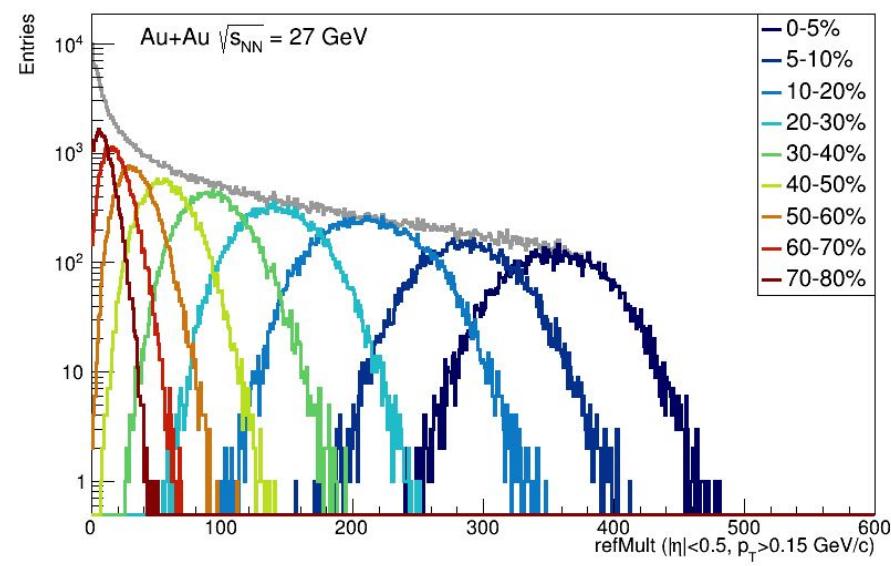
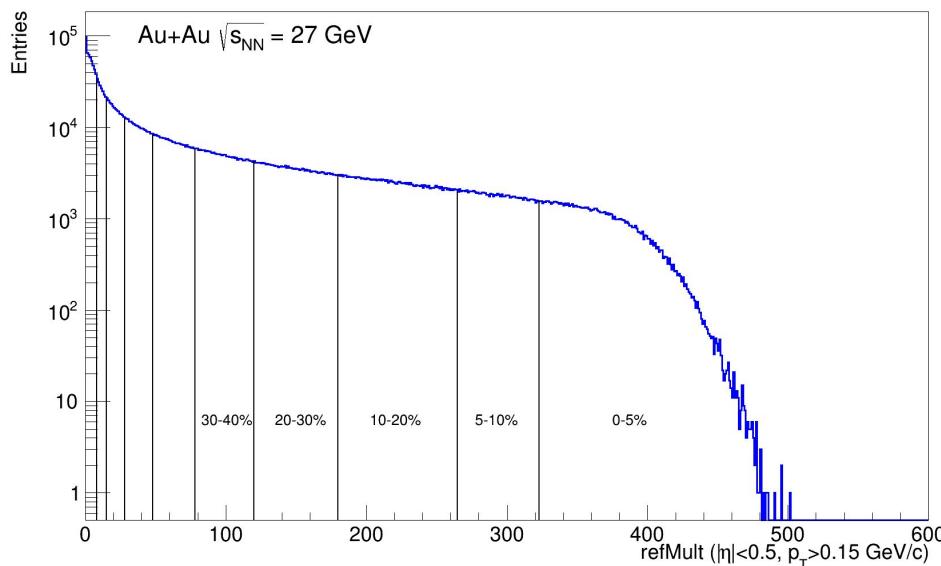


$\pi/K/p$ spectra, Au+Au, $\sqrt{s_{NN}} = 27$ GeV

UrQMD

Cuts for centrality calculation:

- π^\pm, K^\pm, p ($p\bar{p}$)
- $p_T > 0.15 \text{ GeV}/c$
- $|\eta| < 0.5$



Au+Au, 27 GeV

Centrality, %	Multiplicity	Percent in this bin	Cumulative percent
0-5	$322 \leq \text{mult} < 502$	0.0500505	0.0500505
5-10	$264 \leq \text{mult} < 322$	0.0495357	0.0995862
10-20	$179 \leq \text{mult} < 264$	0.100235	0.199822
20-30	$119 \leq \text{mult} < 179$	0.100627	0.300448
30-40	$77 \leq \text{mult} < 119$	0.0987935	0.399242
40-50	$47 \leq \text{mult} < 77$	0.100271	0.499513
50-60	$27 \leq \text{mult} < 47$	0.0990788	0.598592
60-70	$14 \leq \text{mult} < 27$	0.103341	0.701933
70-80	$7 \leq \text{mult} < 14$	0.0972113	0.799144

Statistics: $\sim 2M$

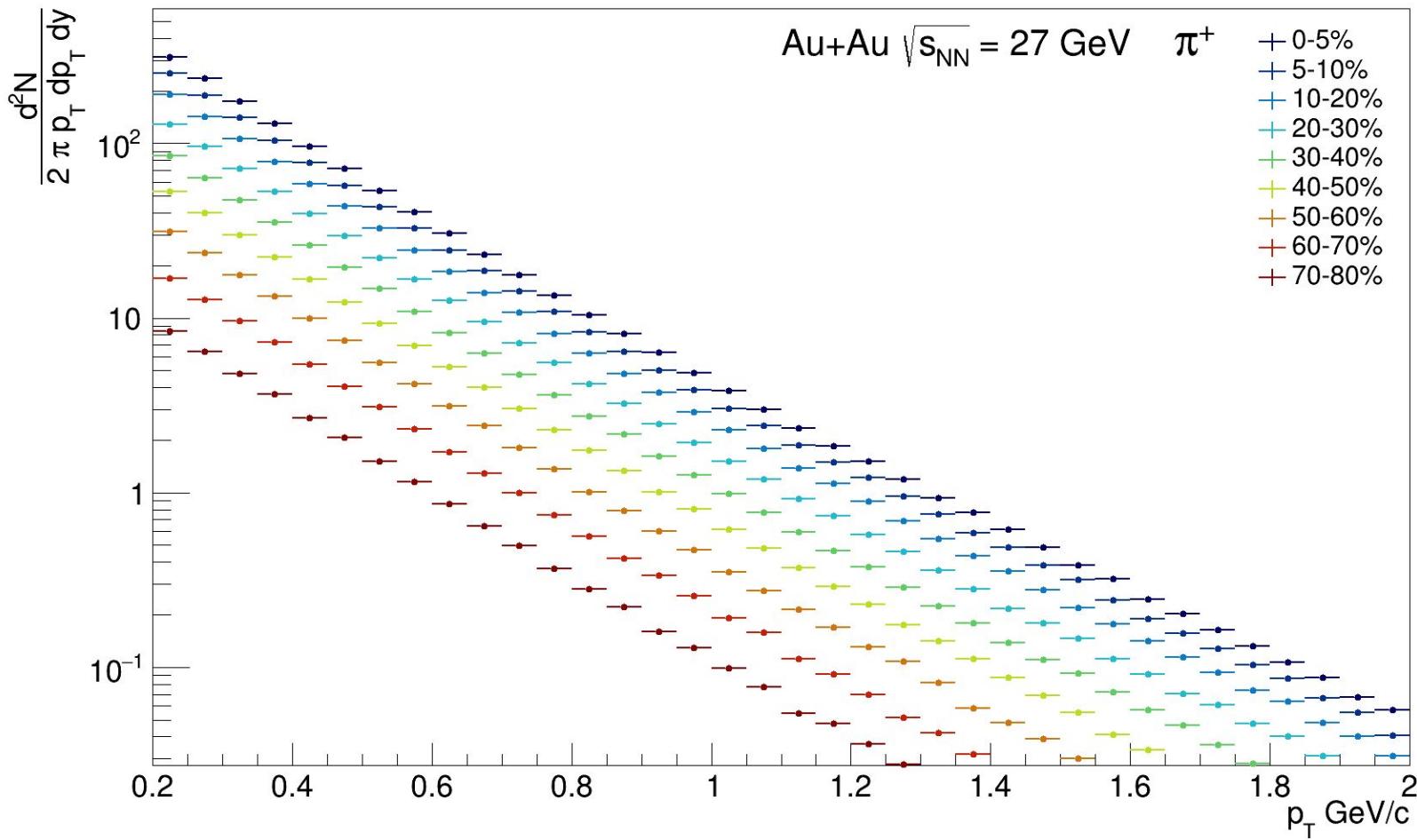
Track cuts:

- PDG ($\pi^\pm = \pm 211$, $K^\pm = \pm 321$, p (p-bar) = ± 2212)
- $|y| < 0.1$,
- $|\eta| < 0.5$
- $p_T > 0.2$ GeV/c

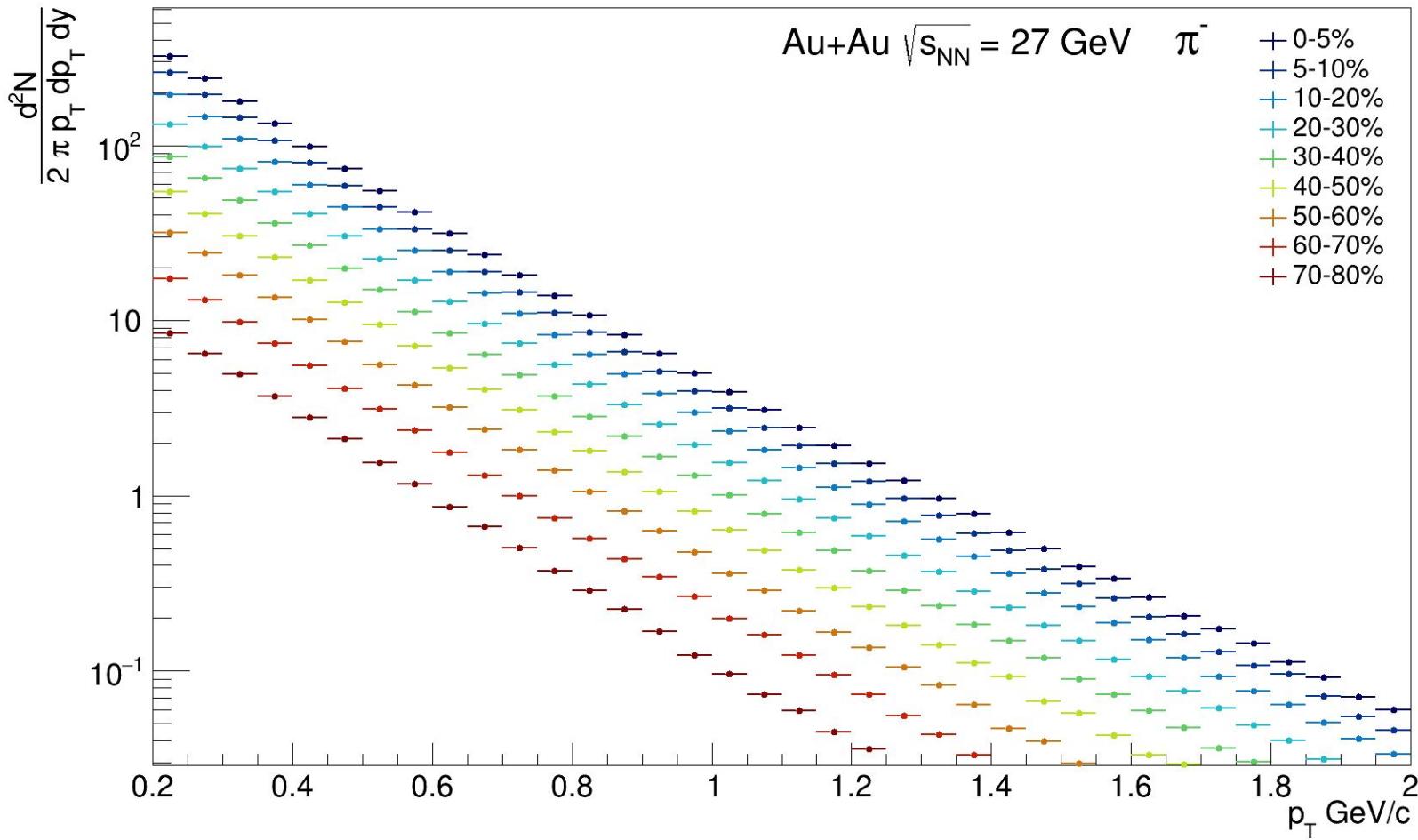
Bin width for spectra: 50 MeV/c

Centrality was calculated using multiplicity.

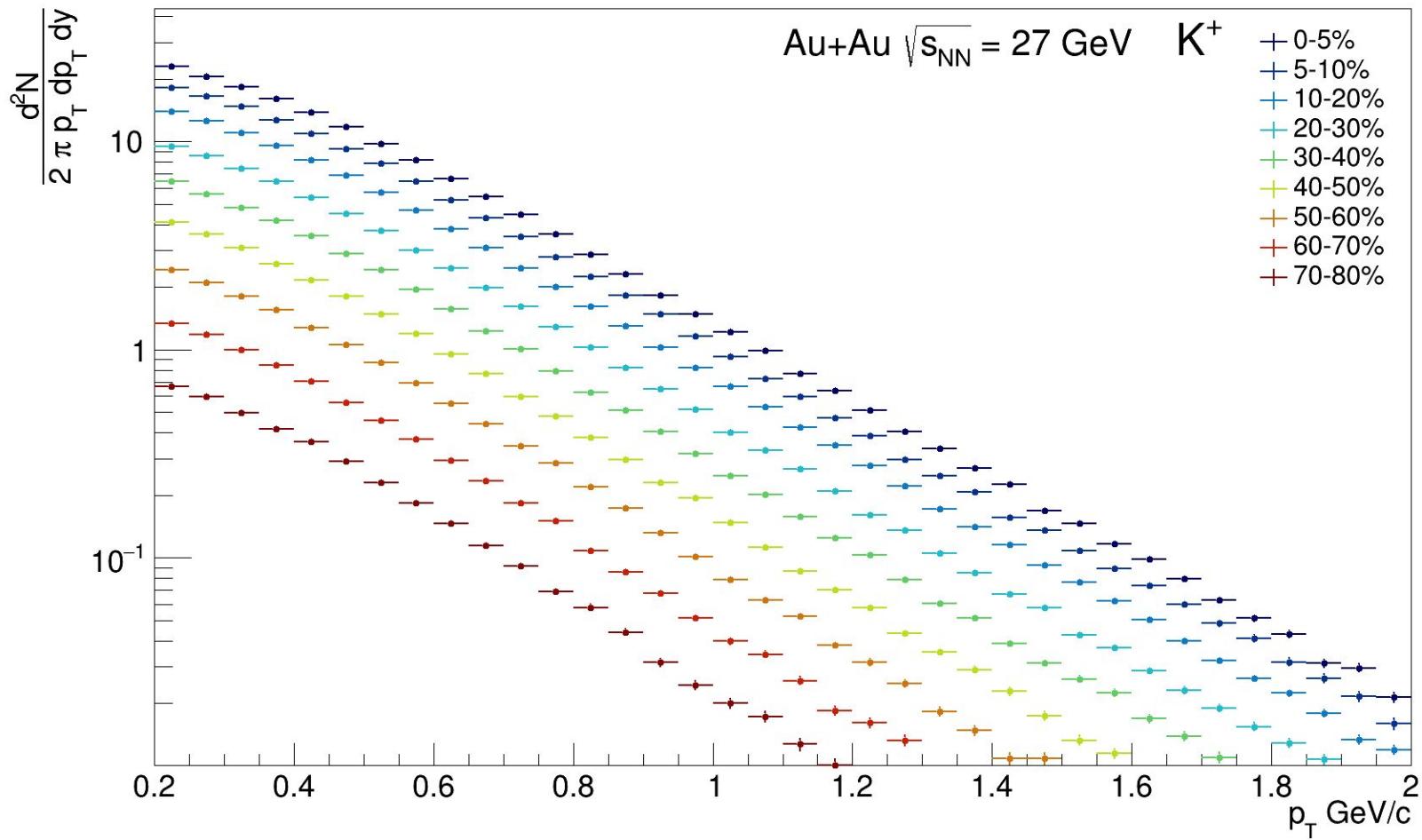
Centrality dependence of π^+ spectra at $\sqrt{s_{\text{NN}}} = 27$ GeV



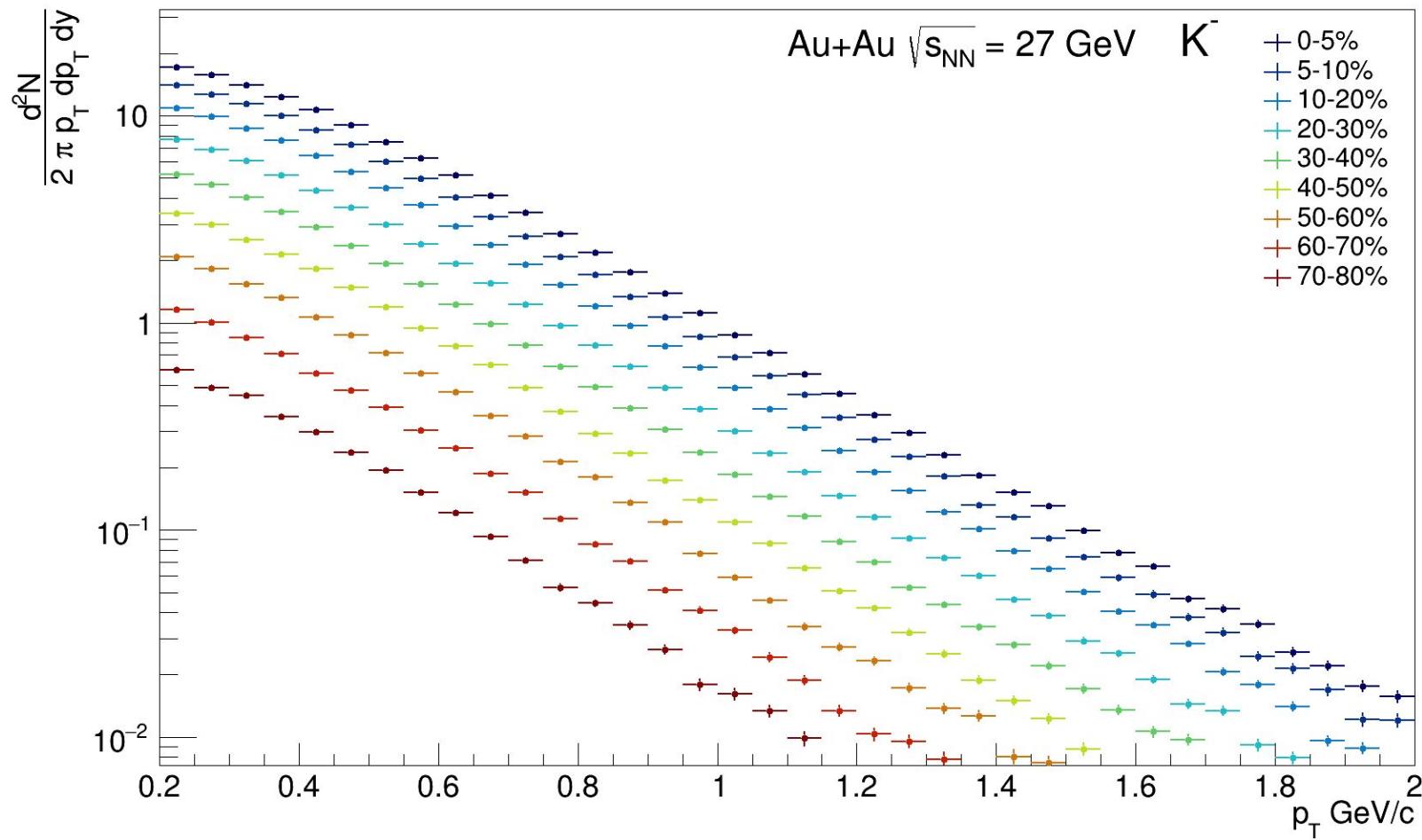
Centrality dependence of π^- spectra at $\sqrt{s_{\text{NN}}} = 27 \text{ GeV}$



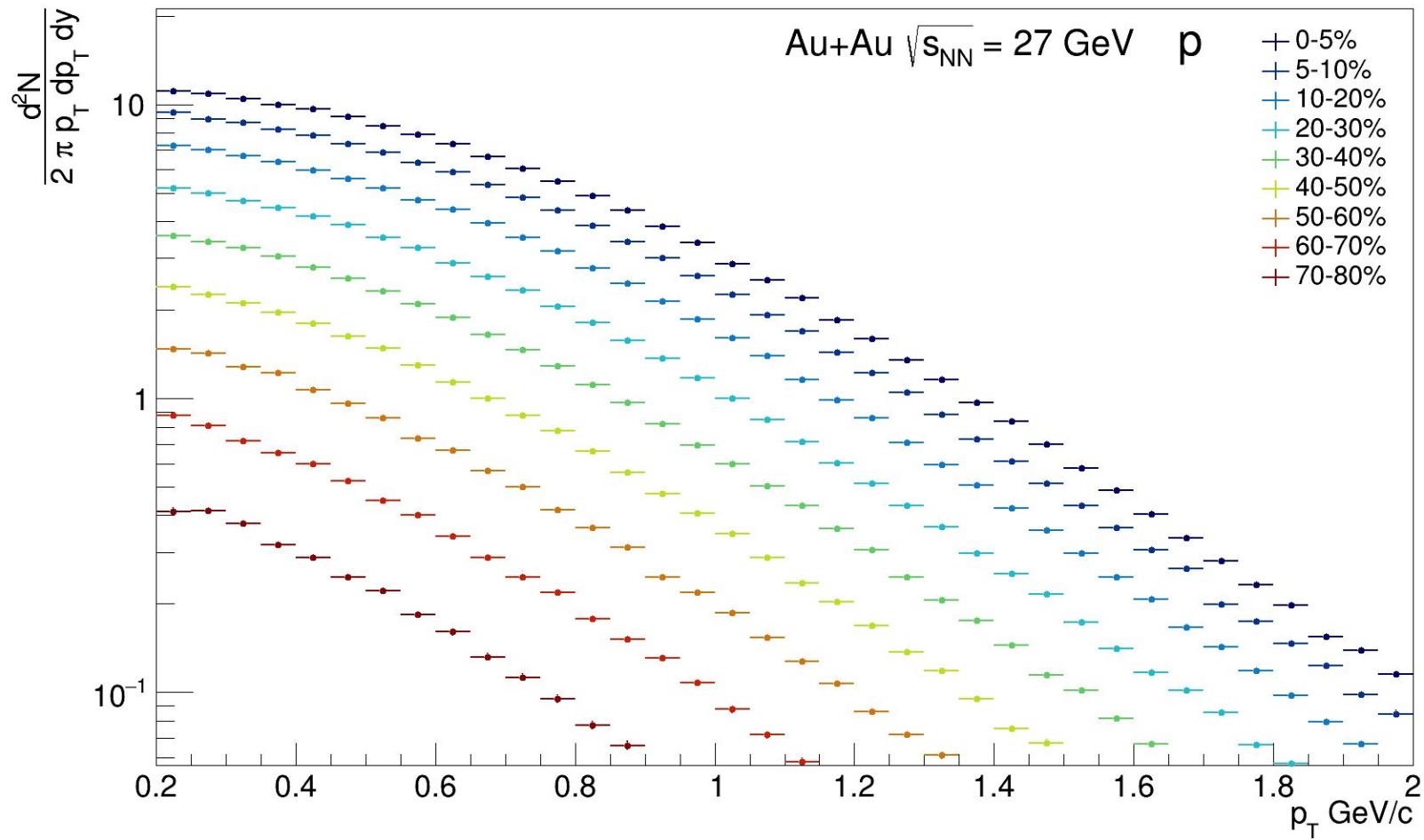
Centrality dependence of K⁺ spectra at $\sqrt{s_{NN}} = 27$ GeV



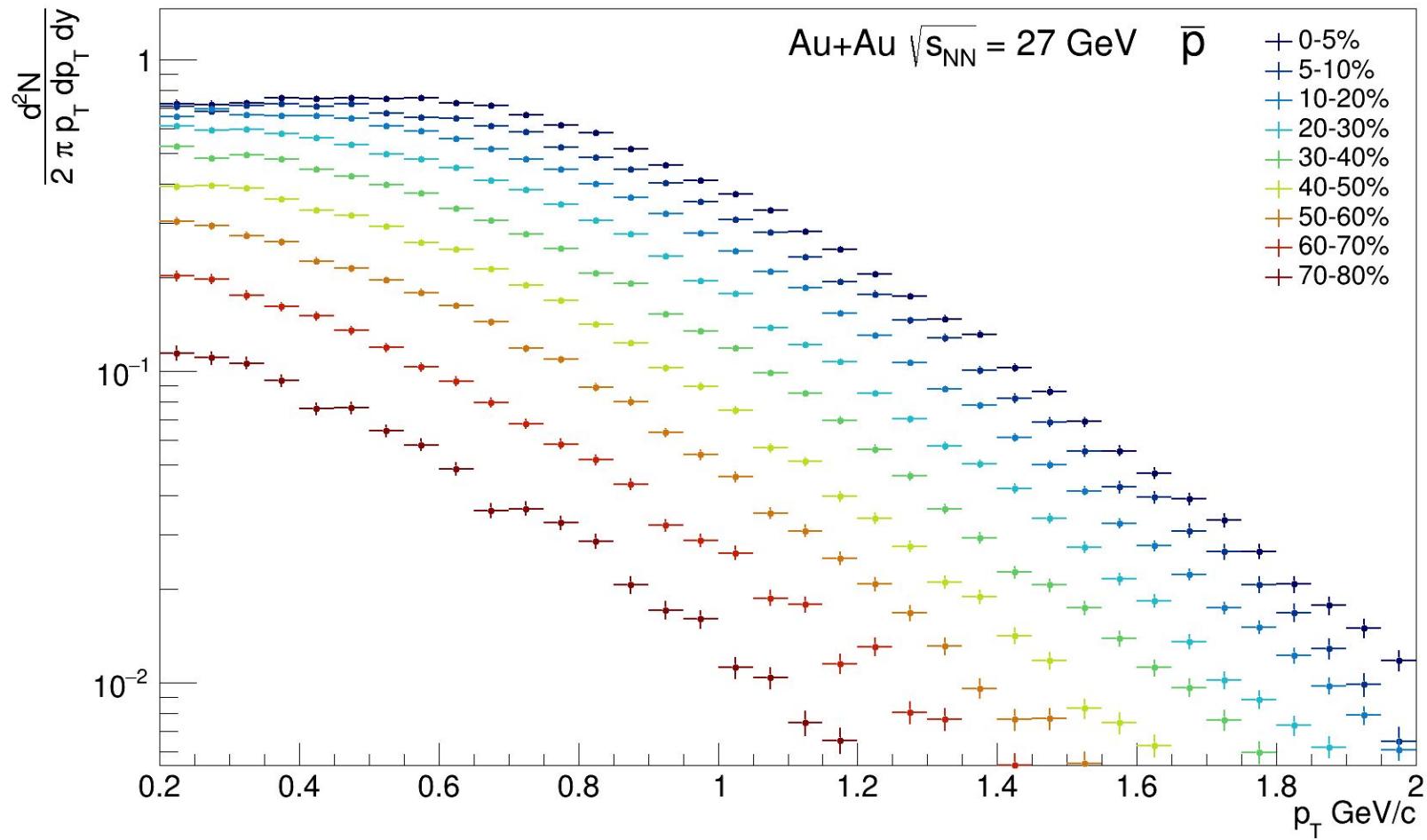
Centrality dependence of K⁻ spectra at $\sqrt{s_{NN}} = 27$ GeV



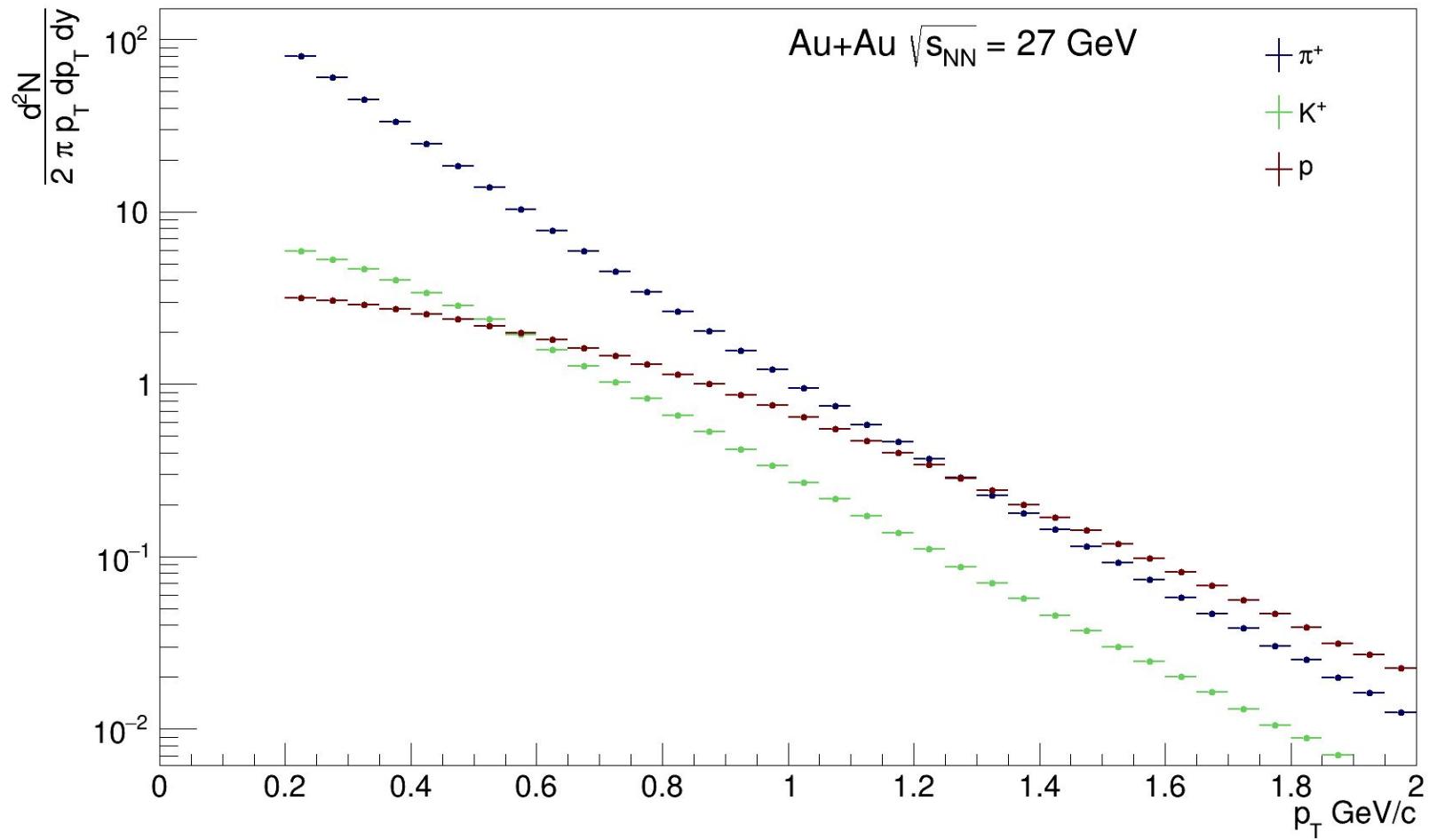
Centrality dependence of p spectra at $\sqrt{s_{NN}} = 27$ GeV



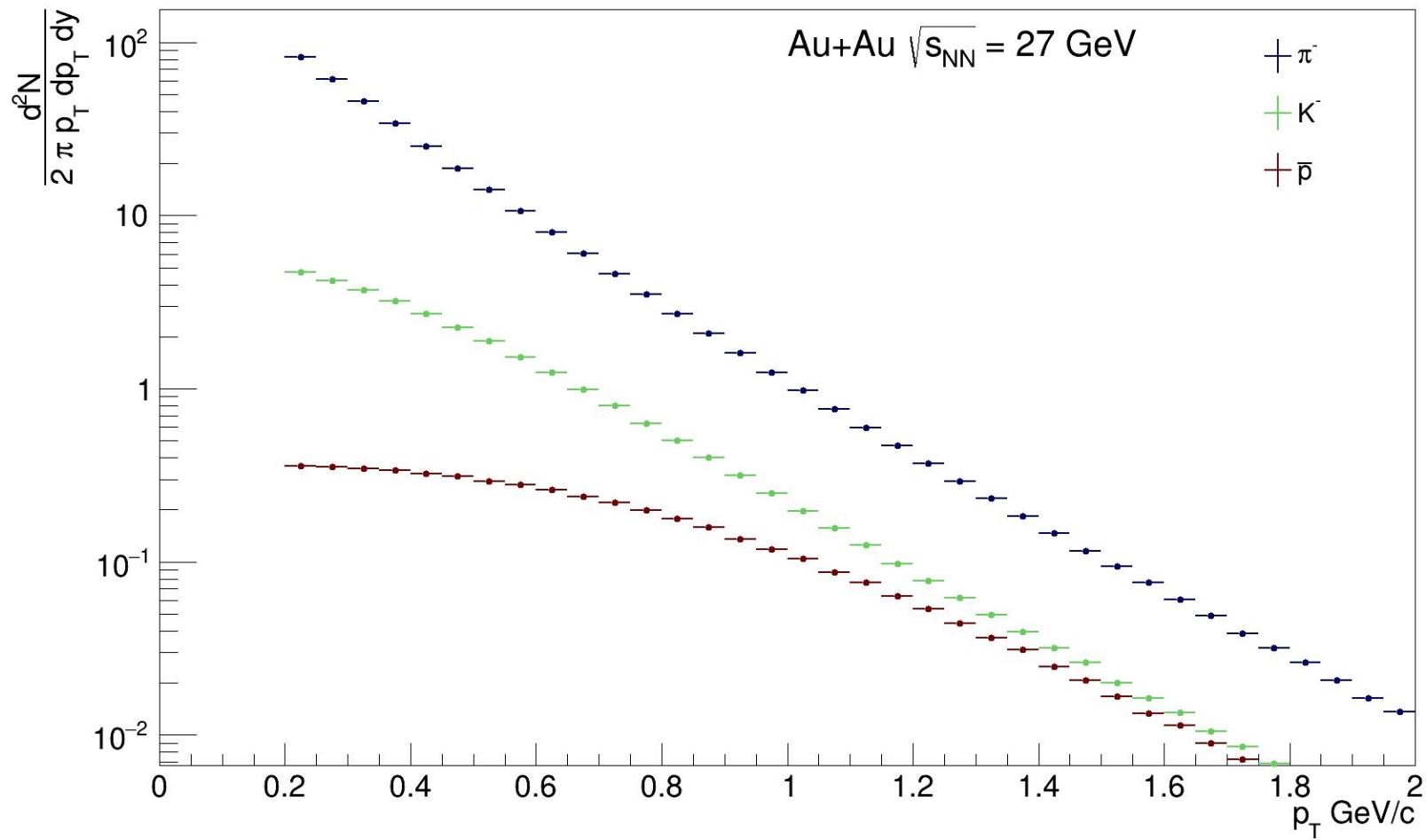
Centrality dependence of p-bar spectra at $\sqrt{s_{NN}} = 27$ GeV



$\pi^+/\text{K}^+/\text{p}$ spectra at $\sqrt{s_{\text{NN}}} = 27$ GeV in integrated centrality range



$\pi^-/\bar{K}/\bar{p}$ -bar spectra at $\sqrt{s_{NN}} = 27$ GeV in integrated centrality range



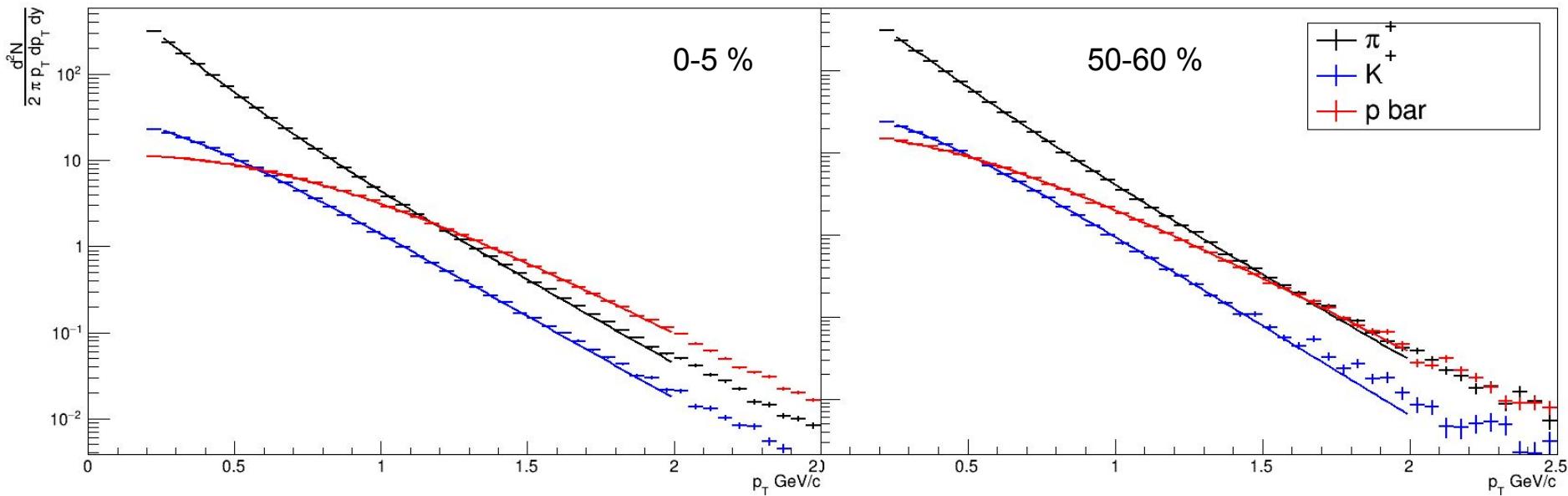
blast wave fit

$$\frac{dN}{p_T dp_T} \propto \int_0^R r dr m_T I_0 \left(\frac{p_T \sinh \rho(r)}{T_{\text{kin}}} \right) \times K_1 \left(\frac{m_T \cosh \rho(r)}{T_{\text{kin}}} \right),$$

where m_T is the transverse mass of a hadron, $\rho(r) = \tanh^{-1} \beta$, and I_0 and K_1 are the modified Bessel functions. We use a radial flow velocity profile of the form

$$\beta = \beta_S (r/R)^n, \quad (14)$$

where β_S is the surface velocity, r/R is the relative radial position in the thermal source, and n is the exponent of flow velocity profile. Average transverse radial flow velocity $\langle \beta \rangle$ can then be obtained from $\langle \beta \rangle = \frac{2}{2+n} \beta_S$.



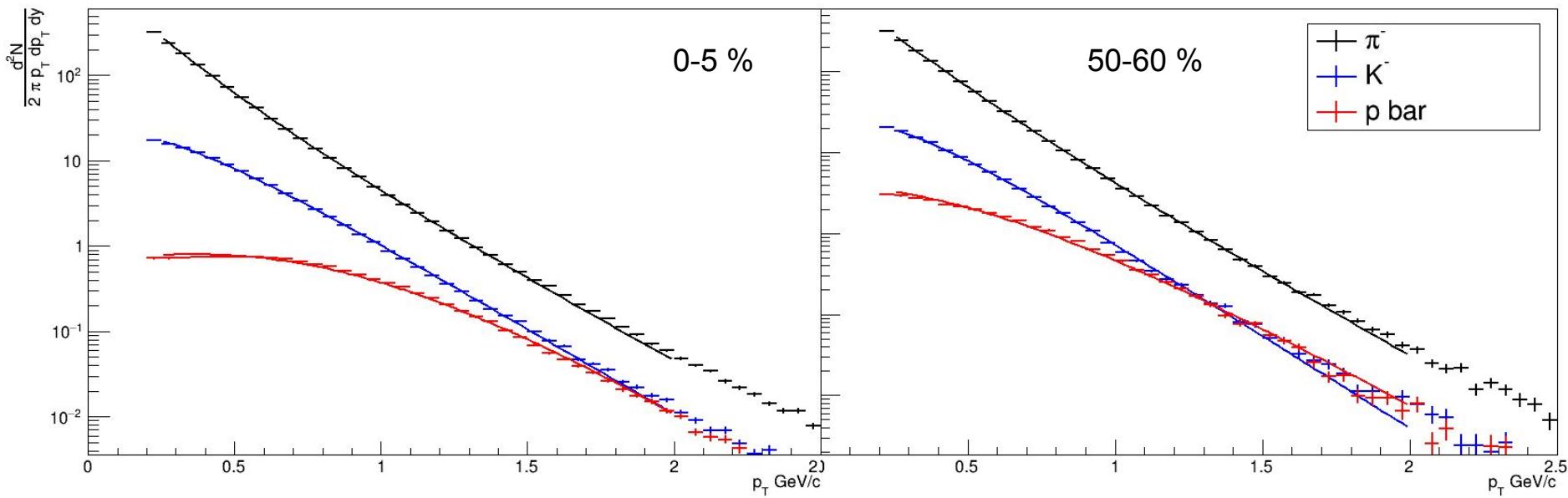
blast wave fit

$$\frac{dN}{p_T dp_T} \propto \int_0^R r dr m_T I_0 \left(\frac{p_T \sinh \rho(r)}{T_{\text{kin}}} \right) \times K_1 \left(\frac{m_T \cosh \rho(r)}{T_{\text{kin}}} \right),$$

where m_T is the transverse mass of a hadron, $\rho(r) = \tanh^{-1} \beta$, and I_0 and K_1 are the modified Bessel functions. We use a radial flow velocity profile of the form

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	Centr	$\langle \beta \rangle$	T_{kin} , MeV
Pion +	0-5 %	0.53	75.3
K +		0.32	140.6
p		0.45	111.8
Pion -		0.53	73.8
K -		0.24	156.6
p Bar		0.53	72.5
Pion +	50-60%	0.42	102.1
K +		0.31	119.9
p		0.26	153.7
Pion -		0.44	98.5
K -		0.31	116.7
p Bar		0.10	184.5

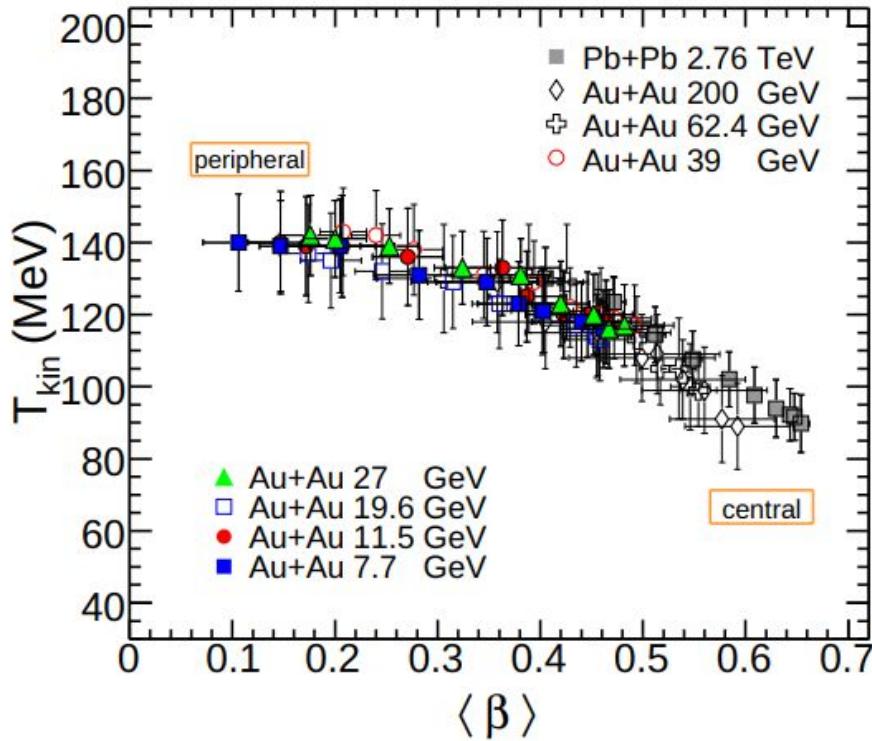


FIG. 37: (Color online) Variation of T_{kin} with $\langle \beta \rangle$ for different energies and centralities. The centrality increases from left to right for a given energy. The data points other than BES energies are taken from Refs. [43, 66]. Uncertainties represent systematic uncertainties.