub>p</sub> **Neutron lifetime -0.11958(21) 0.17% -0.11958(21) 0.17%** $\int \tau^{-1} = G_F^2 |V_{ud}|^2 (1+3\lambda^2) \frac{f^k m_e^3 c^4}{2\sigma^3 h^2}$  $877.75 \pm 0.35$ s  $0.04\%$ **0.9807(30) 0.3% Unitarity CKM**  $\begin{bmatrix} V_{ud} & V_{us} & V_{ub} \ V_{cd} & V_{cs} & V_{cb} \ V_{td} & V_{ts} & V_{tb} \end{bmatrix}$ **-1.2757(5) 0.04%**  $W_L$  $2<sub>1</sub>$  $-\lambda^2$ )  $\lambda^2$ )  $a = \frac{(1 + 2i)^{2}}{(1 + 2i)^{2}}$  $=\frac{(1 - \lambda)^{2}}{(1 - 2 \lambda)^{2}}$  **-0.1049(13) 1.3%**  $\lambda^2$  | - $V_{ud}^{unit} = \sqrt{1 - V_{us}^2 - V_{ub}^2}$  $W_R$  $D = 2 \cdot \frac{\text{Im}(\lambda)}{1 + 3 |\lambda|^2}$  $-1.2$  (2.0) $\times$ 10<sup>-4</sup>  $= 0.97452(18)$ .

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### **Improving the accuracy of measurements and trends in the neutron lifetime**



Experimental results on neutron lifetime since 1990 from [8], discrepancy between 2005 data [9] and 2000 data [10], new magnetic trap results (marked in green) which are decisive [11-14].

#### **Measuring neutron decay asymmetries**



Measurements of the neutrino asymmetry of neutron decay (B) and the averaged PDG result [18].

Measurements of the electron-neutrino asymmetry of neutron decay (a) and the averaged PDG result [18].

year

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#### Superallowed  $0^+ \rightarrow 0^+$  nuclear  $\beta$  decays: 2020 critical survey, with implications for  $V_{ud}$  and CKM unitarity

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Z of daughter

A new critical review of all half-life, decay energy, and branching ratio measurements associated with 23 superallowed  $0+ \rightarrow 0+$  is available. Their average Ft combined with the muon lifetime yields the up-down quark mixing element of the Cabibbo-Kobayashi-Maskawa matrix, **Vud** =  $0.97373 \pm 0.00031$ . This is one standard deviation lower than our 2015 result, and its uncertainty has increased by 50%. This is not a consequence of any shifts in the experimental data, but of new calculations for the radiative corrections. **The lower Vud now leads to a higher tension in the top-row unitarity test in the CKM Matrix.**

This result is given in the last row of Table XVII: where the unitarity sum is  $|Vu|2 = 0.9985(6)$ , indicating **unitarity violation of 2.4σ.**

**Inconsistencies between experimental results and the** *V***−***A* **theory of weak interaction**

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#### **The difference Vud between the matching values and the Vud value from 0+-0+**  s **transitions is 2.6 sigma** .



Dependence of the quark mixing matrix element Vud on  $\lambda$  , calculated using the SM formulas from neutron decay, from experiments with Fermi supereallowed nuclear  $0+ - 0+$ transitions and from the unitarity of the SM matrix, using  $Vus$  measurements [18].

$$
\frac{\Delta V_{ud}}{V_{ud}} = 8.6 * 10^{-4} (2.6 \,\sigma)
$$

**There is a discrepancy between the experimental value of neutrino asymmetry and the SM prediction. The difference between these values is 2.1 sigma**



Comparison of the **experimental neutrino asymmetry** of neutron decay and that calculated within the SM framework depending on the ratio of the axial and vector constants of weak interaction  $\lambda$ .

$$
\frac{\Delta B}{B} = 6.5 * 10^{-3} (2.1 \sigma)
$$
\n
$$
\frac{\Delta B}{B} = \frac{B_{\text{exp}} - B_{\text{SM}}}{B_{\text{SM}}} \qquad B_{\text{SM}} = \frac{2\lambda_n (\lambda_n - 1)}{(1 + 3\lambda_n^2)}
$$

**Data** 
$$
|Vus|
$$
 from PDG  
 $V_{us} = 0.2243(8)$ 



The third element of the top row, **|Vub|, is very small**  and has almost no effect on the unitarity test. Its value from the Particle Data Group (PDG) evaluation is:

$$
|Vub| = (3.94 \pm 0.36) \times 10-3
$$

# *V*<sup>*unit*</sup> from the unitarity **of the CKM matrix**

$$
V_{ud}^{unit} = \sqrt{1 - V_{us}^2 - V_{ub}^2} = 0.97452(18).
$$

however, the matrix element  $V_{ud}^{00}$  from  $0^+ - 0^+$ beta decays is different

$$
V_{ud}^{00} = 0.97367(32)
$$

$$
\frac{V_{ud}^{unit} - V_{ud}^{00}}{V_{ud}^{00}} = 8.6 * 10^{-4} (2.4 \sigma)
$$

# **The description of experimental results within the framework of the V-A version of the theory turns out to be unsatisfactory, since it cannot be represented by a single value of the parameter**  $\lambda = G_{\lambda}/G_{\nu}$  $\begin{array}{c}\n\mathbf{a} \text{ is in } \mathbf{f} \text$





 $a_{\rm exp} = -0.10402(82)$  $A_{\rm exp} = -0.11958(21)$  $B_{\rm exp} = 0.9807(30)$  $V_{ud}^{unit} = 0,97452(18)$ 

Results of calculating the parameter value  $\lambda = G_{\lambda}/G_{\nu}$ within the V-A version of the weak interaction theory, the experiments for a, A, B and  $\tau$  cannot be represented by a single value.

**The observed discrepancy can be analyzed within the framework of a model taking into account right-handed currents. In the simplest leftright manifesto of the model, mixing of left and right vector bosons is considered, and for current states**  $W_L$  $W_R$  **and mass states**  $W_1$  $W_2$ **we can write: pancy can be analy:**<br> **count right-handed**<br> **condel, mixing of light-handed**<br> **rrent states**  $W_L$   $W_L$ <br>  $\frac{W_L}{W_L}$ <br>  $\frac{1}{1} \cos \zeta + W_2 \sin \zeta$ <br>  $\frac{1}{1} \cos \zeta - W_1 \sin \zeta + W_2 \sin \zeta$ <br> **condensizing**  $\zeta$ <br> **condensizing** of cu<br> an be analyzed with<br>ght-handed curren<br>mixing of left and<br>ates  $W_L$   $W_R$  and<br> $W_L$   $W_R$  and<br> $W_L$  and  $\frac{V_1}{V_1}$  sin  $\frac{V_2}{V_1}$ <br>of the squares of the m<br> $W_1$  and  $W_2$ . ancy can be analyzed with<br>unt right-handed currents<br>model, mixing of left and<br>rent states  $W_L$   $W_R$  and n<br> $\overline{\cos \zeta + W_2 \sin \zeta}$ <br> $\overline{\cos \zeta - W_1 \sin \zeta + W_2 \cos \zeta}$ <br>angle of mixing of current state<br>- ratio of the squares of the m ncy can be analyzed within the fr<br>nt right-handed currents. In the<br>nodel, mixing of left and right ve<br>nt states  $W_L$   $W_R$  and mass stat<br> $\overline{\cos \zeta + W_2 \sin \zeta}$ <br> $(-W_1 \sin \zeta + W_2 \cos \zeta)$ <br>ngle of mixing of current states<br>ratio of **Example 1** discrepancy can be analyzed within<br>  $\frac{1}{2}$  into account right-handed currents.<br>
sto of the model, mixing of left and rind for current states  $W_L$   $W_R$  and ma<br>  $W_L = W_1 \cos \zeta + W_2 \sin \zeta$ <br>  $W_R = e^{-i\omega} (-W_1 \sin \zeta + W_2$ **Example 1** discrepancy can be analyzed within the fraction of the model, mixing of left and right vector of the model, mixing of left and right vector and for current states  $W_L$   $W_R$  and mass state e:<br>  $W_L = W_1 \cos \zeta + W_2 \$ can be analyzed within the framework of a<br>right-handed currents. In the simplest left-<br>el, mixing of left and right vector bosons is<br>states  $W_L$   $W_R$  and mass states  $W_1$   $W_2$ <br> $\zeta + W_2 \sin \zeta$ <br> $W_1 \sin \zeta + W_2 \cos \zeta$ <br>e of mix screpancy can be analyzed within the framework of a<br>
in account right-handed currents. In the simplest left-<br>
if the model, mixing of left and right vector bosons is<br>
in current states  $W_L$   $W_R$  and mass states  $W_1$   $W_2$ screpancy can be analyzed within the framework of a<br>
in account right-handed currents. In the simplest left-<br>
if the model, mixing of left and right vector bosons is<br>
in current states  $W_L$   $W_R$  and mass states  $W_1$   $W_2$ 

$$
W_L = W_1 \cos \zeta + W_2 \sin \zeta
$$
  

$$
W_R = e^{-i\omega} (-W_1 \sin \zeta + W_2 \cos \zeta)
$$

where  $\zeta$  is the angle of mixing of current states  $W_l$  and  $W_R$ ,  $\delta$  – ratio of the squares of the masses of states  $W_1$  and  $W_2$ .

 $\omega$  - CP-violating phase



#### **V-A variant of the theory left-right manifest model**

$$
\tau_{\text{exp}} \pm \Delta \tau_{\text{exp}} = \frac{4905,7}{V_{ud}^2 [1 + x^2 + 3\lambda^2 (1 + y^2)]}
$$
\n
$$
a_{\text{exp}} \pm \Delta a_{\text{exp}} = \frac{(1 - \lambda^2)[1 + (\delta + \zeta)^2] - 4\delta\zeta}{(1 + 3\lambda^2)[1 + (\delta + \zeta)^2] - 4\delta\zeta}
$$
\n
$$
A_{\text{exp}} \pm \Delta A_{\text{exp}} = -\frac{2\lambda[\lambda(1 - y^2) + (1 - xy)]}{1 + x^2 + 3\lambda^2 (1 + y^2)}
$$
\n
$$
B_{\text{exp}} \pm \Delta B_{\text{exp}} = \frac{2\lambda[\lambda(1 - y^2) - (1 - xy)]}{1 + x^2 + 3\lambda^2 (1 + y^2)}
$$
\nWhere  $x = \delta - \zeta$ ,  $y = \delta + \zeta$ .

#### **Expansion in**  $\delta$  **and**  $\zeta$  **of order no higher than two can be represented by the following expressions**

$$
\tau_{\exp} \pm \Delta \tau_{\exp} = \frac{4905,7}{V_{ud}^2 [1 + x^2 + 3\lambda^2 (1 + y^2)]}
$$
\n
$$
a_{\exp} \pm \Delta a_{\exp} = \frac{(1 - \lambda^2)[1 + (\delta + \zeta)^2] - 4\delta\zeta}{(1 + 3\lambda^2)[1 + (\delta + \zeta)^2] - 4\delta\zeta}
$$
\n
$$
A_{\exp} \pm \Delta A_{\exp} = -\frac{2\lambda [\lambda (1 - y^2) + (1 - xy)]}{1 + x^2 + 3\lambda^2 (1 + y^2)}
$$
\n
$$
B_{\exp} \pm \Delta B_{\exp} = \frac{2\lambda [\lambda (1 - y^2) - (1 - xy)]}{1 + x^2 + 3\lambda^2 (1 + y^2)}
$$

no higher than two can be represented by the  
\nowing expressions\n
$$
\frac{\tau_{\rm exp} \pm \Delta \tau_{\rm exp} - \tau_{V-A}}{\tau_{V-A}} \simeq -\left[\delta^2 + \zeta^2 + 2\frac{(3\lambda^2 - 1)}{(3\lambda^2 + 1)}\delta\zeta\right]
$$
\n
$$
\frac{a_{\rm exp} \pm \Delta a_{\rm exp} - a_{V-A}}{a_{V-A}} \simeq -\frac{16}{(1-\lambda^2)(1+3\lambda^2)}\delta\zeta
$$
\n
$$
\frac{A_{\rm exp} \pm \Delta A_{\rm exp} - A_{V-A}}{A_{V-A}} \simeq -2\delta^2 - 2\delta\zeta \frac{\left[6\lambda^3 + 3\lambda^2 - 1\right]}{(\lambda+1)(1+3\lambda^2)} - 2\frac{\lambda}{\lambda+1}\zeta^2}
$$
\n
$$
\frac{B_{\rm exp} \pm \Delta B_{\rm exp} - B_{V-A}}{B_{V-A}} \simeq -2\delta^2 - 2\delta\zeta \frac{\left[6\lambda^3 - 3\lambda^2 + 1\right]}{(\lambda-1)(1+3\lambda^2)} - 2\frac{\lambda}{\lambda-1}\zeta^2}
$$
\n14

$$
\frac{a_{\exp} \pm \Delta a_{\exp}}{a_{V-A}} \simeq -\frac{10}{\left(1 - \lambda^2\right)\left(1 + 3\lambda^2\right)} \delta \zeta
$$

$$
\frac{[\lambda(1-y^2)-(1-xy)]}{(1+x^2+3\lambda^2(1+y^2))}\n\begin{array}{|l|}\nA_{\exp}\pm\Delta A_{\exp}-A_{V-A} & \Delta\\
\hline\nA_{V-A} & \Delta\end{array}\n\simeq -2\delta^2 - 2\delta\zeta \frac{\left[6\lambda^3+3\lambda^2-1\right]}{(\lambda+1)\left(1+3\lambda^2\right)} - 2\frac{\lambda}{\lambda+1}\zeta^2
$$

$$
\boxed{\frac{B_{\text{exp}} \pm \Delta B_{\text{exp}} - B_{\text{V-A}}}{B_{\text{V-A}}} \simeq -2\delta^2 - 2\delta \zeta \frac{\left[6\lambda^3 - 3\lambda^2 + 1\right]}{(\lambda - 1)\left(1 + 3\lambda^2\right)} - 2\frac{\lambda}{\lambda - 1}\zeta^2}
$$

#### The search for optimal values was done using the method  $\chi^2$



# **vector bosons can be successfully described**



#### **Optimal values of the parameters δ and ζ obtained by the χ2 method using**  experimental neutron decay data for  $a, A, B$  and  $\tau$



#### **Analysis of Fermi superallowed ^+−^+ transitions taking into account the influence of right-handed currents**

$$
(V_{ud}^{00})^2 = (V_{ud,SM}^{00})^2 [1 + (\delta + \zeta)^2]
$$

from neutron decay  $V_{ud}^n = 0.97477(37)$ , from the unitarity  $V_{ud}^{unit} = 0.97452(18)$ , from  $0^+ - 0^+$ transitions,  $V_{ud}^{00} = 0.97426(31)$ 

Dependence of the quark mixing matrix element  $V_{\mu}ud$  on calculated using the SM formulas from neutron decay (blue area). Determination of  $\lambda$  from the electron asymmetry of neutron  $decay - A (green area)$ . Determination of  $V_{ud}$  from the unitarity of the CM matrix, using  $V_{\perp}$ us measurements [18] (orange area). Determination of *V\_ud* from experiments with Fermi **superresolved nuclear 0+-0+ transitions after introducing a correction for the optimal parameters**  $\delta$  **and**  $\zeta$  **obtained in the analysis of neutron decay (shaded area).**  $\lambda$ 



#### **It is important to note that the coincidence was obtained when studying different objects - from neutron decay and from nuclear transitions.**



**It should be additionally noted that the contradiction noted in [22] as a violation of unitarity is eliminated. We explain this discrepancy within the framework of the left-right manifest model.**

Dependence of the parameter  $\delta$  on the parameter  $\zeta$  from equations (5) for the results of measuring the quantities  $a, A$ , B and  $\tau$  at the value  $\lambda$  opt=−1.2738 with an additional analysis for  $0^{\wedge}$ + $-0^{\wedge}$ + transitions from equation (11). Purple lines correspond to this additional analysis.

#### **Taking into account the accuracy of calculation of radiative corrections**

$$
\frac{1}{\tau_n} = \frac{G_F^2 |V_{ud}|^2}{2\pi^3} m_e^5 (1 + 3\lambda_n^2)(1 + RC)f
$$
  

$$
f = 1.6887(1)
$$

Radiative corrections in the form of a multiplier

 $(1 + RC)$  can be represented as a work  $(1 + RC) = (1 + \delta_R)(1 + \Delta_R)$ , where is the contribution

 $\delta_R = 0.01505$  arises from the exchange or emission of one photon exclusively, the contribution

 $\Delta_R = 0.02381$  – that part of the radiative corrections that is due to the exchange of electroweak boson and QCD corrections [24].

**Table** of radiative corrections to the neutron lifetime **and correlation coefficients of neutron decay asymmetries in percent.**



#### **Taking into account the accuracy of the unitarity condition**

The above analysis does not yet take into account the accuracy of determining  $V_{ud}^{unit}$ from experiments

with strange and charmed mesons

$$
V_{ud}^{unit} = \sqrt{1 - V_{us}^2 - V_{ub}^2} = 0.97452(18).
$$

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**Therefore, a final analysis was carried out, taking into account the accuracy of the calculation of**  radiative corrections for the neutron lifetime and the accuracy of determining  $V_{ud}^{unit}$ .



**Final result of the analysis**

**As a result of the analysis, it was found that there are indications of the existence of a right vector boson with mass and mixing angle**

$$
M_{W_R} = 304^{+28}_{-22} \text{ GeV}
$$

$$
\zeta = -0.038 \pm 0.014.
$$

**Comparison with the constraints on the mass of**  $W_R$  **and the mixing angle that were obtained earlier - in 1998 in [19] and in 2012 in [36].**

**Our work 1998 in [19]**

More precise constraints on the  $W_R$  mass and **mixing angle were obtained in muon decay:**



mixing angle, which were obtained earlier in [19] (shaded area C.L. 90%).

Permitted areas  $1\sigma$ ,  $2\sigma$  for the masses  $W_R$  and mixing angle  $\zeta$  in the LRS model from [36].

#### PHYSICAL REVIEW D 84, 032005 (2011) Precise measurement of parity violation in polarized muon decay

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#### (TWIST Collaboration)

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We present a new high precision measurement of parity violation in the weak interaction, using polarized muon decay. The TWIST Collaboration has measured  $P_{\mu}^{\pi} \xi$ , where  $P_{\mu}^{\pi}$  is the polarization of the muon in pion decay and  $\xi$  describes the intrinsic asymmetry in muon decay. We find  $P_{\mu}^{\pi}\xi$  =  $1.00084 \pm 0.00029$  (stat.)<sup>+0.00165</sup> (syst.), in good agreement with the standard model prediction of  $P_{\mu}^{\pi} = \xi = 1$ . Our result is a factor of 7 more precise than the pre-TWIST value, setting new limits in left-right symmetric electroweak extensions to the standard model.

### **Why**  $W_R$ **(300 GeV**,  $\zeta = -0.038$ ) **wasn't detected at FermiLab and CERN?**

#### **Why**  $W_R$ **(300GeV,**  $\zeta = -0.038$ **) wasn't detected at FermiLab?**



Search for  $W'$  bosons decaying to an electron and a neutrino with the D0 detector

 $\mathbf{S}$ 

#### **Estimation of the contribution of <b>W**  $R(300 \text{ GeV}, \zeta = -0.038)$  to the ATLAS collaboration data

This analysis was done in HEPD of PNPI





#### The effect of the right boson  $304^{+28}_{-22}$  GeV with a mixing angle of - 0.038 **is at the level of systematic calculation errors.**

#### **With WR**



#### **wo WR**



#### **Conclusion from the analysis of experiments at colliders**

**The results of our work**

$$
M_{W_R} = 304^{+28}_{-22} \text{ GeV}
$$
  

$$
\zeta = -0.038 \pm 0.014.
$$

**do not contradict the results of experiments at colliders**

# **Prospects for neutron decay experiments**

#### **Project of the installation for measuring neutrino asymmetry at the PIK reactor**

There is a possibility of further increasing the accuracy of measurements in neutron decay. For example, the PNPI NRC KI project "Neutron Beta Decay" for the PIK reactor is aimed at this [29-31], in which it is planned to use a superconducting solenoid with a long flight base for neutron decay in order to increase the statistics of decay events and with a magnetic mirror-collimator to isolate the direction of electron emission. It is a development of the PNPI RAS experiment of 1998 [19], in which it is planned to achieve a relative measurement accuracy of 10-3 for neutrino and electron decay asymmetries.







#### Testing of the installation at NIIEFA 31.05.24. A current of 1050 A was introduced into the superconducting solenoid.



#### **Project of the installation for measuring neutrino asymmetry at the PIK reactor**

**Relative measurement accuracy of 10-3 for neutrino and electron decay asymmetries**



#### **Increasing the measurement accuracy by 3 times can already provide an answer to the question posed.**

#### **Separate detection of protons and electrons**





**1. As a result of the analysis, it was established that there are indications of the existence of a right vector boson with mass**   $= 304^{+28}_{-22}$  GeV **and the mixing angle with**  $W_L$ :  $\zeta = -0.038 \pm 0.014$ .

**2. This result does not contradict experiments at colliders.**

**3. However, it is necessary to conduct even more precise measurements of neutron decay and its theoretical analysis.**

**4. Measuring neutrino asymmetry with an accuracy of 10-3 is our goal.**



## **SCIENTIFIC RESEARCH PROGRAM Neutron decay at the reactor PIK**

